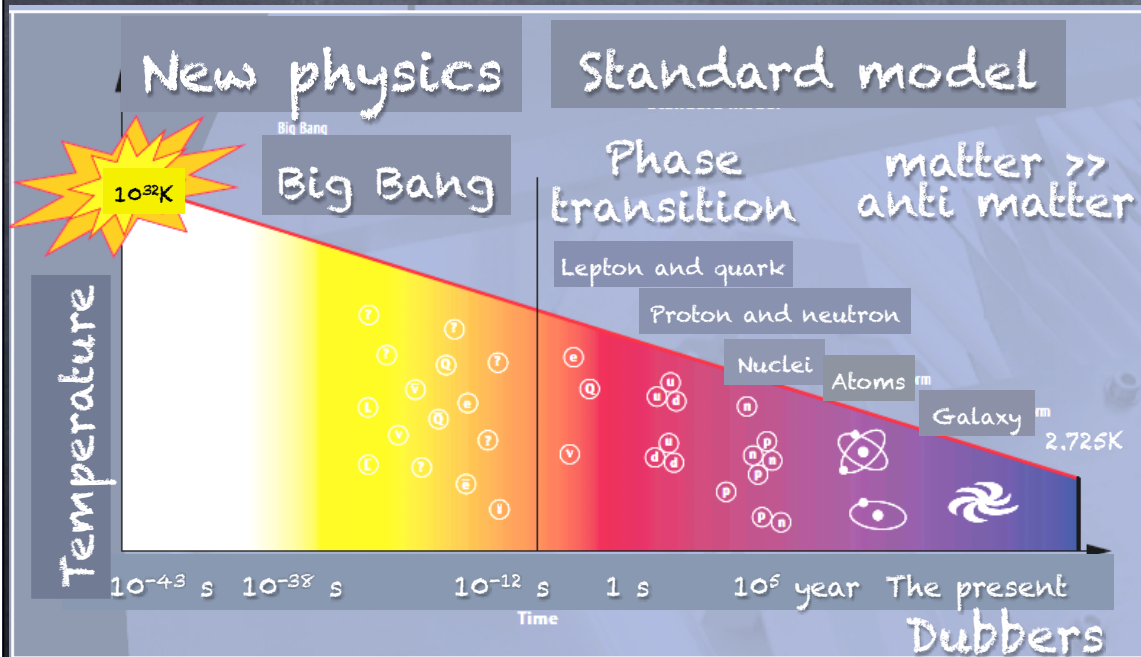
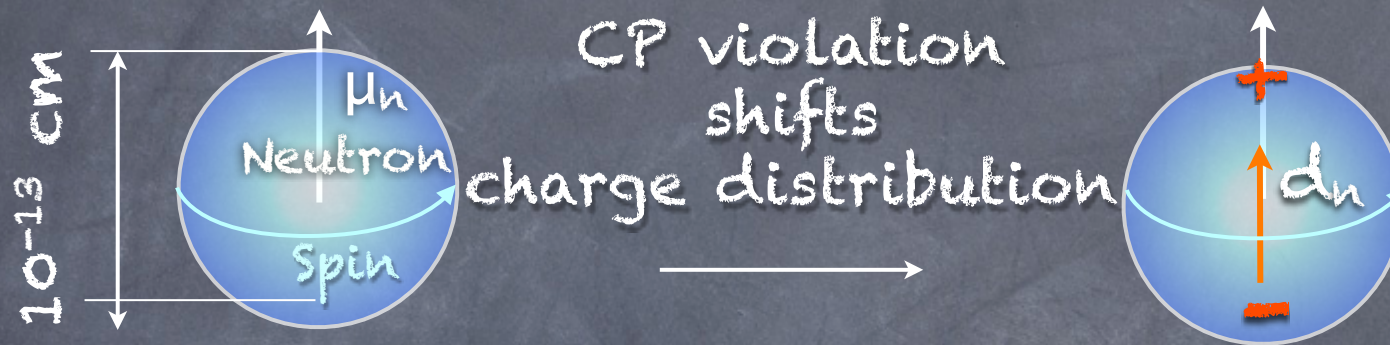


A new EDM measurement with a new generation UCN source

May 22, 2012, ISINN20



Mystery of baryogenesis :
CP violation creates matter
but
Standard model
can't explain the baryon
asymmetry

KEK-RCNP EDM measurement

Y. Masuda, S. Jeong, Y. Watanabe, T. Adachi, S. Kawasaki, K. Matsuta,
M. Mihara, K. Hatanaka, R. Matsumiya, and K. Asahi

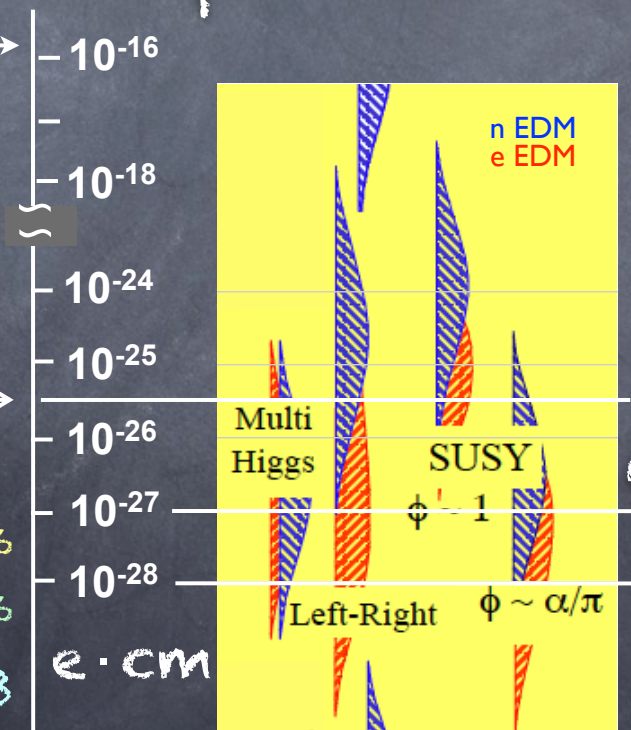
New physics
predicts EDM

Earth's field
($\mu_n H_0 / E$)



EDM upper limit

statistical 1.5×10^{-26}
systematic 0.7×10^{-26}
related with B



statistical
high density UCN



10^{-27} e · cm



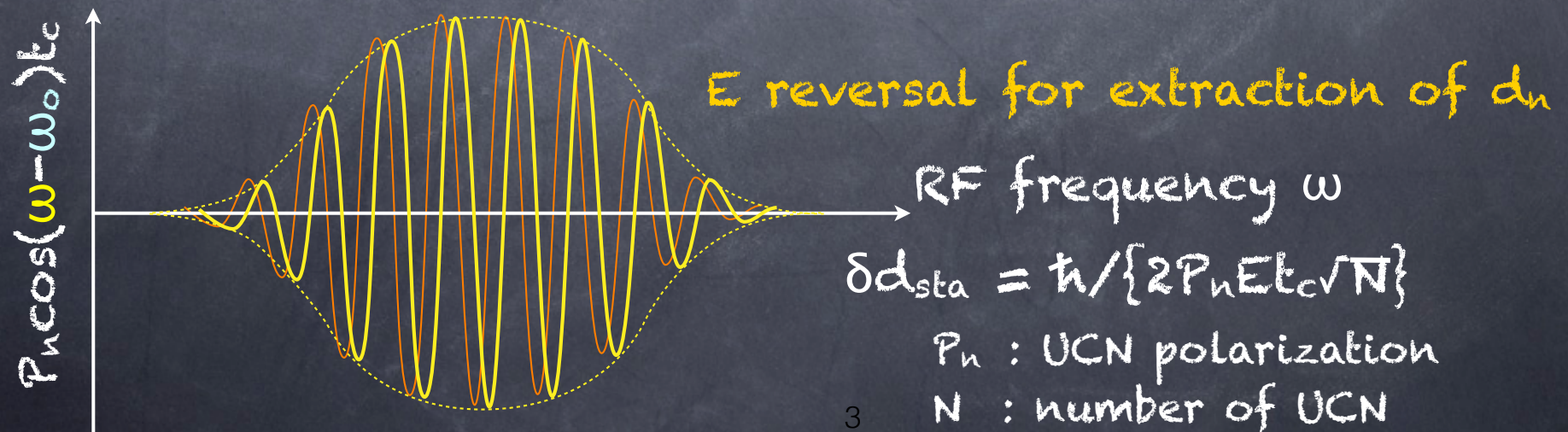
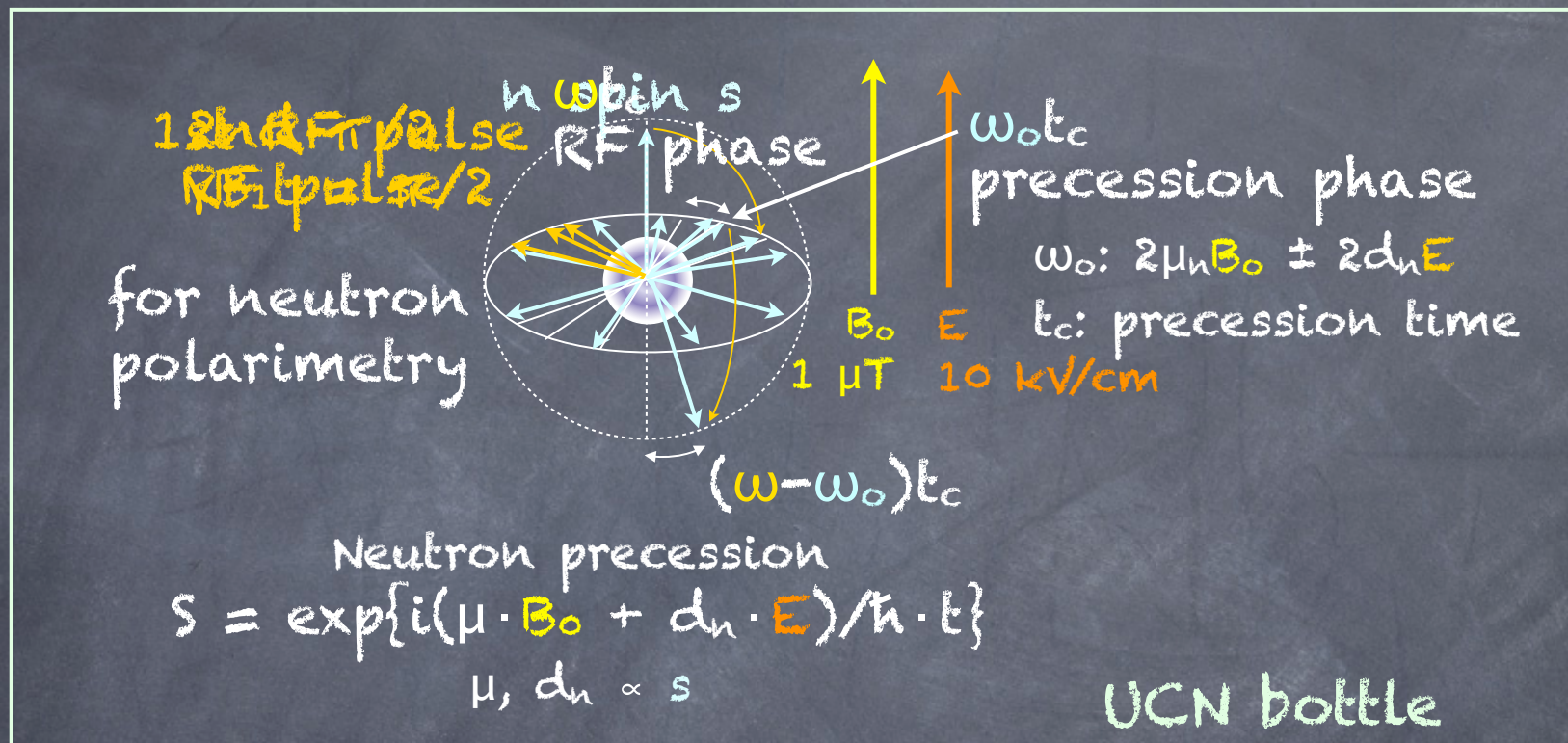
10^{-28} e · cm

Ours

systematic
new magnetometer

10^{-28} e · cm

EDM measurement



Ramsey resonance apparatus

Spherical coil

EDM cell

Door valve

$E_c = 90 \text{ neV}$

$\pi/2$ RF coil

UCN valve

Spin flipper

Magnetized iron foil

Rotary valve

UCN detector

UCN production for EDM

$$T_s = 81 \text{ s}$$

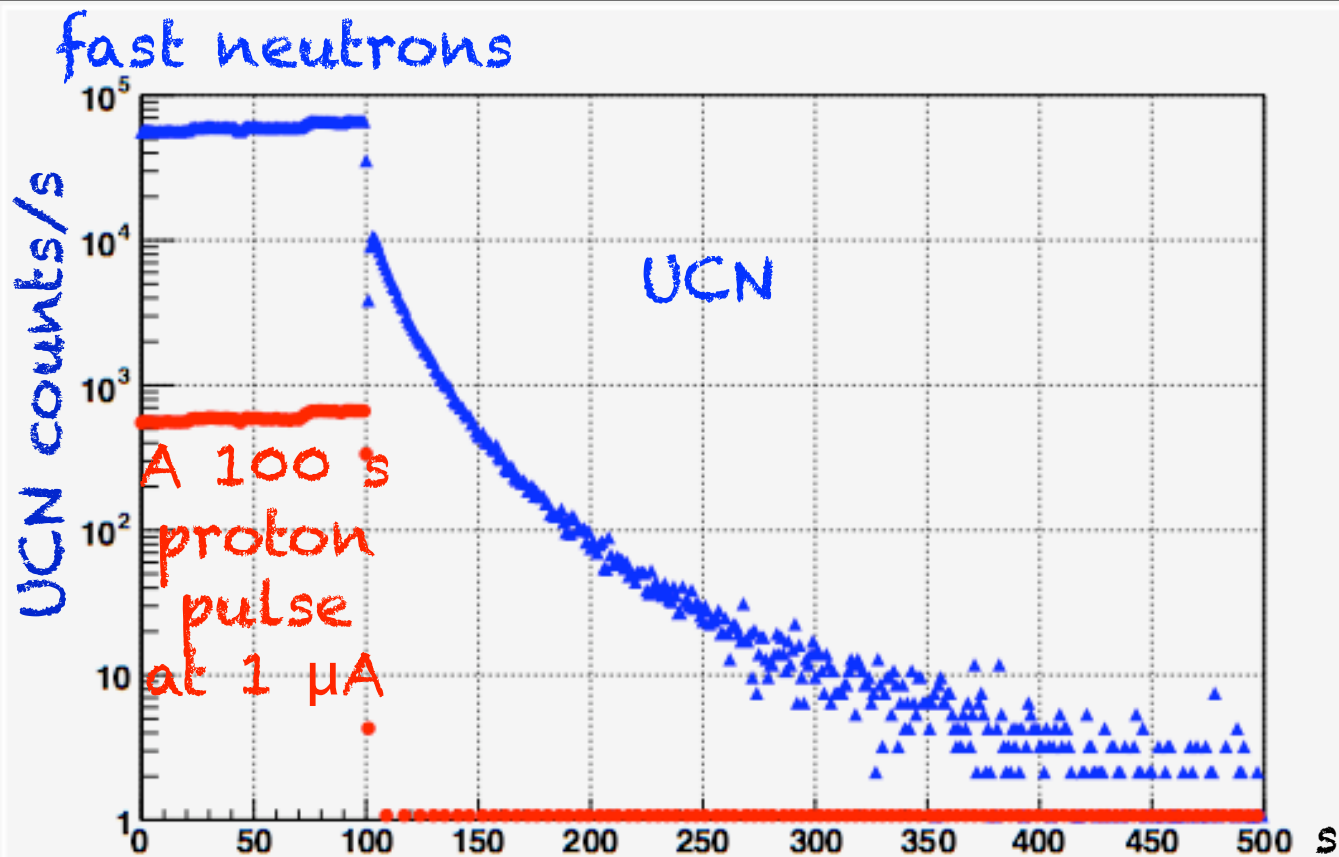
Phys. Rev. Lett. 108(2012)134801

240 s 390W proton irradiation

26 UCN/cm³ at $E_c = 90 \text{ neV}$, $P = 4 \text{ UCN/cm}^3 \cdot \text{s}$ at $E_c = 210 \text{ neV}$

$$\rho_{\text{UCN}} \propto E_c^{3/2}$$

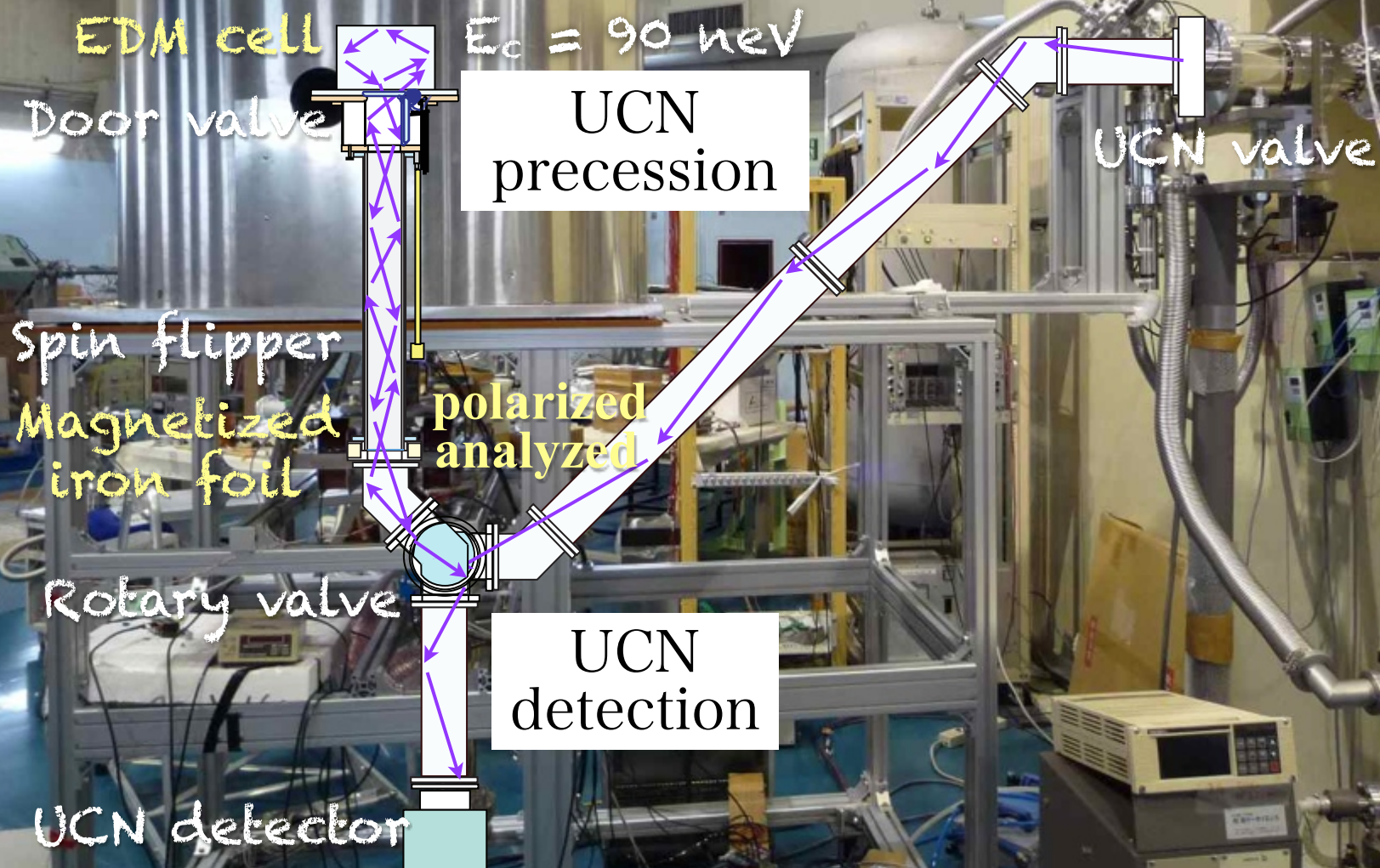
will be increased in the new source



UCN detector

He
ulator

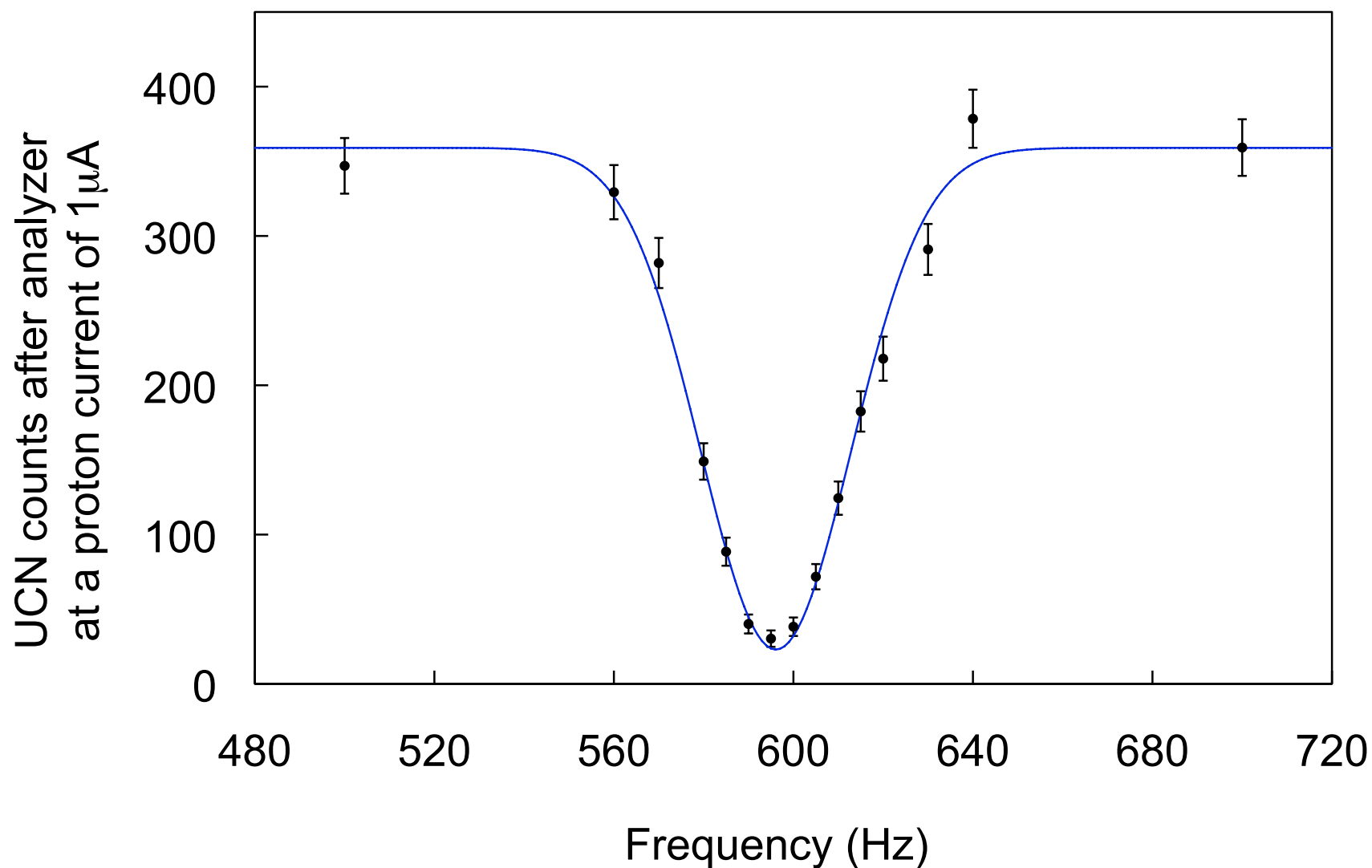
Ramsey resonance experiment



UCN NMR in EDM cell



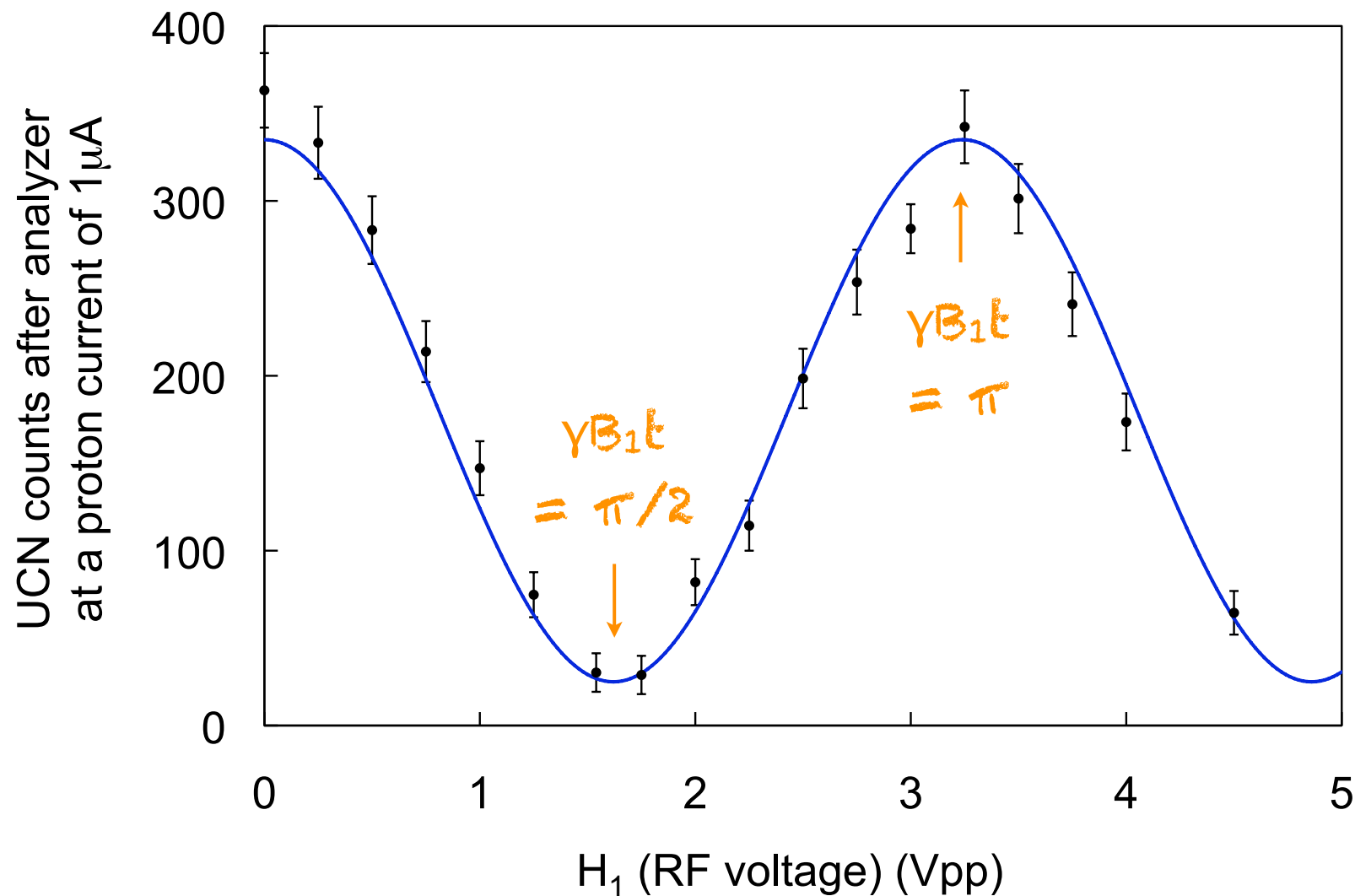
two coherent $\pi/2$ RF pulses



UCN spin rotation by $\gamma B_1 t$

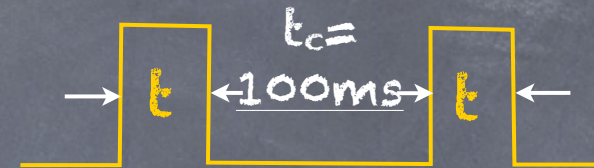


two coherent $\pi/2$ RF pulses

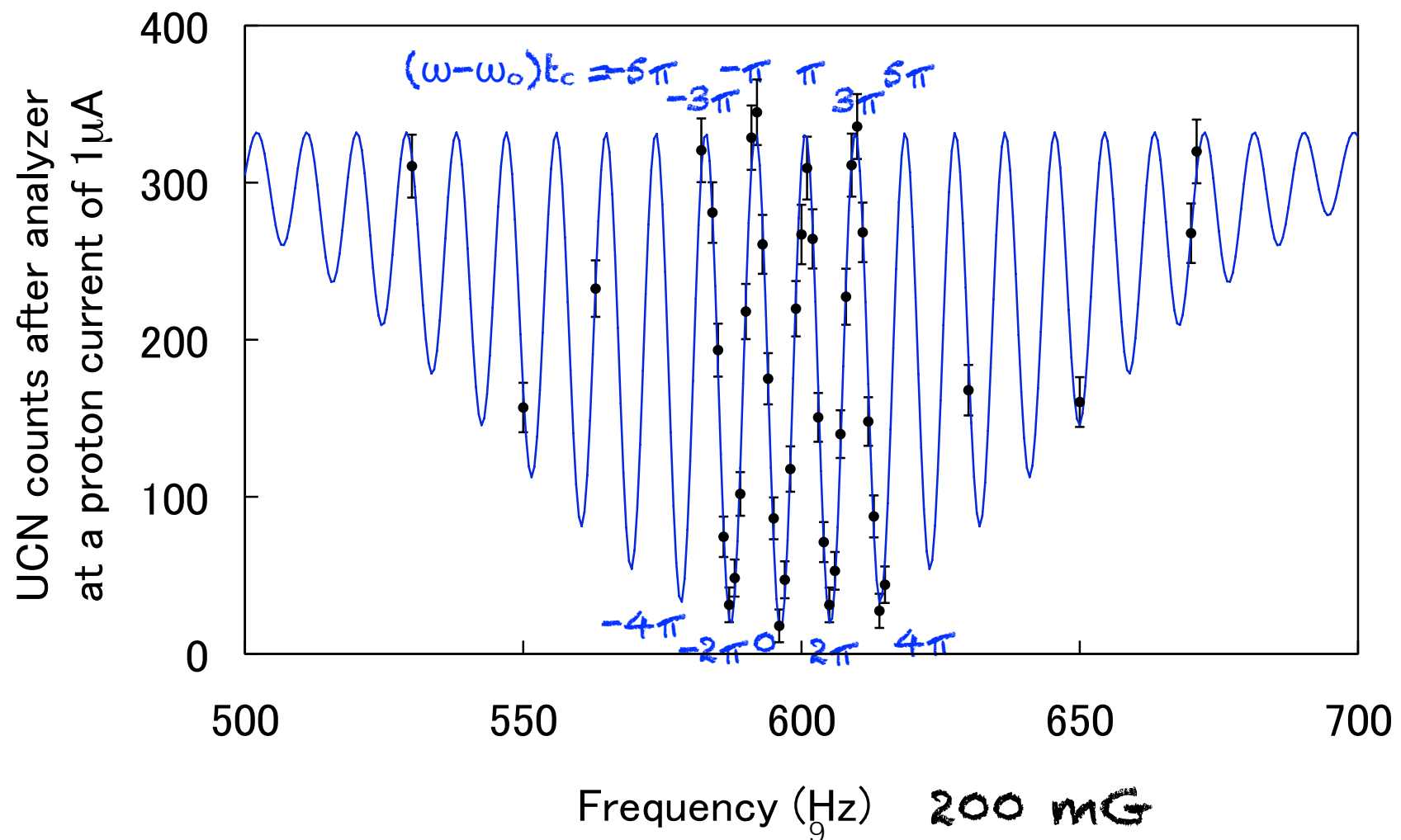


Ramsey resonance

visibility $\alpha =$
 $(N_{\max} - N_{\min}) / (N_{\max} + N_{\min}) = 0.9$

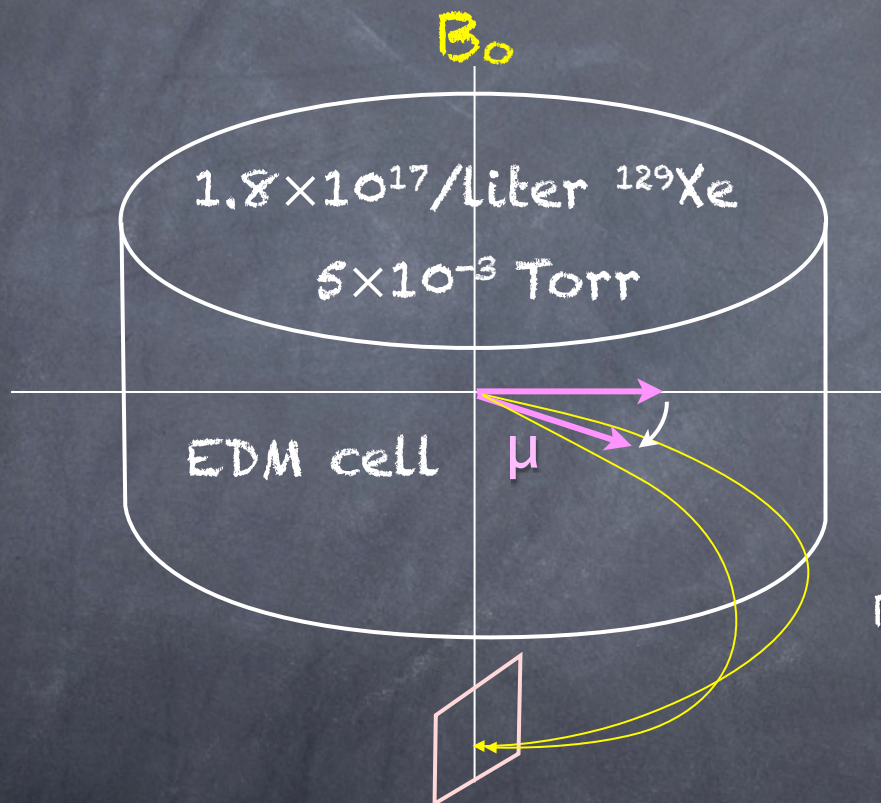


two coherent $\pi/2$ RF pulses



^{129}Xe spin magnetometer

Upon E reversal



^{129}Xe magnetization

$$P_{^{129}\text{Xe}} = 0.5$$

$$B = 0.4 \text{ pT}$$

in the EDM cell

Rb- ^{129}Xe
van der Waals
molecule

$$B = 0.15 \text{ pT}$$

at 0.1 m

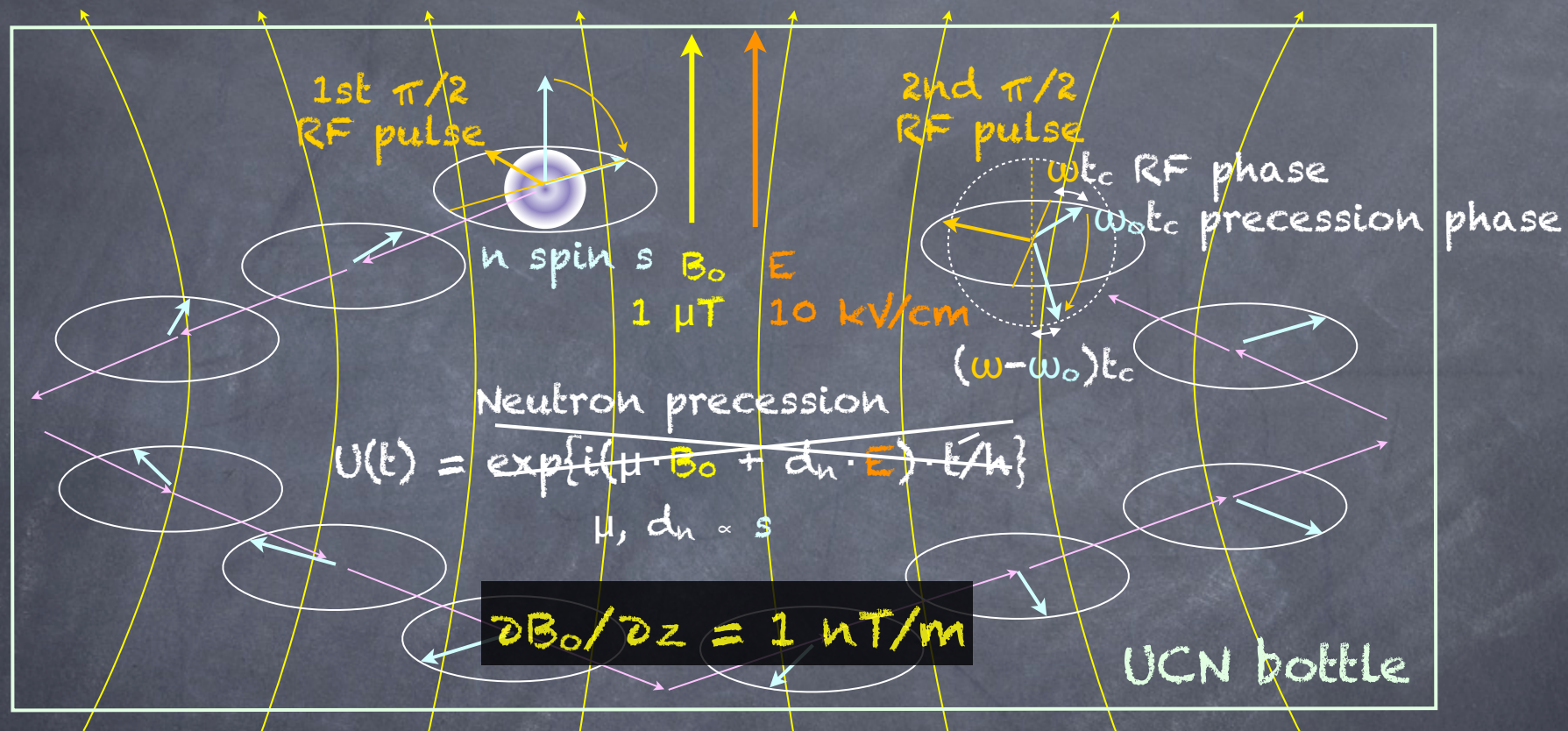
from the center

SQUID 1fT, $5\mu\Phi_0/\sqrt{\text{Hz}}$

B_0 monitor by nuclear spin precession

Isotope	J_N	$g(\gamma_N = g\mu_N/\hbar)$	σ_a at 2200 m/s	ρ for $\tau = 1/(\sigma_a \rho v)$ ≈ 500 s
^{129}Xe Ours	1/2	<u>-0.777</u>	<u>21 b</u>	$2.5 \times 10^{14}/\text{cc}$, Rb-Xe
^{199}Hg ILL	1/2	0.5026	2150 b	$(3 \times 10^{10}/\text{cc})$, photon)
n cryoEDM	1/2	-1.913		
^3He SNS	1/2	-2.128	5333 b	$10^{12}/\text{cc}$, SQUID
^{133}Cs PSI	7/2	2.579	29 b	

Geometric phase effect (GPE)



Geometric phases arise from the transverse fields,

$$B_{or} = (\partial B_0 / \partial z) r / 2, B_v = E \times v / c^2$$

Pendlebury, Phys. Rev. A70(2004)032102.
 Lamoreaux, Phys. Rev. A71(2005)052115.

Effect of time dependent interaction

$$U(t) = \exp(-iH_0 t / \hbar)$$

$$H_0 = -\boldsymbol{\mu} \cdot \mathbf{B}_0 - d_n \cdot \mathbf{E}$$

Phys.Lett. A376(2012)1347

$$H = H_0 + V(t)$$

$$V(t) = -\boldsymbol{\mu} \cdot \mathbf{B}_{xy}(t) = -\gamma s \cdot (\mathbf{B}_v(t) + \mathbf{B}_{0r}(t))$$

$$\mathbf{B}_v = \mathbf{E} \times \mathbf{v} / c^2 \quad \mathbf{B}_{0r} = -(\partial \mathbf{B}_{0z} / \partial z) \mathbf{r} / 2$$

$$U_I(t) = 1 + \left(\frac{-i}{\hbar}\right) \int_0^t dt' V_I(t') + \left(\frac{-i}{\hbar}\right)^2 \int_0^t dt' \int_0^{t'} dt'' V_I(t') V_I(t'') + \dots$$

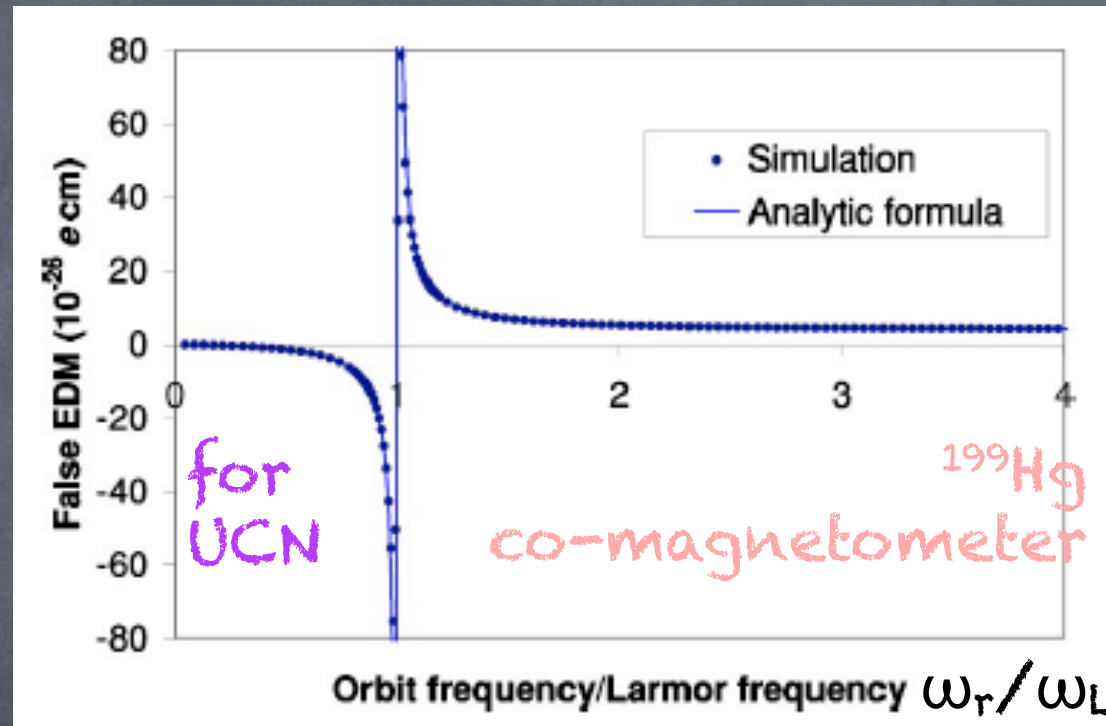
$$V_I(t) = \exp(iH_0 t / \hbar) \{-\boldsymbol{\mu} \cdot \mathbf{B}_{xy}(t)\} \exp(-iH_0 t / \hbar)$$

$$U_I(t) = 1 + \frac{is_z}{\hbar} \frac{1}{4} \gamma^2 \frac{E}{c^2} \frac{\partial \mathbf{B}_{0z}}{\partial z} \int_0^t dt' \int_0^{t'} d\tau \cos(\omega_0 \tau)$$

$$\{\mathbf{x}(t') \mathbf{v}_x(t' - \tau) - \mathbf{x}(t' - \tau) \mathbf{v}_x(t') + \mathbf{y}(t') \mathbf{v}_y(t' - \tau) - \mathbf{y}(t' - \tau) \mathbf{v}_y(t')\}$$

GPE problem

$$\int_0^t d\tau \cos(\omega_0 \tau) \{x(t')v_x(t'-\tau) - x(t'-\tau)v_x(t') + y(t')v_y(t'-\tau) - y(t'-\tau)v_y(t')\}$$



$$d_{\text{afn}} = -\hbar/4 \cdot (\partial B_{0z}/\partial z)/B_{0z}^2 \cdot v_{xy}^2/c^2 \sim 1 \times 10^{-27} \text{ e} \cdot \text{cm}$$

$$d_{\text{afHg}} = \hbar/8 \cdot \gamma_n \gamma_{\text{Hg}} (\partial B_{0z}/\partial z) R^2/c^2 \sim 5 \times 10^{-26} \text{ e} \cdot \text{cm}$$

Pendlebury, Phys. Rev. A70(2004)032102.
Lamoreaux, Phys. Rev. A71(2005)052115.

Solution for GPE problem

$$\int_0^t d\tau \cos(\omega_0 \tau) \{x(t')v_x(t'-\tau) - x(t'-\tau)v_x(t') + y(t')v_y(t'-\tau) - y(t'-\tau)v_y(t')\}$$

^{129}Xe buffer gas

$$\lambda = 1/n\sigma \ll 0.7 \text{ mm}$$

$$n = 1.8 \times 10^{14} / \text{cc at } 5 \times 10^{-3} \text{ torr}$$

$$\sigma_{\text{Xe-Xe}} \gg 838 \text{ \AA}^2$$

random walk



$r(t)$ does not change so much
 $v(t-\tau)$ changes rapidly

$$\langle r(t)v(t-\tau) \rangle \rightarrow \ll 1$$

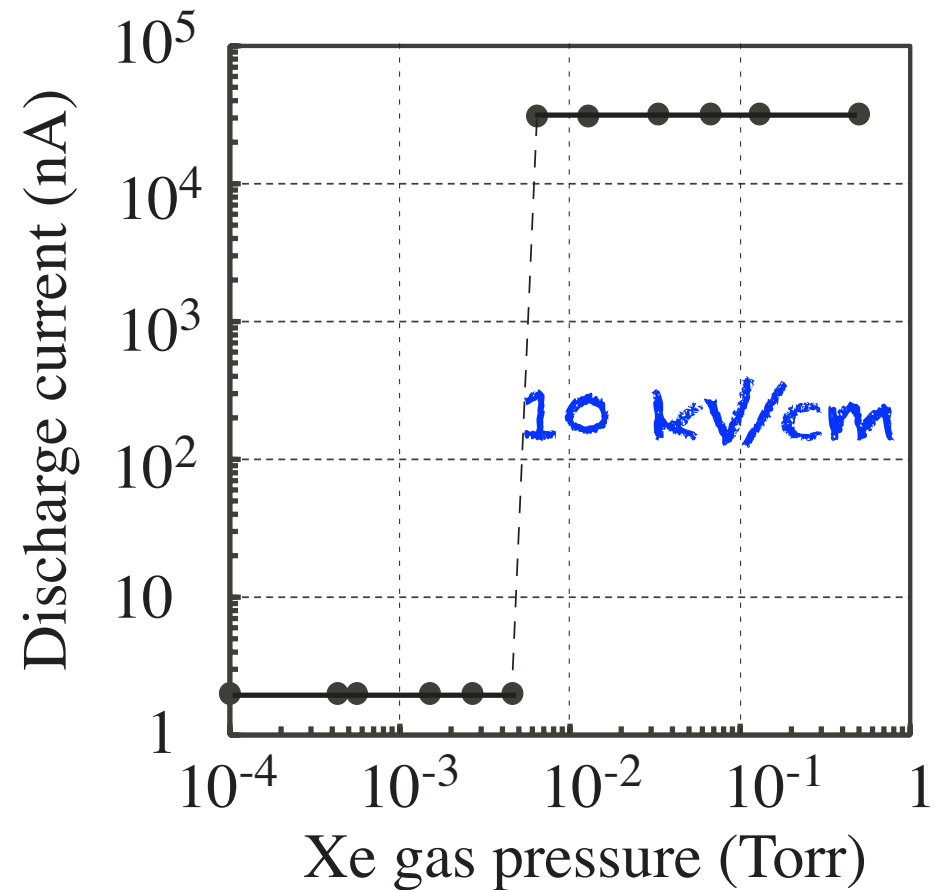
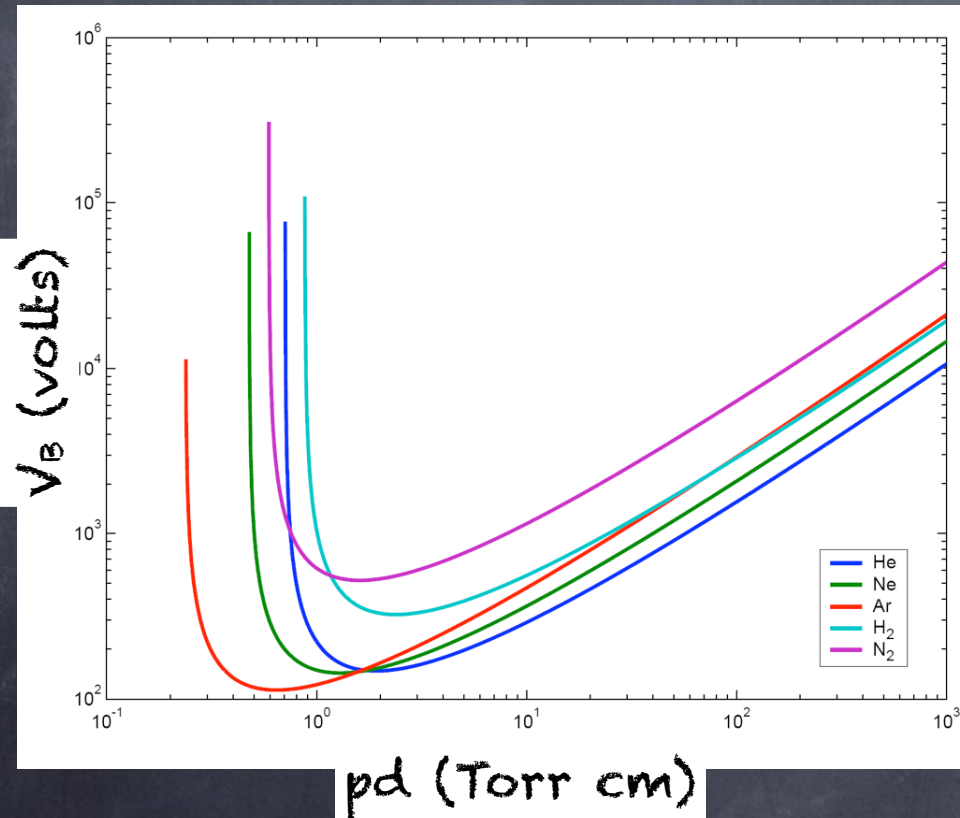
$$d_{\text{afHgn}} \sim \underline{5 \times 10^{-26} \text{ e} \cdot \text{cm}}$$

$$\underline{d_{\text{afXen}} \rightarrow < 1 \times 10^{-28} \text{ e} \cdot \text{cm}}$$

Phys.Lett. A376(2012)1347

High voltage break down at 5×10^{-3} torr ?

Paschen's law

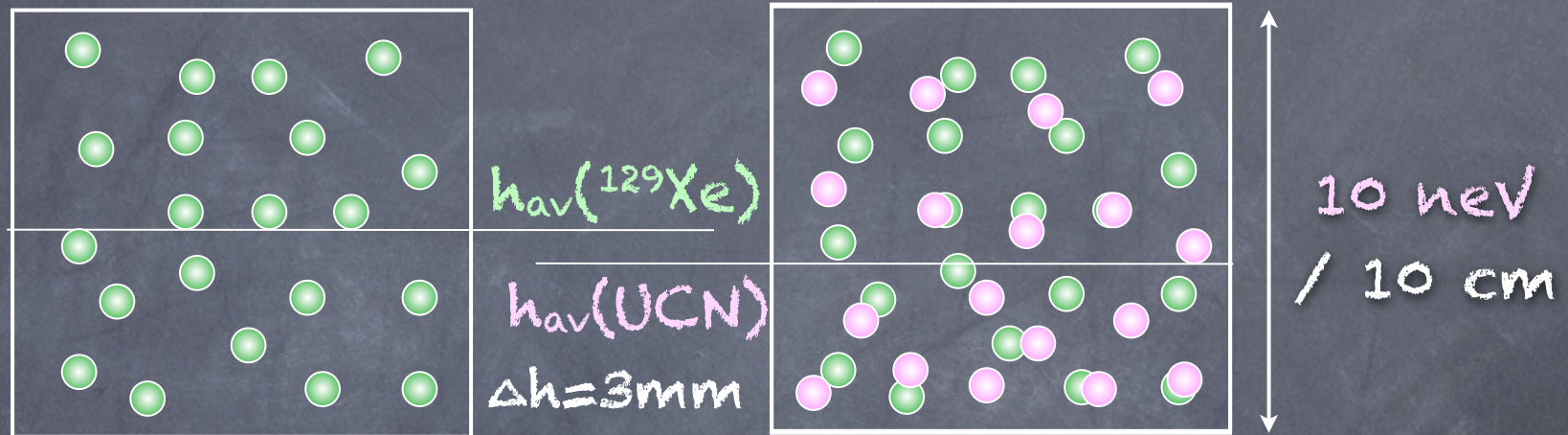


Phys. Lett. A 376(2012)1347

Field gradient control

^{129}Xe : $v_{^{129}\text{Xe}}$ 240 m/s

UCN: v_{UCN} 5 m/s



$$\omega_{^{129}\text{Xe}} = \gamma_{^{129}\text{Xe}} B(h_{^{129}\text{Xe}}), \quad \omega_n = \gamma_n B(h_n)$$

Assuming cylindrical symmetry

$$(\omega_n / \gamma_n) / (\omega_{^{129}\text{Xe}} / \gamma_{^{129}\text{Xe}}) = 1 + \delta\gamma + \delta z (\partial B_{0z} / \partial z) / B_{0z}$$

Precision g factor measurement

Phys.Lett. A376(2012)1347

$$(\omega_n/\gamma_n)/(\omega_{129\text{Xe}}/\gamma_{129\text{Xe}}) = 1 + \delta\gamma + \delta z(\partial B_{0z}/\partial z)/B_{0z}$$

The value of $\delta\gamma$ can be obtained with 10^{-7}
in a 3 cm high cell
at $B_{0z} = 50 \mu\text{T}$, $\partial B_{0z}/\partial z = 5 \text{ nT/m}$
the effect of residual magnetization is small.

The uncertainty of 10^{-7} is much smaller than ^{199}Hg .

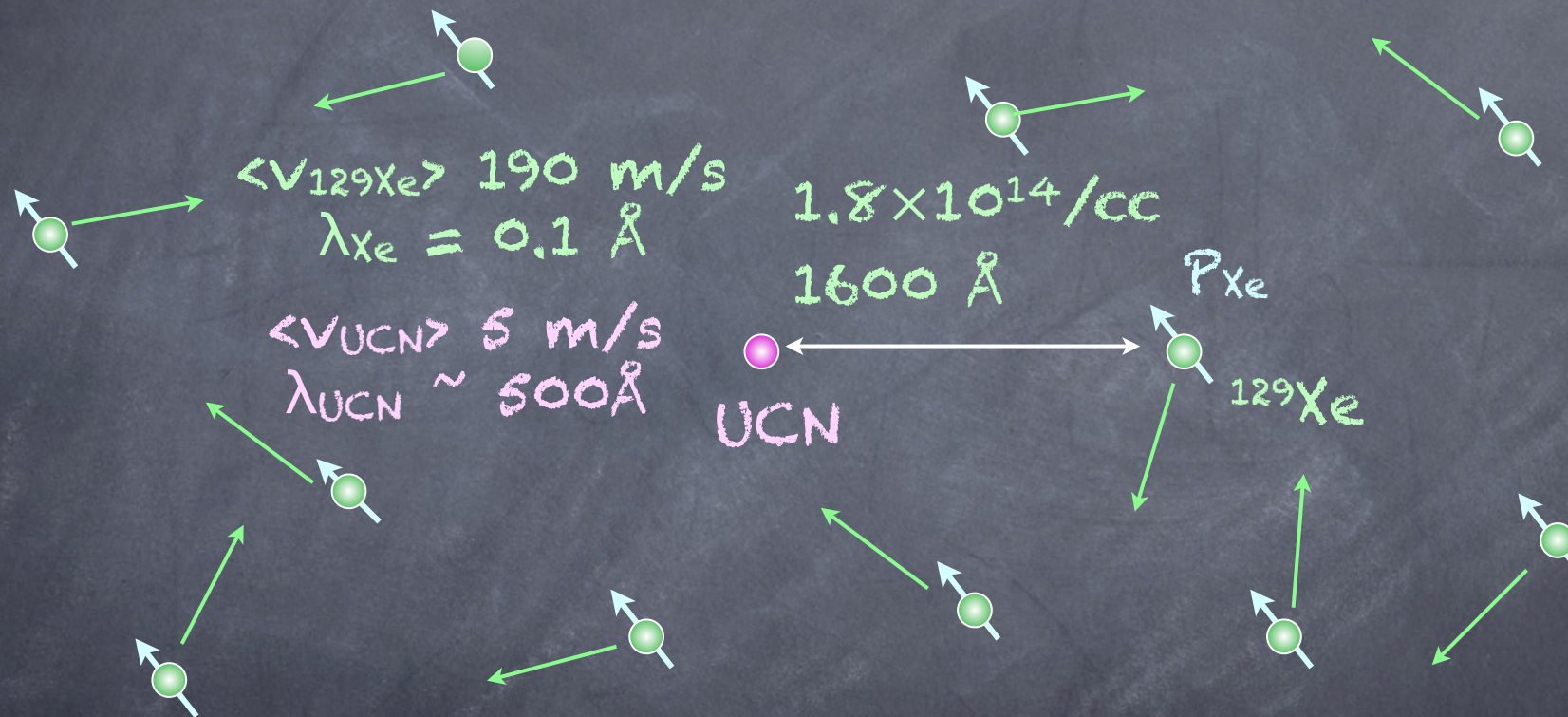
$$\omega_n/\omega_{129\text{Xe}} \text{ in the EDM cell} \rightarrow \partial B_{0z}/\partial z \times 1/10$$

$$d_{\text{afn}} = -\hbar/4 \cdot (\partial B_{0z}/\partial z)/B_{0z}^2 \cdot v_{xy}^2/c^2 \rightarrow < 1 \times 10^{-28} \text{ e} \cdot \text{cm}$$

Effect of ^{129}Xe pseudo magnetism ?

$$0.4 \text{ nT} > B_{\text{or}} = 0.1 \text{ nT}$$

spin dependent coherent scattering



Even if coherence remain,
time variation of B_p , $\omega_{^{129}\text{Xe}}$ and $E \times v/c^2$, ω_r are out of

Rotating field originated from nuclear forth

$$\mathbf{B}_{xy}(t) = \exp\left(-\frac{is_z \omega_r t}{\hbar}\right) \mathbf{B}_v(0) \exp\left(\frac{is_z \omega_r t}{\hbar}\right) + \exp\left(-\frac{is_z \omega'_0 t}{\hbar}\right) \mathbf{B}_p(0) \exp\left(\frac{is_z \omega'_0 t}{\hbar}\right)$$

$$V_I(t) = -\frac{\mu \mathbf{B}_v}{\hbar} \left[\exp\left\{\frac{is_z(\omega_0 - \omega_r)t - \delta}{\hbar}\right\} s_x \exp\left\{\frac{-is_z(\omega_0 - \omega_r)t + \delta}{\hbar}\right\} \right]$$

$$-\frac{\mu \mathbf{B}_p}{\hbar} \left[\exp\left\{\frac{is_z(\omega_0 - \omega'_0)t}{\hbar}\right\} s_x \exp\left\{\frac{-is_z(\omega_0 - \omega'_0)t}{\hbar}\right\} \right]$$

ω_r : $\mathbf{E} \times \mathbf{v}/c^2$ rotation, $\omega'_0 = \omega_{129\text{Xe}}$, $\omega_0 = \omega_n$, δ : phase at $\pi/2$

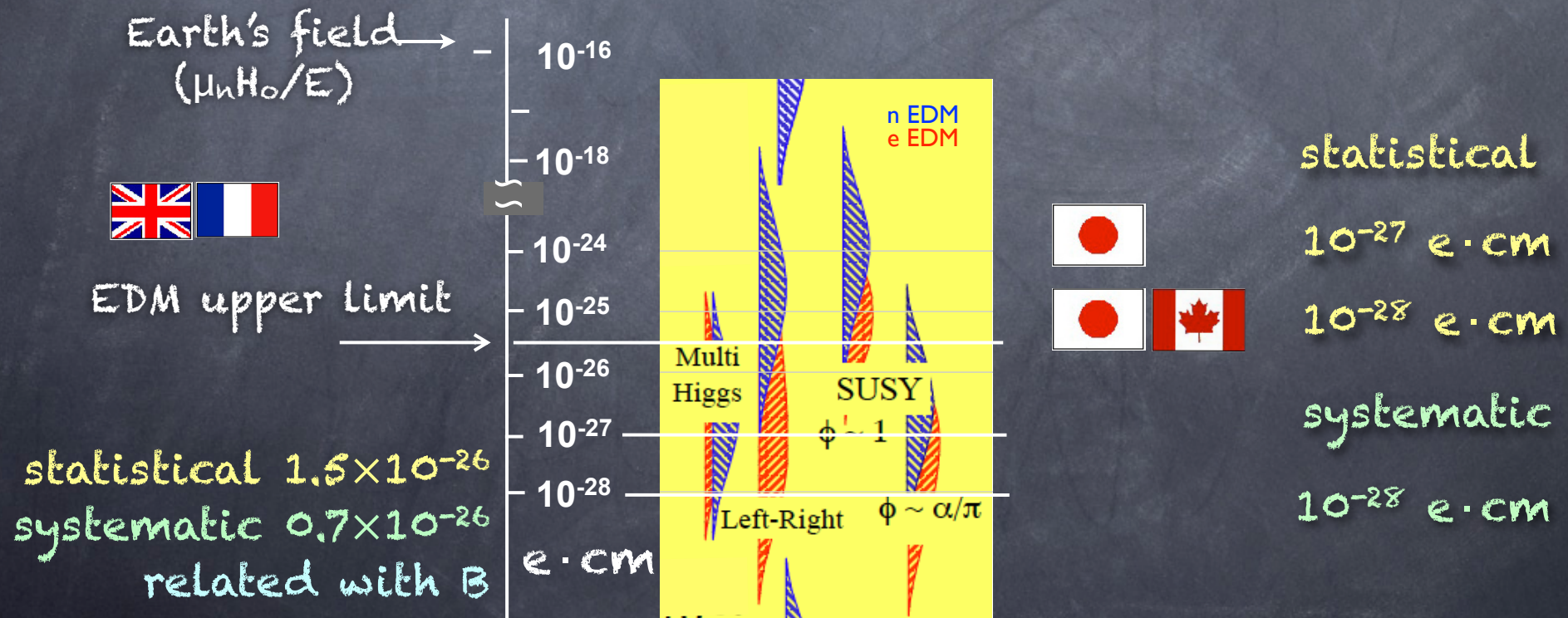
$$\begin{aligned} \delta\omega = & (\gamma_n^2 \mathbf{B}_v \mathbf{B}_p / 2t) [\{1 / (\omega_0 - \omega'_0) + 1 / (\omega_0 - \omega_r)\} \\ & \times \{1 / (\omega'_0 - \omega_r)\} \{\sin((\omega'_0 - \omega_r)t - \delta) + \sin \delta\} \\ & - (\gamma_n^2 \mathbf{B}_v \mathbf{B}_p / 2t) [\{1 / (\omega_0 - \omega'_0)\} \{1 / (\omega_0 - \omega_r)\} \\ & \times \{\sin((\omega_0 - \omega_r)t - \delta) + \sin((\omega_0 - \omega'_0)t + \delta)\}] \end{aligned}$$

These oscillating term $< 10^{-28}$ e·cm

KEK-RCNP EDM measurement

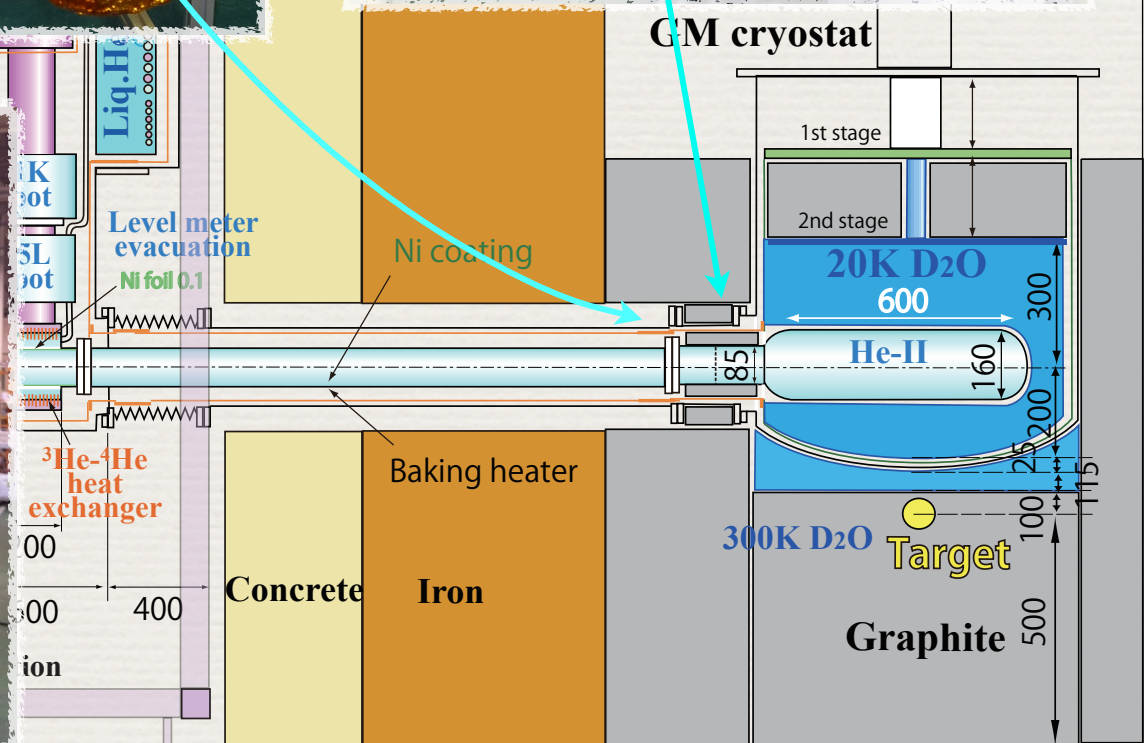
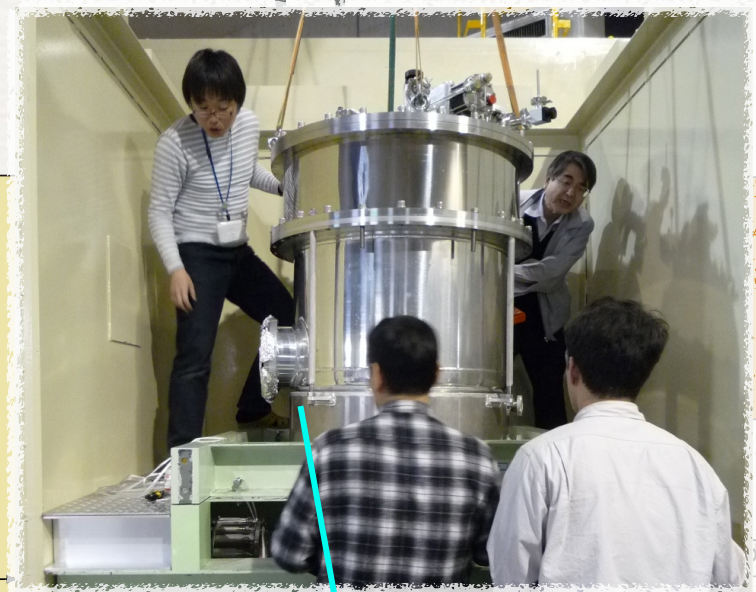
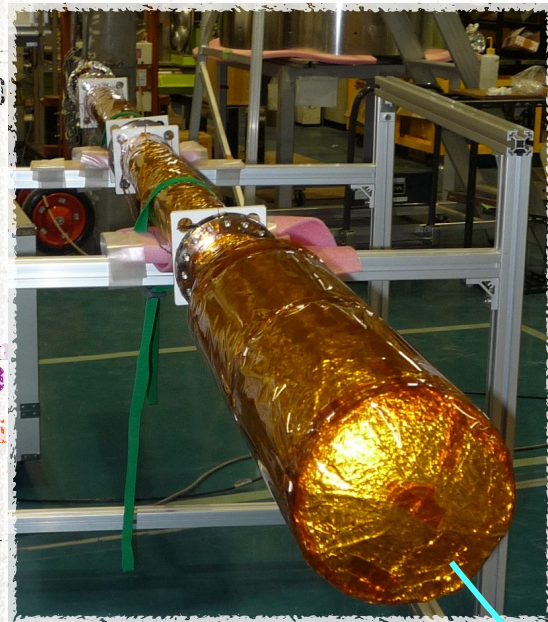
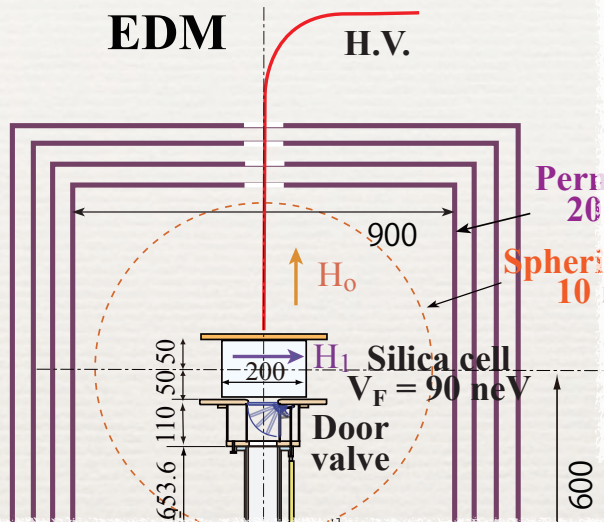
magnetic field	magnetometer	EDM cell	UCN polarization	polarized UCN density
spherical coil Finemet	^{129}Xe comagnetometer	$E_c = 90 \text{ neV}$	90%	600* UCN/cm ³

* 3000 UCN/cm³ at TRIUMF

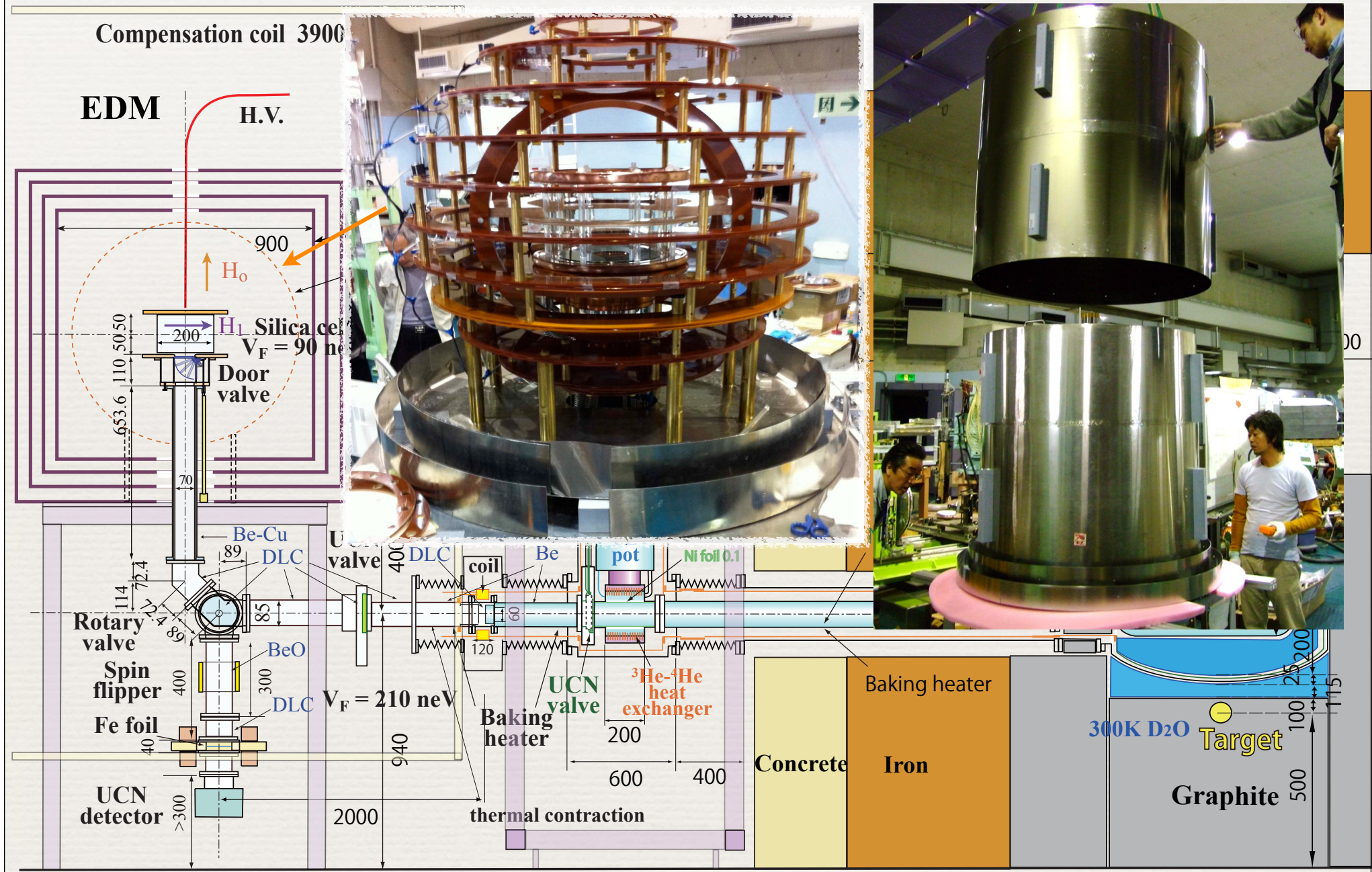


We are constructing the apparatus

Compensation coil 3900×3900



We are constructing the apparatus



Thanks