

Nuclear spin-orbit interaction in ternary fission

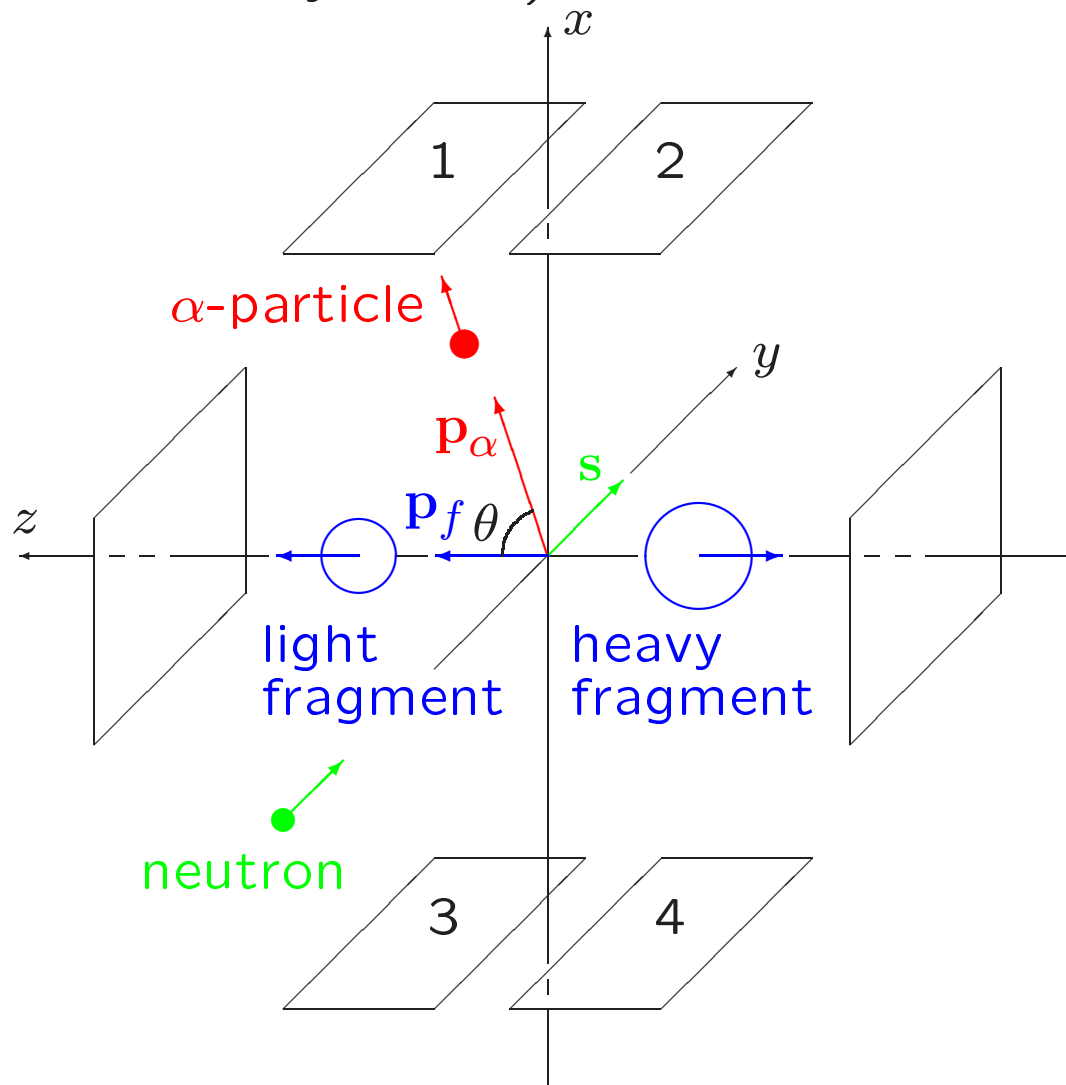
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Plan:

- 1) T-odd angular correlations in ternary fission
- 2) model description of these correlations by spin-orbit interaction
- 3) evidences for target spin effects and their interpretation
- 4) specification for the model with spin-orbit interaction for ternary fission

Asymmetries ($\sim 10^{-3}$) in fission with third particle emission (α -particle, if in ternary fission):



$$D_{13}(\theta) = \frac{N_1 - N_3}{N_1 + N_3}$$

$$D_{24}(\pi - \theta) = \frac{N_2 - N_4}{N_2 + N_4}$$

$$D = \frac{(N_1 + N_2) - (N_3 + N_4)}{(N_1 + N_2) + (N_3 + N_4)} =$$

$$= \frac{D_{13} + \lambda D_{24}}{1 + \lambda} \approx \frac{D_{13} + D_{24}}{2}$$

$$\lambda = \frac{N_2 + N_4}{N_1 + N_3} \approx 1$$

(y,z)-Transverse asym.: T-effect ("TRI-effect"): $D_{13} \simeq D_{24} \simeq D$ ($n+^{233}\text{U}$)

z-Reverse asym.: R-effect ("ROT-effect"): $D_{13} \simeq -D_{24} \Rightarrow D \simeq 0$ ($n+^{235}\text{U}$)

Short list of articles:

T-effect (in ternary fission for α particle as the third particle):

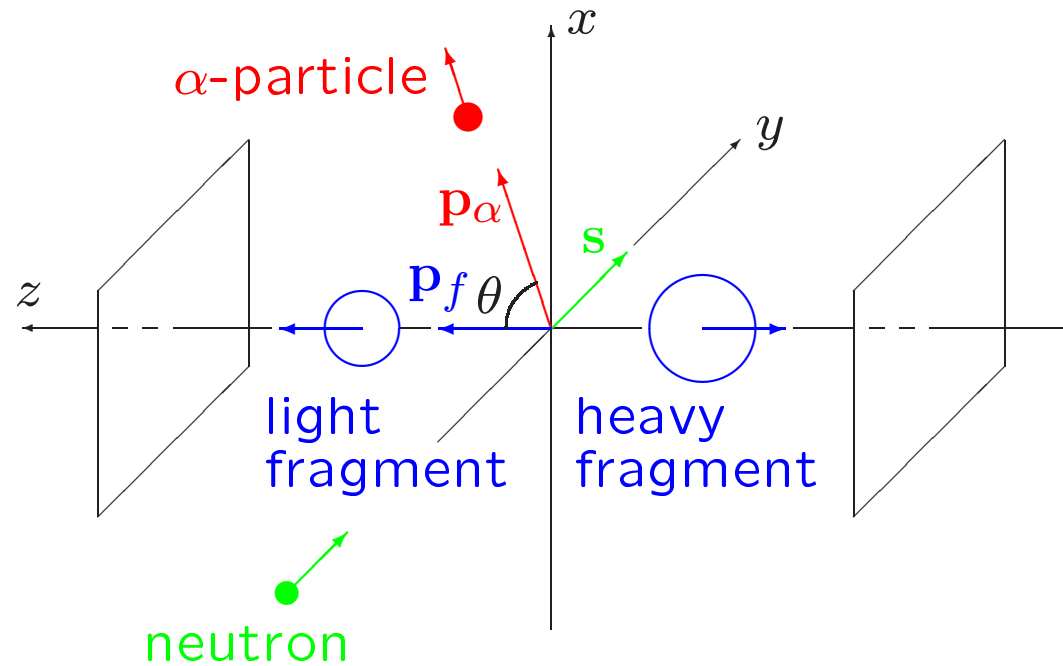
P. Jesinger et al., Nucl. Instr. Meth. Phys. Res. A. **440**, (2000) 618

P. Jesinger et al. Phys. At. Nucl. **65**, (2002) 630

R-effect: (in ternary fission for α particle as the third particle):

F. Goennenwein et al., Phys. Lett. B. **652**, (2007) 13

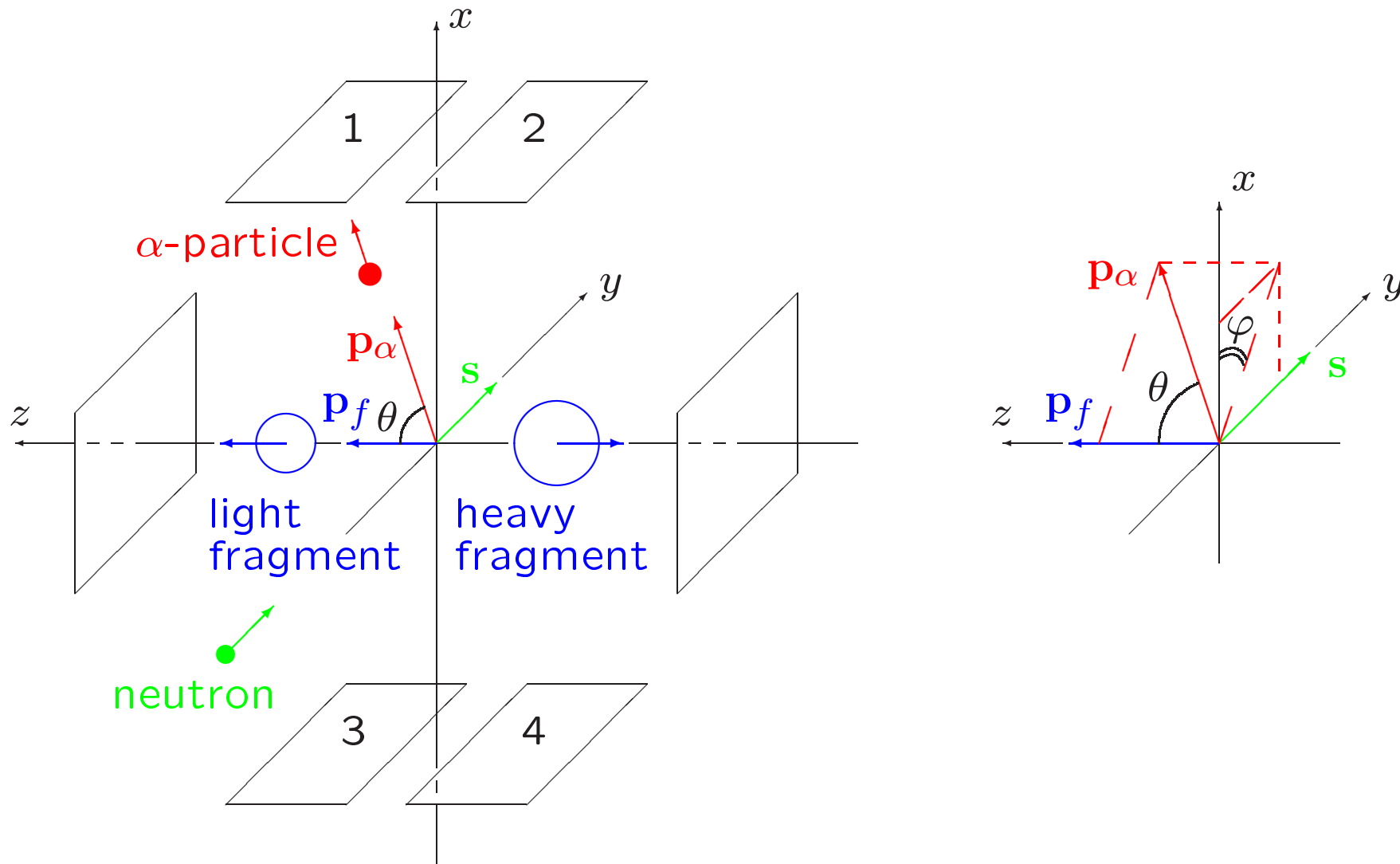
Model with spin-orbit interaction (A.L.Barabanov. Proc. 9-th ISINN, Dubna, 2001, P. 93; arXiv: 0712.3543):



$$\frac{dw(\theta, \varphi, E_\alpha)}{d\Omega dE_\alpha} \sim \sum_{M_f} \left| \int_0^\infty \langle \psi_{\mathbf{p}_\alpha}^{(-)} | \Psi_{JM_f} | \frac{dV(t)}{dt} | \psi_{\alpha i} \Psi_J \rangle e^{i\omega_{fi}t} dt \right|^2,$$

$$V = V_0 + (V_{Jl}(\mathbf{J}\mathbf{l}_\alpha) + (\mathbf{J}\mathbf{l}_\alpha)V_{Jl}), \quad \mathbf{J} = \mathbf{J}_L + \mathbf{J}_H + \mathbf{L} \sim \mathbf{s}$$

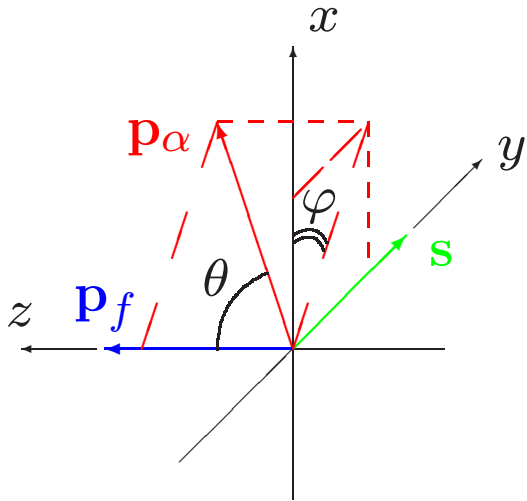
Model with spin-orbit interaction (A.L.Barabanov. Proc. 9-th ISINN, Dubna, 2001, P. 93; arXiv: 0712.3543):



$$\frac{dw(\theta, \varphi, E_\alpha)}{d\Omega dE_\alpha} = W_0(\theta) + p(J) \cos \varphi (W_1(\theta) + W_2(\theta)),$$

T- and R-effects are of the same nature:

$$\frac{dw(\theta, \varphi, E_\alpha)}{d\Omega dE_\alpha} = W_0(\theta) + p(J) \cos \varphi (W_1(\theta) + W_2(\theta)),$$



$$W_0(\theta) = \sum_{Q=0,1,2,\dots} (2Q+1) a_Q(E_\alpha) P_Q(\cos \theta),$$

$$W_1(\theta) = \sum_{Q=1,3,\dots} (2Q+1) b_Q(E_\alpha) P_Q^1(\cos \theta),$$

$$W_2(\theta) = \sum_{Q=2,4,\dots} (2Q+1) b_Q(E_\alpha) P_Q^1(\cos \theta),$$

$$P_Q^1(\cos \theta) = \sin \theta \frac{dP_Q(\cos \theta)}{d \cos \theta}, \quad \text{therefore:} \quad W_1(\theta) = W_1(\pi - \theta)$$

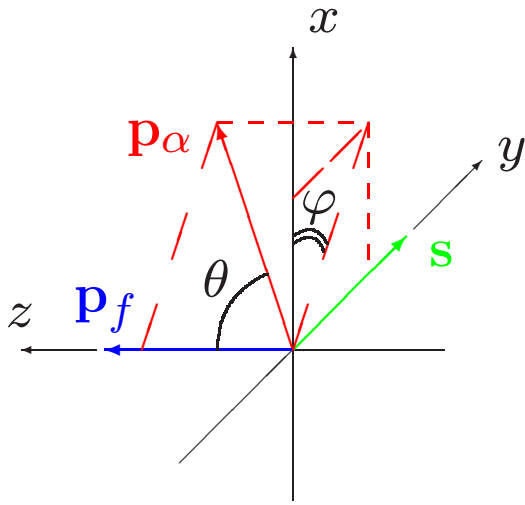
$$W_2(\theta) = -W_2(\pi - \theta)$$

$$\text{if } W_2(\theta) \simeq 0 \quad \Rightarrow \quad D_{13} \simeq D_{24} \simeq D \quad \Rightarrow \quad \text{T-effect (n+}^{233}\text{U)}$$

$$\text{if } W_1(\theta) \simeq 0 \quad \Rightarrow \quad D_{13} \simeq -D_{24} \quad (D \simeq 0!) \quad \Rightarrow \quad \text{R-effect (n+}^{235}\text{U)}$$

3- and 5-fold correlations (T-odd):

$$\frac{dw(\theta, \varphi, E_\alpha)}{d\Omega dE_\alpha} = W_0(\theta) + p(J) \cos \varphi (W_1(\theta) + W_2(\theta)),$$



$$W_1(\theta) = \sum_{Q=1,3,\dots} (2Q+1) b_Q(E_\alpha) P_Q^1(\cos \theta),$$

$$W_2(\theta) = \sum_{Q=2,4,\dots} (2Q+1) b_Q(E_\alpha) P_Q^1(\cos \theta),$$

$$\cos \varphi W_1(\theta) = 3 b_1(E_\alpha) \underbrace{\cos \varphi \sin \theta}_{\text{|||}} + \dots \Rightarrow \text{T-effect (n+}^{233}\text{U)}$$

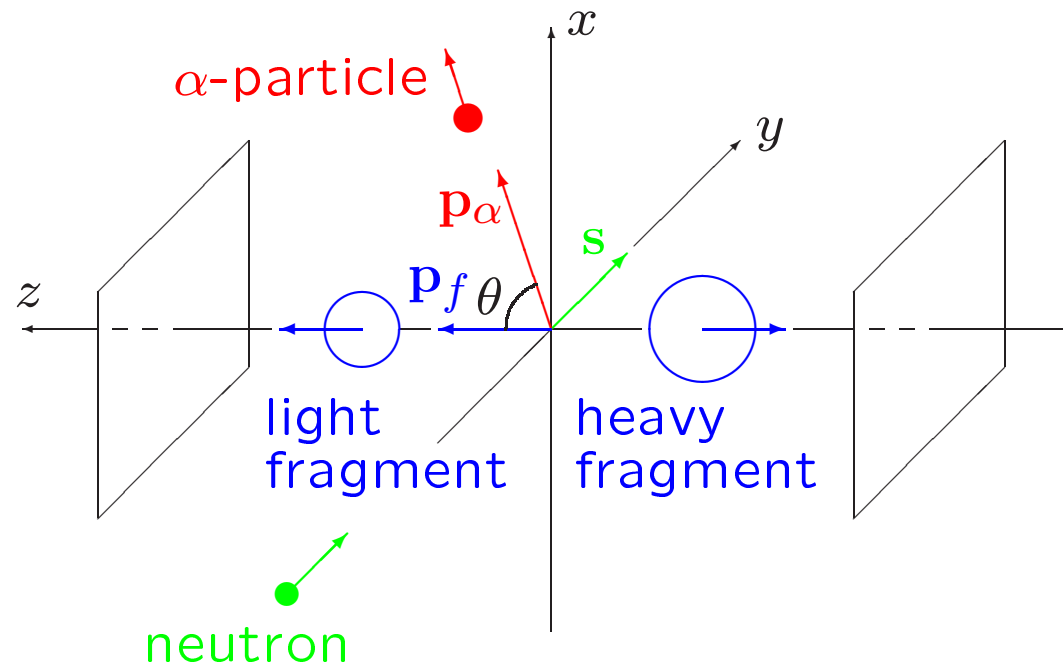
$$\underbrace{(\mathbf{p}_\alpha [\mathbf{s} \times \mathbf{p}_f])}_{\text{|||}}$$

$$\cos \varphi W_2(\theta) = 15 b_2(E_\alpha) \underbrace{\cos \varphi \sin \theta \cos \theta}_{\text{|||}} + \dots \Rightarrow \text{R-effect (n+}^{235}\text{U)}$$

$$\underbrace{(\mathbf{p}_\alpha [\mathbf{s} \times \mathbf{p}_f]) (\mathbf{p}_\alpha \mathbf{p}_f)}_{\text{|||}}$$

Are there some evidences for spin (target spin!) – orbit interaction?

What is a nature of such interaction?



$$\frac{dw(\theta, \varphi, E_\alpha)}{d\Omega dE_\alpha} \sim \sum_{M_f} \left| \int_0^\infty \langle \psi_{\mathbf{p}_\alpha}^{(-)} \psi_{JM_f} | \frac{dV(t)}{dt} | \psi_{\alpha i} \psi_J \rangle e^{i\omega_{f_i} t} dt \right|^2,$$

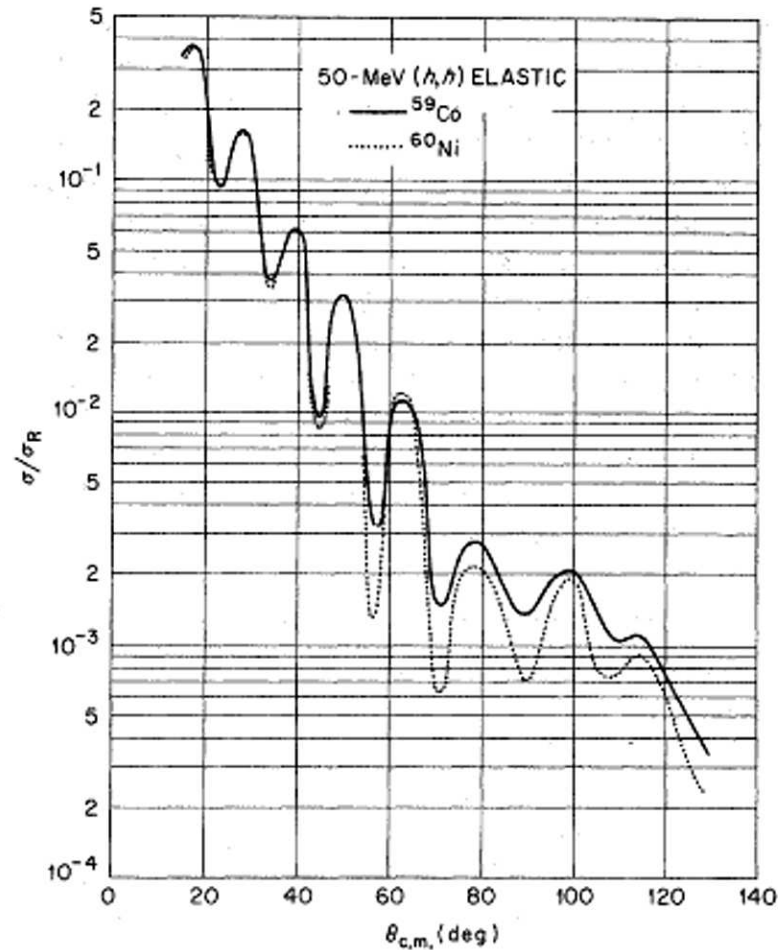
$$V = V_0 + (V_{Jl}(\mathbf{J}\mathbf{l}_\alpha) + (\mathbf{J}\mathbf{l}_\alpha)V_{Jl}), \quad \mathbf{J} = \mathbf{J}_L + \mathbf{J}_H + \mathbf{L} \sim \mathbf{s}$$

Spin-Orbit and Target Spin Effects in Helion Elastic Scattering*

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$$I = \frac{7}{2} \text{ for } {}^{59}\text{Co}$$

$$I = 0 \text{ for } {}^{60}\text{Ni}$$

FIG. 5. Angular distributions for 50-MeV helion elastic scattering from ${}^{59}\text{Co}$ and ${}^{60}\text{Ni}$. The plot for each target was obtained by drawing a smooth curve through the experimental data.

TARGET-SPIN EFFECTS ON ELASTIC SCATTERING CROSS SECTIONS*

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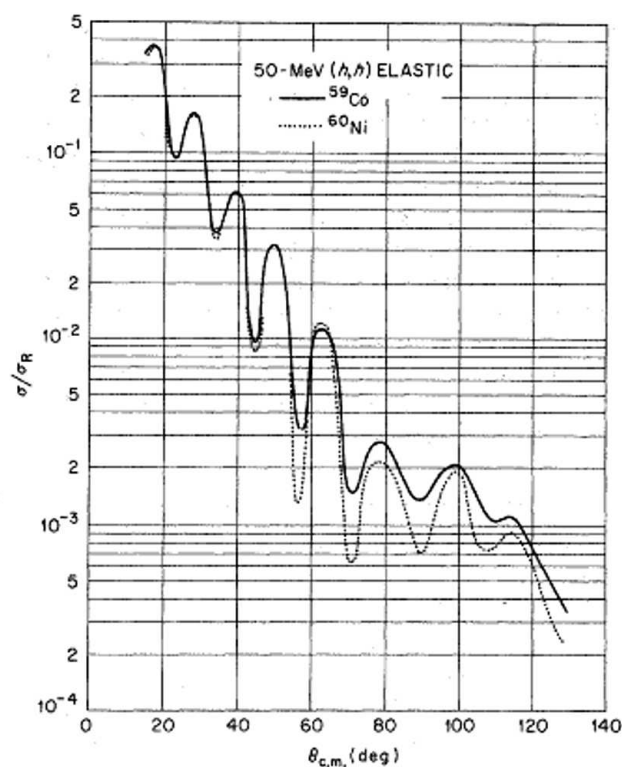
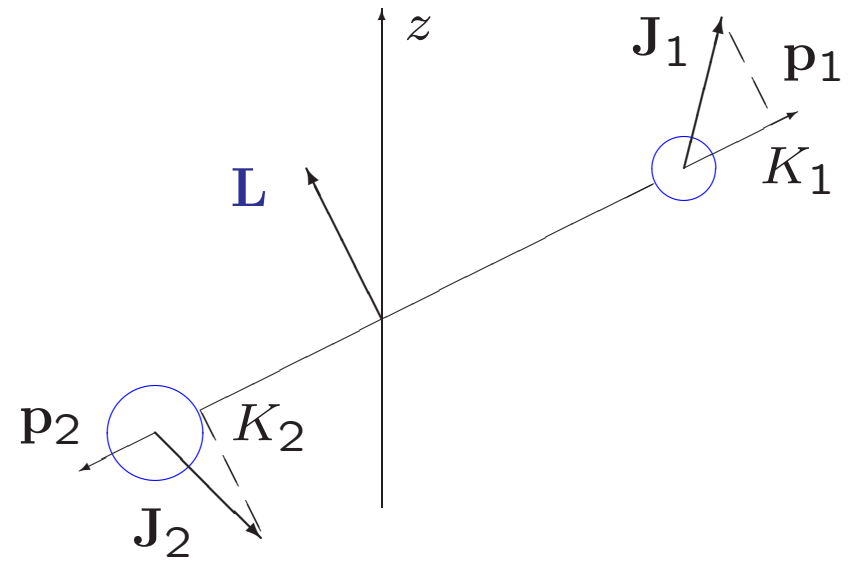
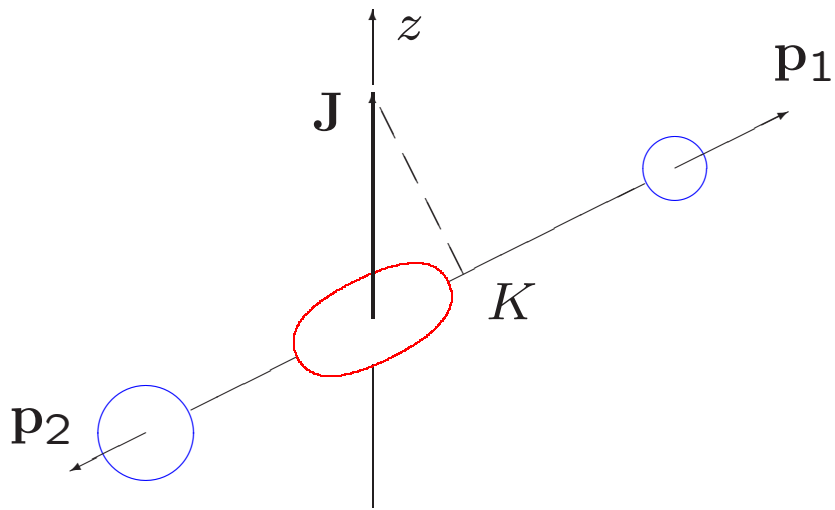


FIG. 5. Angular distributions for 50-MeV helion elastic scattering from ^{59}Co and ^{60}Ni . The plot for each target was obtained by drawing a smooth curve through the experimental data.

$$|\text{odd}, IM\rangle = \alpha |(\text{even}, 0), j = I; M\rangle$$

$$+ \sum_{L,j} \beta_{Lj} |(\text{even}, L), j; IM\rangle. \quad (1)$$

$$\sigma_{el}(\text{odd}) \approx \sigma_{el}(\text{even}) + \sum_L \frac{(2\alpha\beta_{LI})^2}{2L+1} \sigma_{inel}(\text{even}, 0^+ \rightarrow L^+) \quad (2)$$



A. Bohr, in *Proc. Int. Conf. on the Peaceful Uses of Atomic Energy, Geneva, 1955*

V.M. Strutinsky, *ZhETF* **30**, 606 (1956)

$$\begin{aligned} \mathbf{J} &\rightarrow \mathbf{J}_1 + \mathbf{J}_2 + \mathbf{L} \\ \mathbf{K} &\rightarrow \mathbf{K}_1 + \mathbf{K}_2 \end{aligned}$$

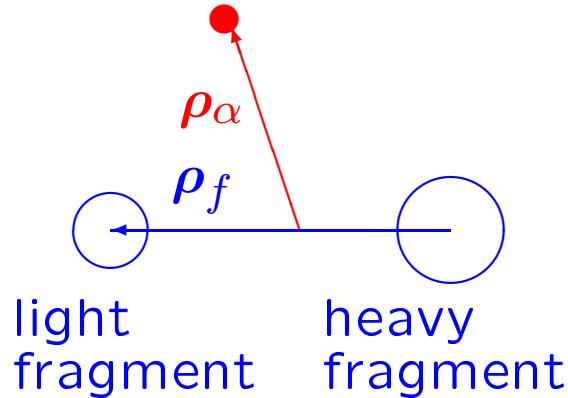
$$V = V_0 + (V_{Jl}(\mathbf{J}\mathbf{l}_\alpha) + (\mathbf{J}\mathbf{l}_\alpha)V_{Jl}), \quad \mathbf{J}\mathbf{l}_\alpha = (\mathbf{J}_1 + \mathbf{J}_2 + \mathbf{L})\mathbf{l}_\alpha \rightarrow \mathbf{L}\mathbf{l}_\alpha$$

Model with hyperspherical harmonics for three particle in the output channel:

Usual description of reactions (input wave + output waves):

$$\Psi \rightarrow \sum_{\lambda_\alpha} a(\lambda_\alpha) \left(\frac{u_{\alpha l}^{(-)}}{r_\alpha} \varphi_{\lambda_\alpha}^\alpha + \sum_{\beta} \left(\frac{m_\beta k_\beta}{m_\alpha k_\alpha} \right)^{1/2} \sum_{\lambda_\beta} S_J(\lambda_\alpha \rightarrow \lambda_\beta) \frac{u_{\beta l_\beta}^{(+)}}{r_\beta} \varphi_{\lambda_\beta}^\beta + \sum_{\gamma} \frac{k_\gamma^2 (\pi/2)^{1/2}}{(m_\alpha k_\alpha)^{1/2}} \sum_{\lambda_\gamma} S_J(\lambda_\alpha \rightarrow \lambda_\gamma) \frac{H_{N+2}^{(+)}(k_\gamma \rho_\gamma)}{(k_\gamma \rho_\gamma)^2} |JM; \lambda_\gamma\rangle \right).$$

α -particle



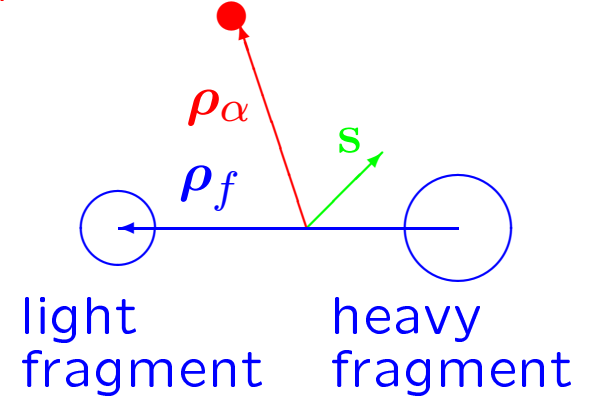
$$\rho_f = \rho \sin \vartheta, \quad \rho_\alpha = \rho \cos \vartheta,$$

$$|JM; \lambda_\gamma\rangle \sim (\sin \vartheta)^{l_f} (\cos \vartheta)^{l_\alpha} P_n^{L_f + \frac{1}{2}, l_\alpha + \frac{1}{2}}(\cos 2\vartheta) \times \sum C_{JM l_\alpha m_\alpha}^{J_0 M_0} C_{L_f m_f F f}^{JM} Y_{l_\alpha m_\alpha}(\rho_\alpha) Y_{L_f m_f}(\rho_f) |F f\rangle$$

$$N = 2n + l_\alpha + l_f$$

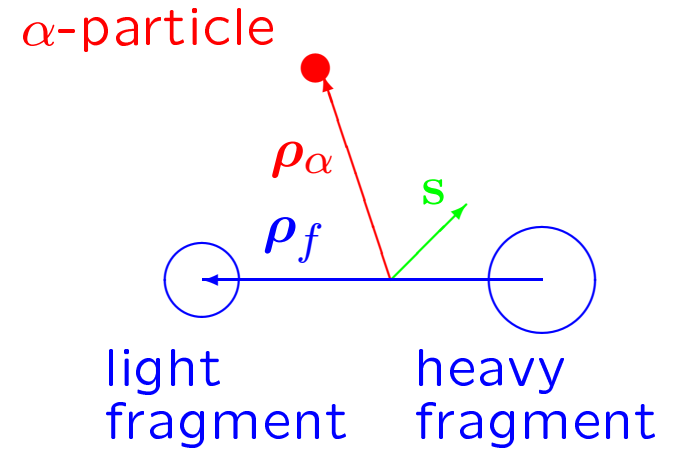
Differential cross section
for ternary fission:

α -particle



$$\begin{aligned} \frac{d\sigma_f}{d\Omega_f d\Omega_\alpha} = & \frac{\pi\lambda^2}{4\pi} \sum_{J_0 J'_0} g_{J'_0} \sum_{Q=0,1} U(IsJ'_0Q, J_0s) \tau'_Q(s) \times \\ & \sum_{L_f L'_f l_\alpha l'_\alpha n n' J J' F} S_{J_0}(0\frac{1}{2} \rightarrow L_f l_\alpha N J F) S_{J'_0}^*(0\frac{1}{2} \rightarrow L'_f l'_\alpha N' J' F) \times \\ & \sum_{\Lambda_f \Lambda_\alpha} \hat{\Lambda}_f \hat{\Lambda}_\alpha \hat{l}'_\alpha \hat{J}' \hat{J} C_{L'_f 0 \Lambda_f 0}^{L_f 0} C_{l'_\alpha 0 \Lambda_\alpha 0}^{l_\alpha 0} U(F J' L_f \Lambda_f, L'_f J) \begin{Bmatrix} J_0 & J & l_\alpha \\ J'_0 & J' & l'_\alpha \\ Q & \Lambda_f & \Lambda_\alpha \end{Bmatrix} \times \\ & \left(\int_0^{\pi/2} (\sin \vartheta)^{L_f + L'_f + 2} (\cos \vartheta)^{l_\alpha + l'_\alpha + 2} P_n^{L_f + \frac{1}{2}, l_\alpha + \frac{1}{2}} (\cos 2\vartheta) P_{n'}^{L'_f + \frac{1}{2}, l'_\alpha + \frac{1}{2}} (\cos 2\vartheta) d\vartheta \right) \\ & (4\pi)^{1/2} \sum_{q \lambda_f \lambda_\alpha} C_{\Lambda_f \lambda_f \Lambda_\alpha \lambda_\alpha}^{Qq} Y_{Qq}^*(s) Y_{\Lambda_f \lambda_f}(\rho_f) Y_{\Lambda_\alpha \lambda_\alpha}(\rho_\alpha) \end{aligned}$$

Differential cross section
for ternary fission:



$$\frac{d\sigma_f}{d\Omega_f d\Omega_\alpha} = \dots \underbrace{(4\pi)^{3/2} \sum_{q\lambda_f\lambda_\alpha} C_{\Lambda_f\lambda_f\Lambda_\alpha\lambda_\alpha}^{Qq} Y_{Qq}^*(\mathbf{s}) Y_{\Lambda_f\lambda_f}(\rho_f) Y_{\Lambda_\alpha\lambda_\alpha}(\rho_\alpha)}_{\phi_{\Lambda_f\Lambda_\alpha}^Q(\mathbf{s}, \rho_f, \rho_\alpha)}$$

$$\phi_{11}^1(\mathbf{s}, \rho_f, \rho_\alpha) \sim (\rho_\alpha [\mathbf{s} \times \rho_f]) \Rightarrow \text{T-effect (n+}^{233}\text{U)}$$

$$\phi_{22}^1(\mathbf{s}, \rho_f, \rho_\alpha) \sim (\rho_\alpha [\mathbf{s} \times \rho_f]) (\rho_\alpha \rho_f) \Rightarrow \text{R-effect (n+}^{235}\text{U)}$$

Summary

1. Both T- and R-effects (in particular, 3- and 5-fold correlations) in ternary fission, apparently, are due to spin-orbit interaction that mixes exit states (final state interaction).
2. Interaction of α -particle with collective degrees of freedom of the fissioning nucleus (interaction of \mathbf{l}_α and \mathbf{L}_f), apparently, gives dominant contribution to the spin-orbit interaction.