

Observation of shape isomeric states in fission fragments

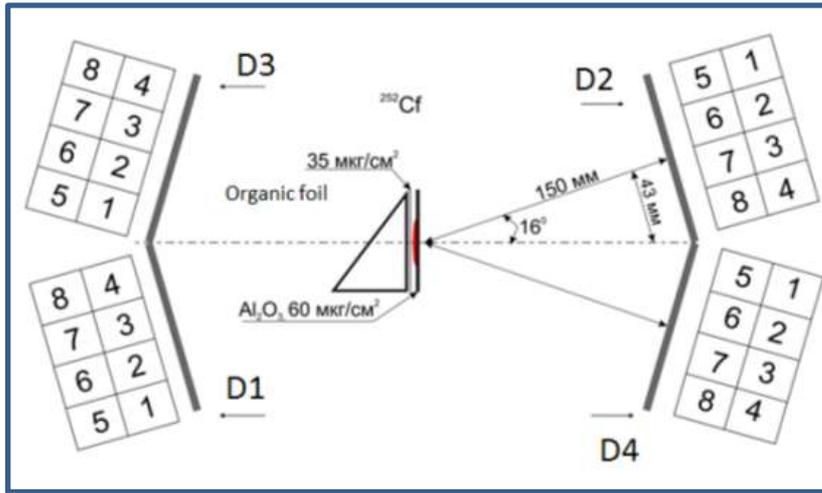
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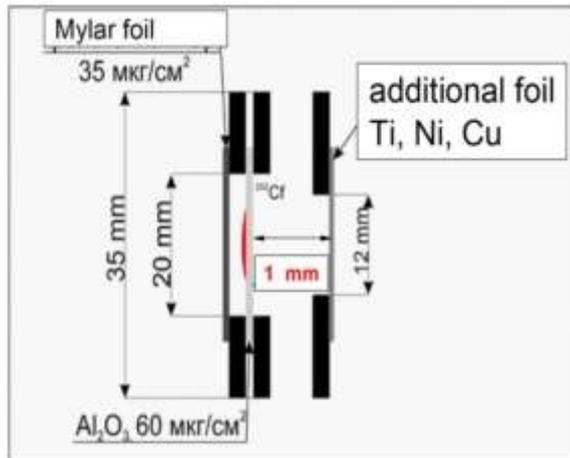
³University of Stellenbosch, Faculty of Military Science, Military Academy, Saldanha 7395, South Africa

Just to remind: our previous experimental results

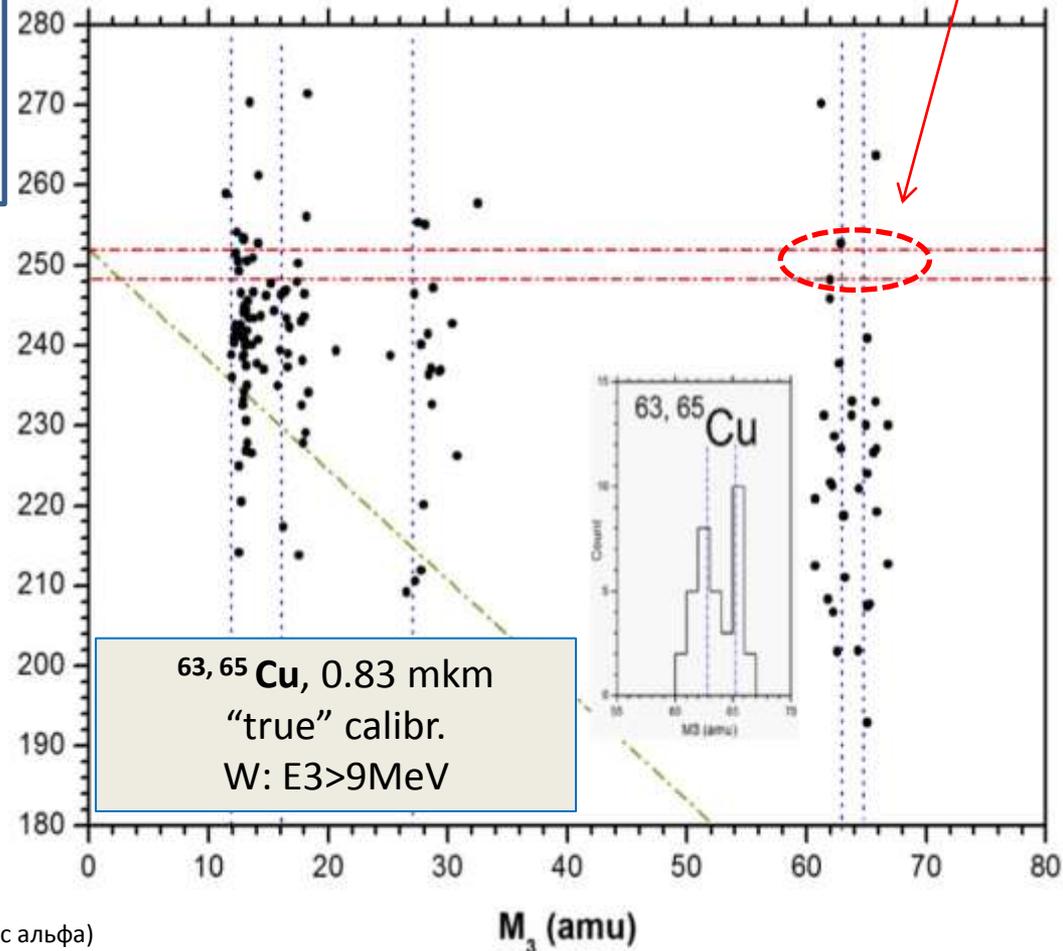


In each ternary event:
 $M_1 > M_2 > M_3$

region of elastic
 FF scattering



Layout of the source module



$^{63,65}\text{Cu}$, 0.83 mkm
 "true" calibr.
 W: $E_3 > 9\text{MeV}$

Rutherford scattering - L_FF/Ti

2 solutions are possible: the same scattering angle at two different impact parameters

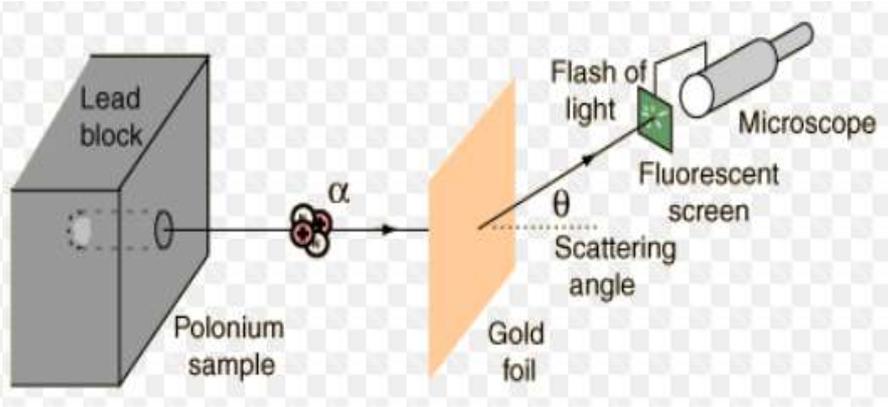
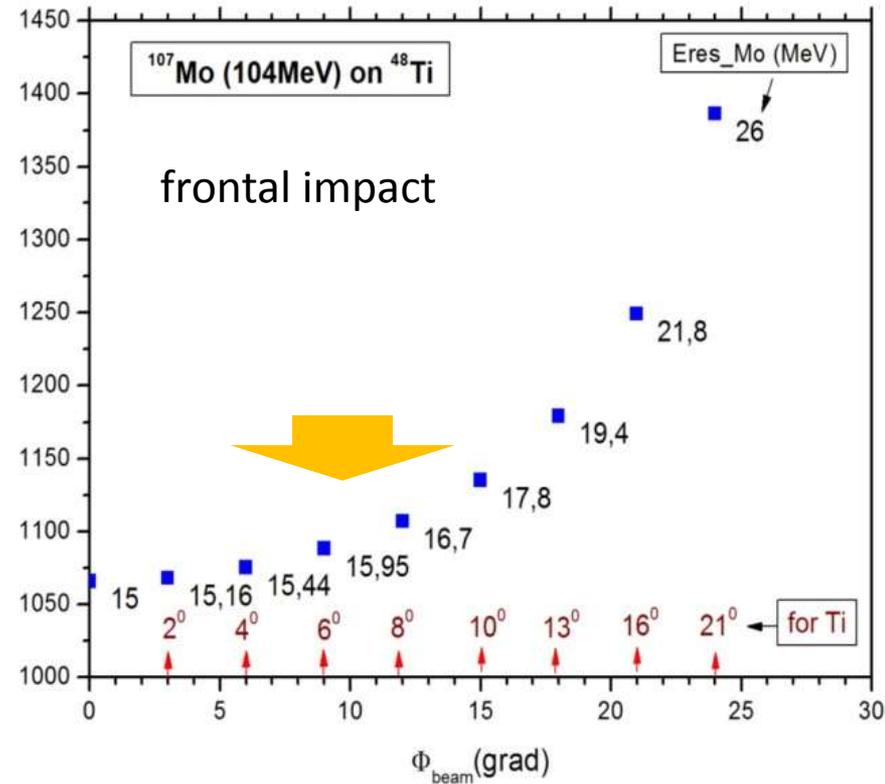
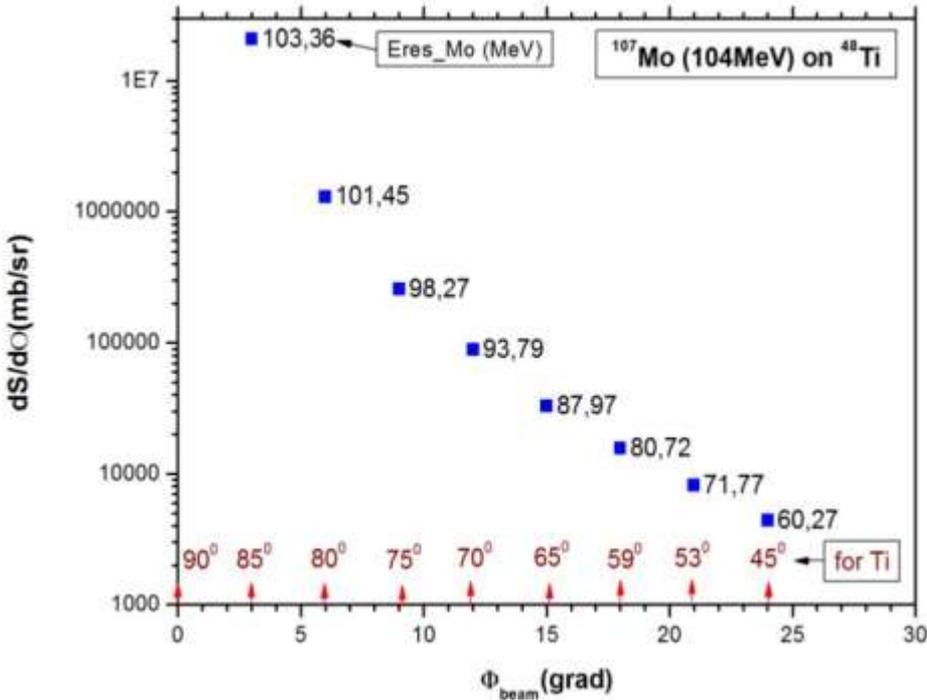
$$\sigma(\theta) = \left(\frac{1}{4\pi\epsilon_0}\right)^2 \frac{Z^2 e^4}{M^2 v^4} \times \frac{1}{\sin^4(\theta/2)}$$

Ze = the positive charge of the target atom,

M = the mass of the α particle,

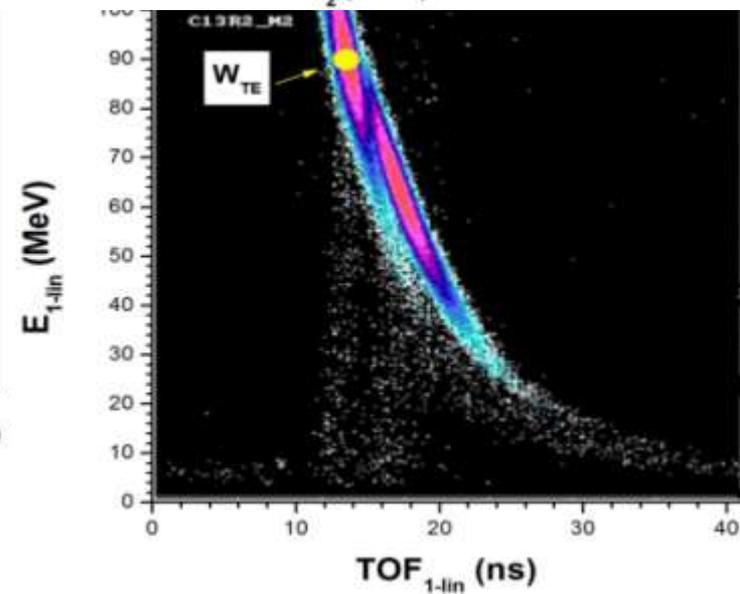
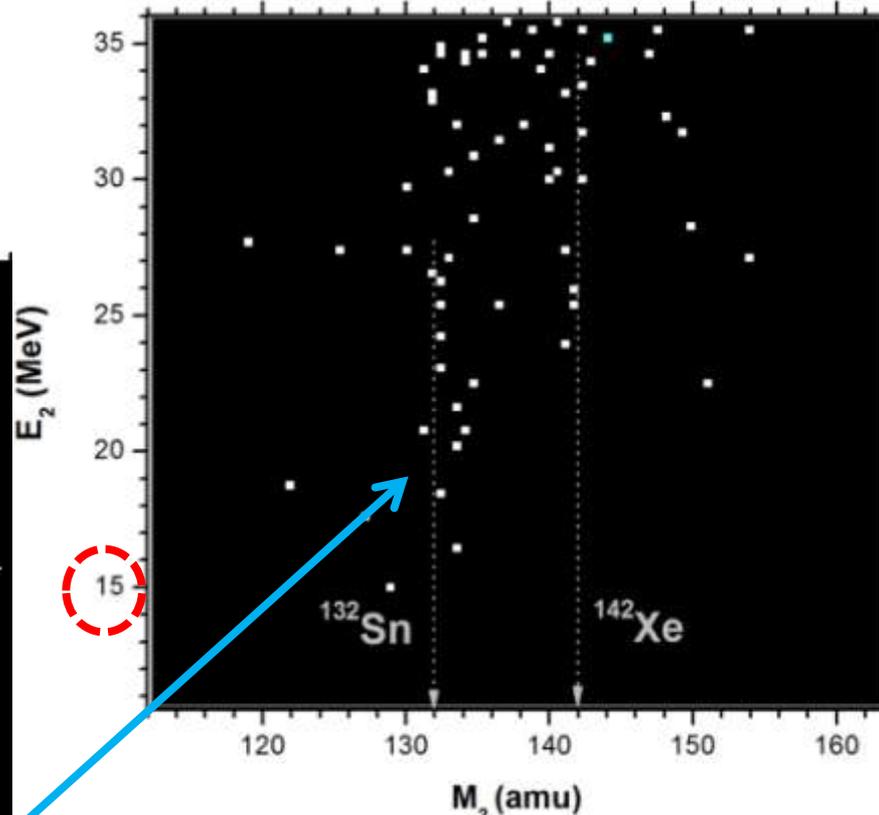
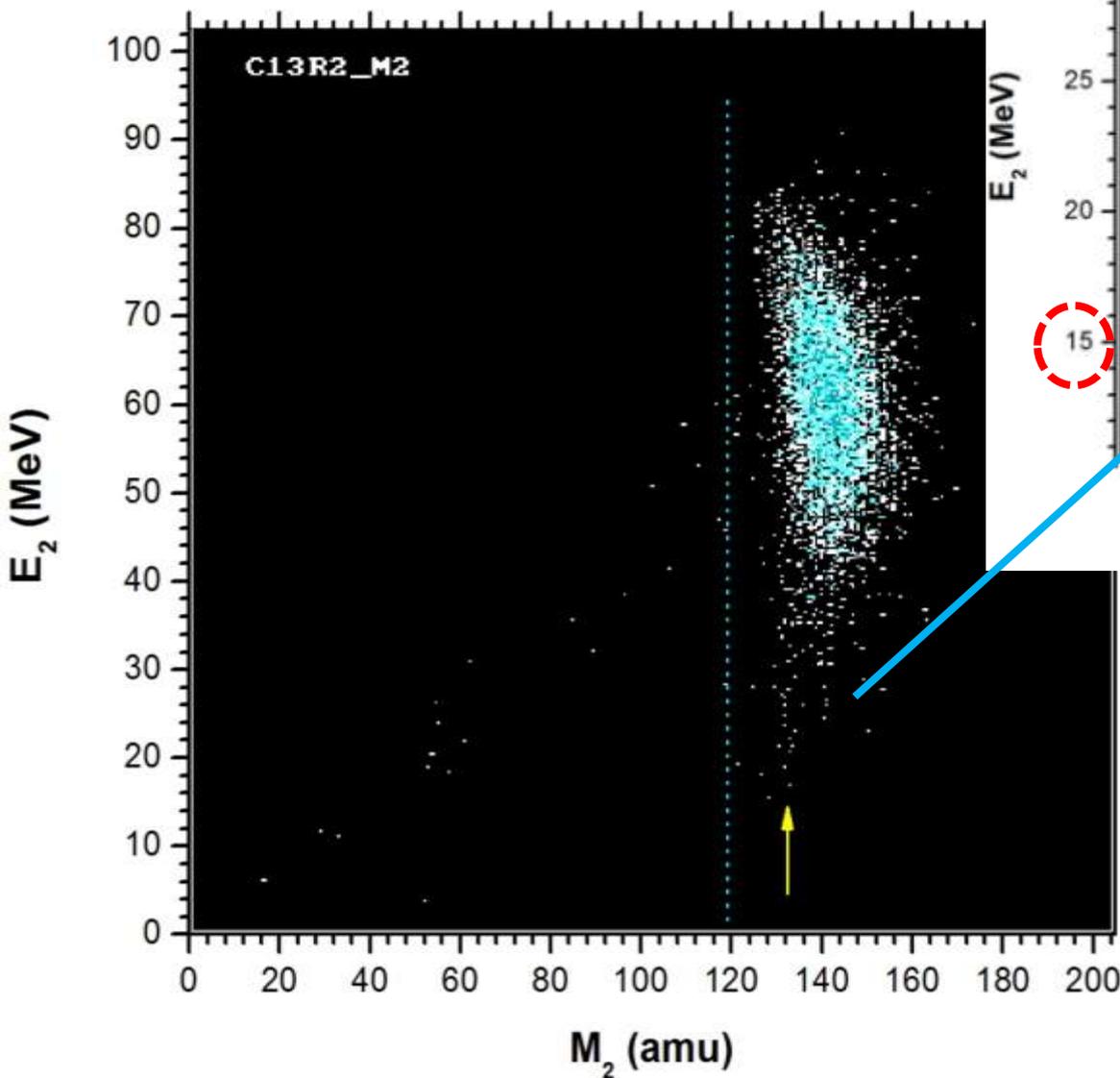
v = incident speed of the α particle,

θ = scattering angle,



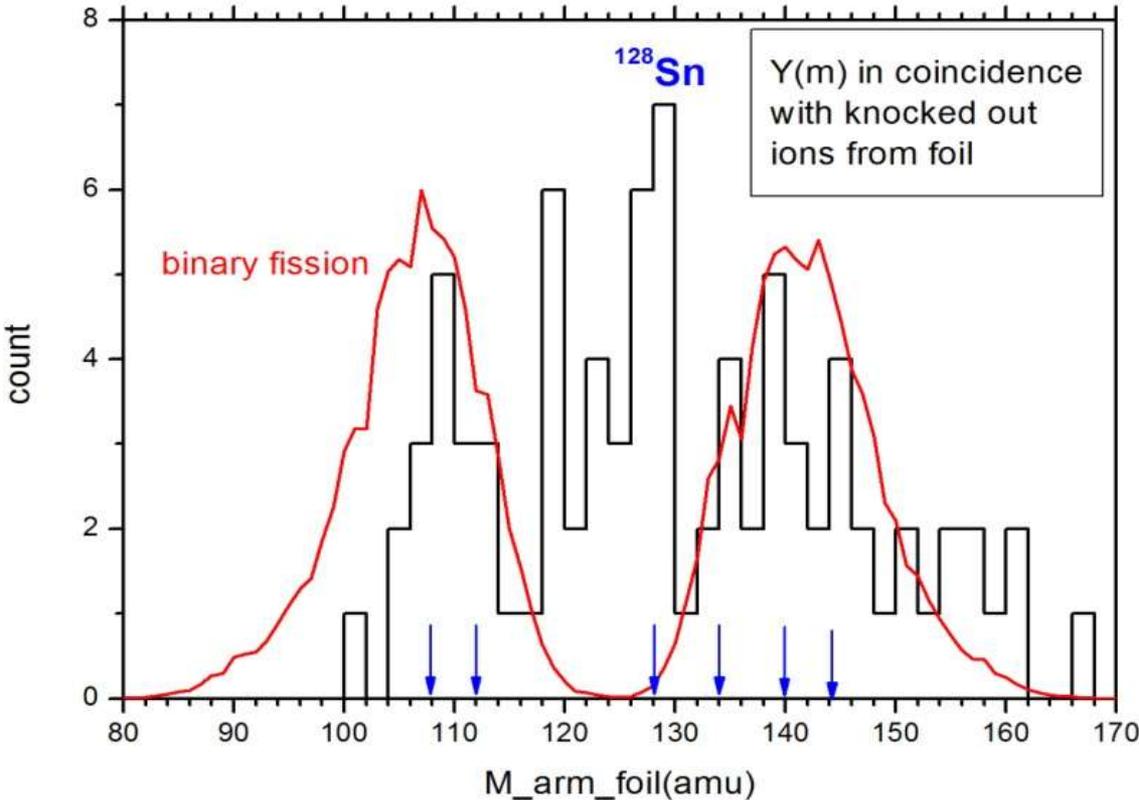
Original Rutherford experiment

Low energy heavy FFs – good mass reconstruction



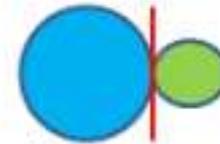
Mass spectrum of the fragments detected in coincidence with the knocked out ions of Ti, Ni, Cu

Really:



Hypothesis:

we suppose each FF, at least just after scission, looks like a di-nuclear system "magic cluster + light ion"



Due to the brake-up in the foil both constituents become free. Thus one of them should be magic nucleus.

Conclusions presented at the previous ISINN21 meeting in 2013.

Presumably:

1. Inelastic impact, at least the frontal one, makes free the constituents of the di-nuclear system (fission fragment) formed in the binary fission .
2. Bearing in mind the distance between the Cf source and the generating foil (~1mm) the lower limit of the life-time of this di-nuclear system (shape isomer) is about 0.1ns.
3. Relative probability of elastic Rutherford scattering of fission fragments i.e. taking place without missing mass is much less than those in the inelastic channel. In other words, **the bulk of the fragments from the conventional binary fission are born as shape-isomers .**

Principal question to be answered :

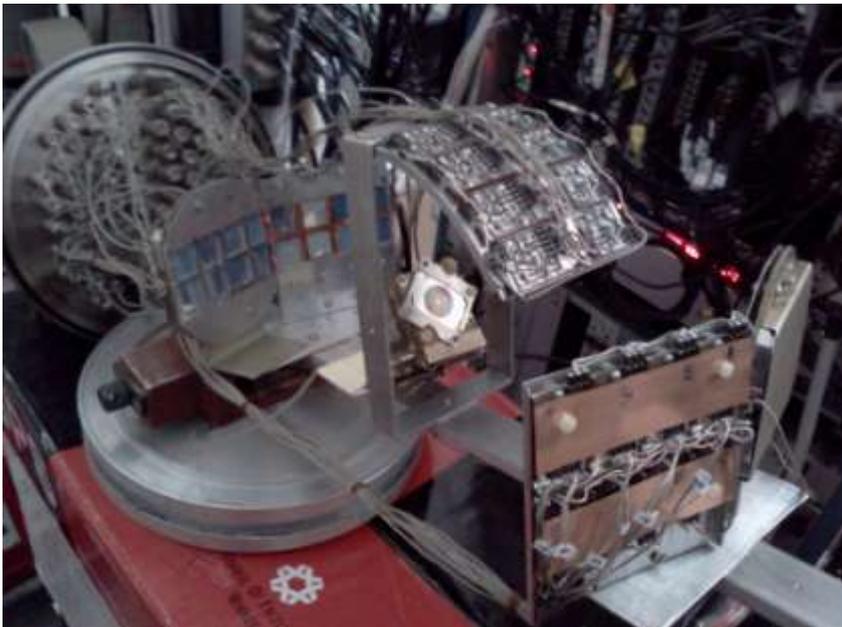
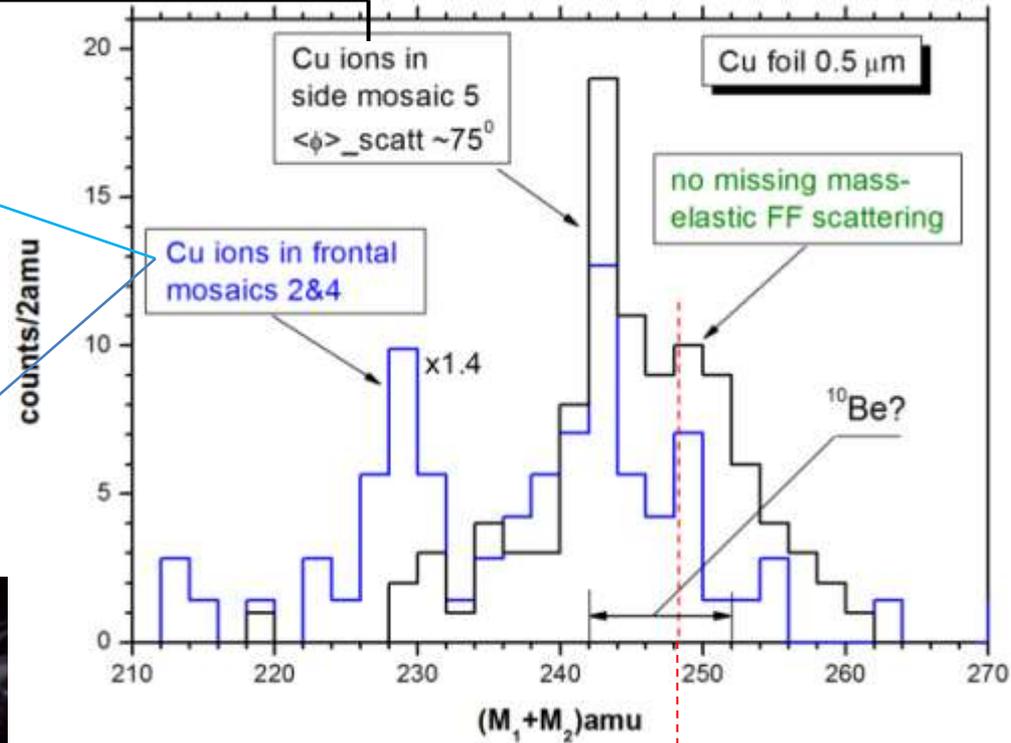
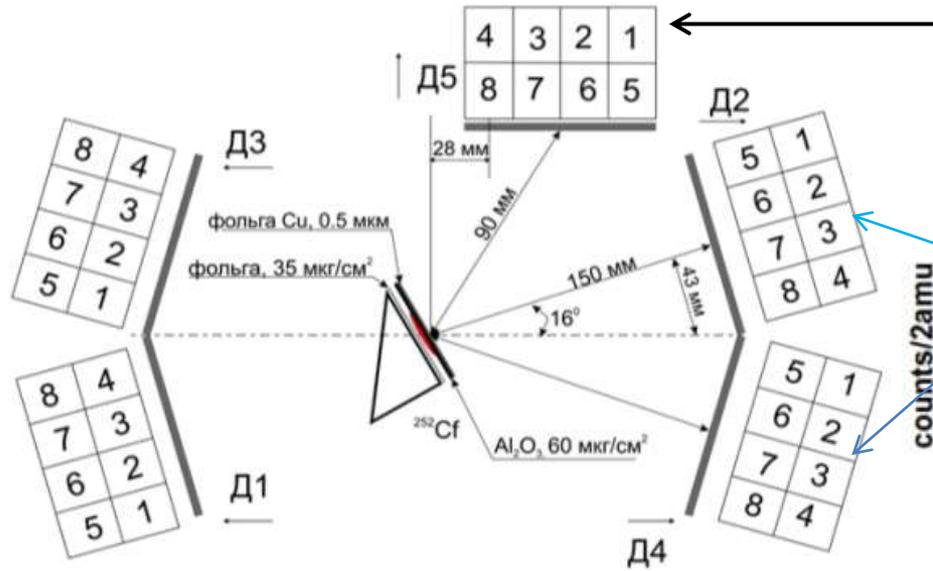
Are there contradictions of the experimental results obtained and model proposed with well known features of the fission process?

1. We suppose that the bulk of the fragments of conventional binary fission look like di-nuclear systems (are in shape isomeric states) just after scission.

Does it mean that the total yield of the effect is sufficient to change the macroscopical fission constants?

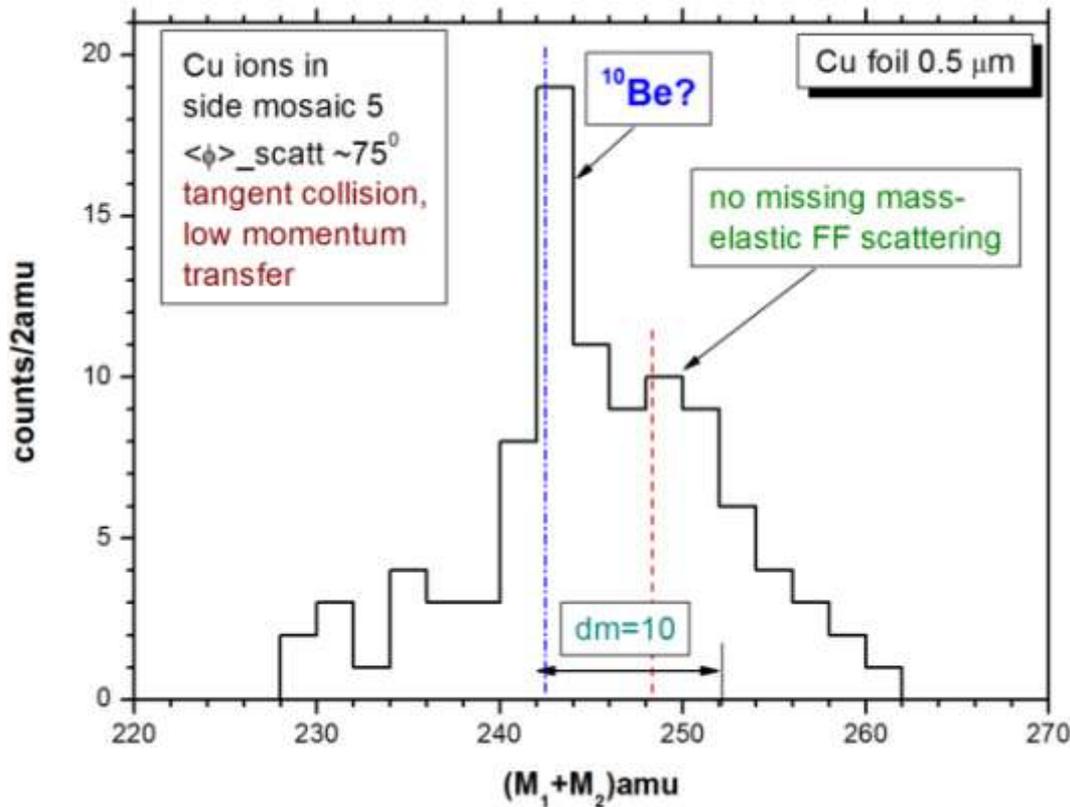
Our answer: **No!** The brake-up of di-nuclear system (FF) appears to occur in almost frontal inelastic collisions which takes place with relatively low probability.

Angular dependence of the brake-up yield



At the large angles of scattering of the knocked-out ions from the foil conventional elastic Rutherford scattering takes place.

Manifestation of ^{10}Be based di-nuclear system



J. Phys. G: Nucl. Part. Phys. **26** (2000) L97–L102
Nuclear quasi-molecular states in ternary fission

D N Poenaru†‡§, B Dobrescu†, W Greiner‡, J H Hamilton§ and A V Ramayya§



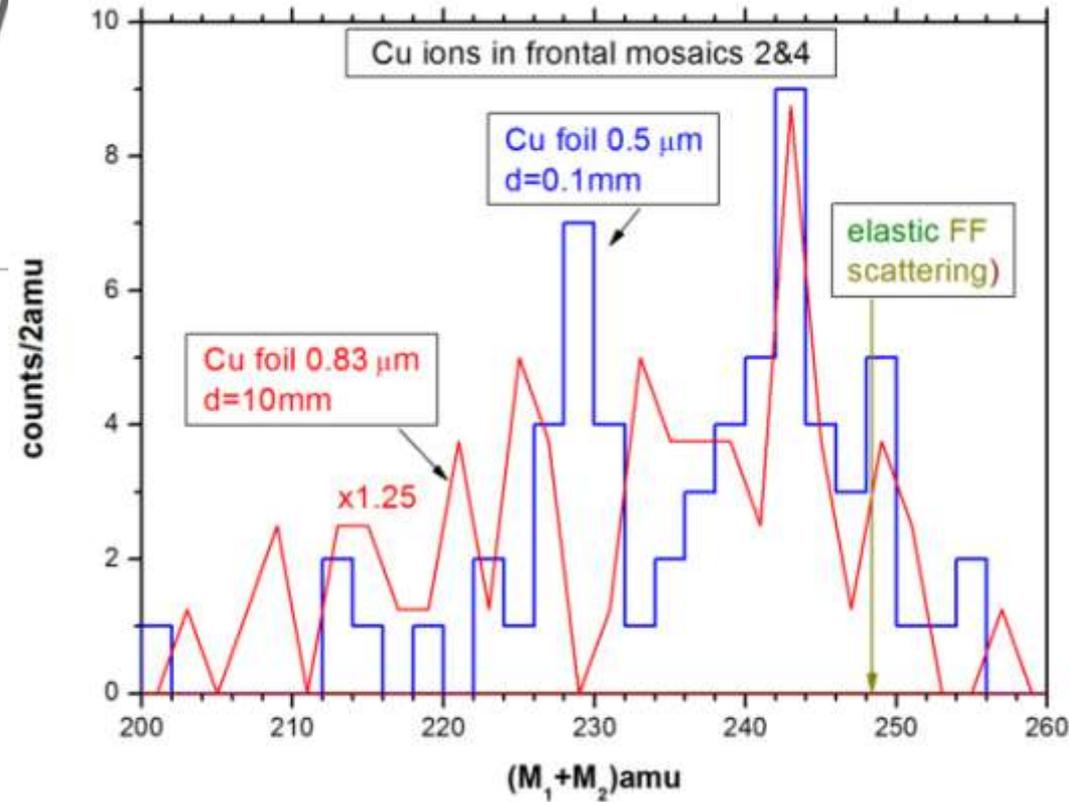
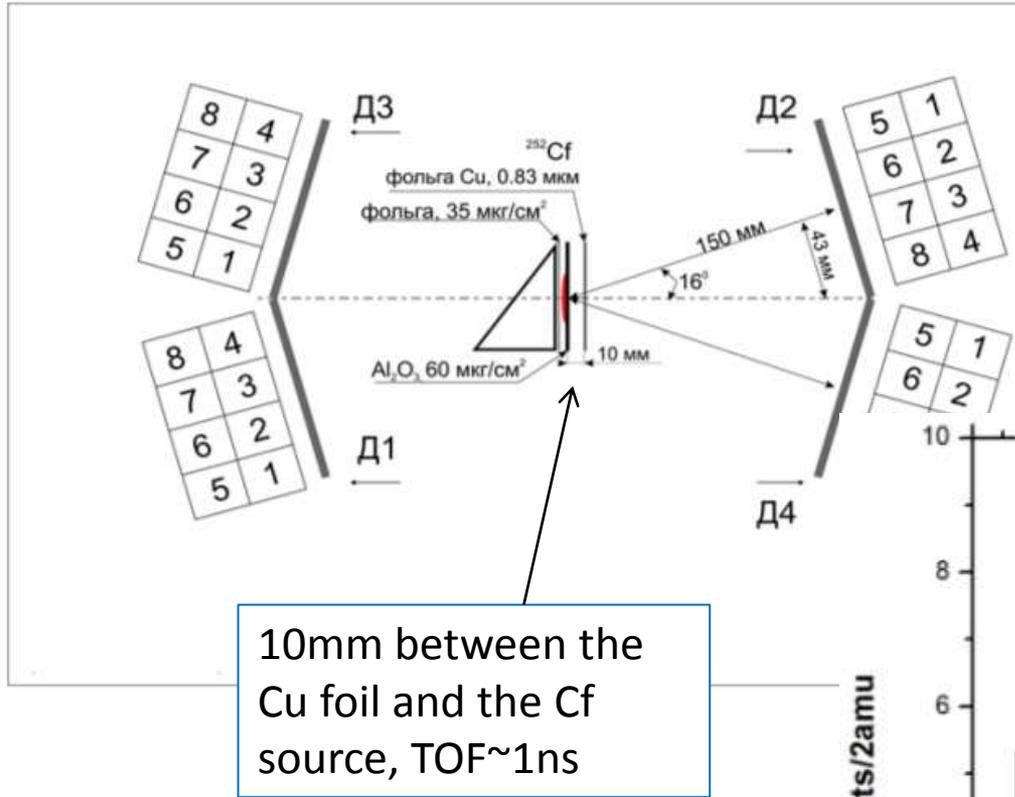
Table 1. Calculated half lives of some quasi-molecular states ^{252}Cf .

| Particle | Fragments | | Q_{exp} (MeV) | K | $\log T$ (s) |
|------------------|-------------------|-------------------|---------------------------|-------|--------------|
| ^{10}Be | ^{132}Sn | ^{110}Ru | 220.183 | 19.96 | -11.87 |
| | ^{138}Te | ^{104}Mo | 209.682 | 25.23 | -9.59 |
| | ^{138}Xe | ^{104}Zr | 209.882 | 26.04 | -9.23 |
| | ^{146}Ba | ^{96}Sr | 201.486 | 22.98 | -10.56 |

There is no such peak in M_1+M_2 spectrum FF from conventional binary fission.
 Likely loosely coupled (see Poenary predictions)
 ^{10}Be based di-nuclear system experiences a breakdown even in tangent collisions.

ps life times are predicted,
low barrier against the decay

Estimation of the life time



Are there contradictions with known features of the fission process?

2. According to our experiments fission fragments in the shape isomeric states show life time at least more than 1ns. What about the neutrons emitted from the fission fragments orders of magnitude faster?

Observation of new microsecond isomers among fission products from in-flight fission of 345 MeV/nucleon ^{238}U

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Abstract

A search for isomeric γ decays among fission fragments from 345 MeV/nucleon ^{238}U has been performed at the [RIKEN Nishina Center RI Beam Factory](#). Fission fragments were selected and identified using the superconducting in-flight separator BigRIPS and were implanted in an aluminum stopper. Delayed γ rays were detected using three clover-type high-purity germanium detectors located at the focal plane within a time window of 20 μs following the implantation. We identified a total of 54 microsecond isomers with half-lives of $\sim 0.1\text{--}10\ \mu\text{s}$, including the discovery of 18 new isomers in very neutron-rich nuclei: $^{59}\text{Ti}^m$, $^{90}\text{As}^m$, $^{92}\text{Se}^m$, $^{93}\text{Se}^m$, $^{94}\text{Br}^m$, $^{95}\text{Br}^m$, $^{96}\text{Br}^m$, $^{97}\text{Rb}^m$, $^{108}\text{Nb}^m$, $^{109}\text{Mo}^m$, $^{117}\text{Ru}^m$, $^{119}\text{Ru}^m$, $^{120}\text{Rh}^m$, $^{122}\text{Rh}^m$, $^{121}\text{Pd}^m$, $^{124}\text{Pd}^m$, $^{124}\text{Ag}^m$, and $^{126}\text{Ag}^m$, and

The fast isotopic separation and identification of reaction products, which take place in several hundred nanoseconds, allow event-by-event detection of isomeric γ rays at the focal plane of the separator with small decay losses in flight. The γ decays are observed under low-background conditions after ion implantation. ...

In-flight fission of a uranium beam have been used as production reactions to populate isomers. In-flight fission is known to be an excellent mechanism for producing neutron-rich exotic nuclei...



Atomic Data and Nuclear Data Tables

journal homepage: www.elsevier.com/locate/adt

Nuclear shape isomers

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^a *Theoretical Division, Los Alamos National Laboratory, Los Alamos, NM 87545, United States*

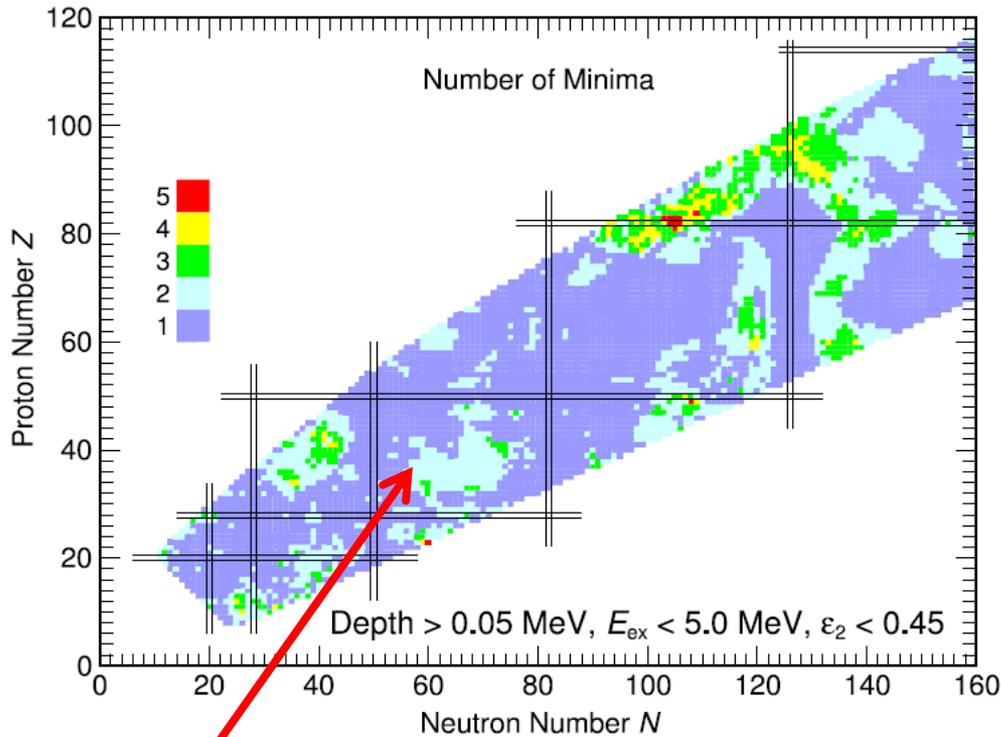
^b *Department of Mathematical Physics, Lund Institute of Technology, P.O. Box 118, SE-22100 Lund, Sweden*

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We calculate potential-energy surfaces as functions of spheroidal (ϵ_2), hexadecapole (ϵ_4), and axial asymmetry (γ) shape coordinates for 7206 nuclei from $A = 31$ to $A = 290$. We tabulate the deformations and energies of all minima deeper than 0.2 MeV and of the saddles between all pairs of minima. The tabulation is terminated at $N = 160$ We also present potential-energy contour plots versus ϵ_2 and γ for 1224 even–even nuclei in the region studied. We can identify nuclei for which a necessary condition for shape isomers occurs, namely multiple minima in the calculated potential-energy surface.

Some results



FF region

Apparently, we and RIKEN group deal with **the same shape isomers** and they observe the de-excitation of isomeric states by detecting of delayed γ -quanta. We fill the presence of di-nuclear system (shape isomer) via brake-up.

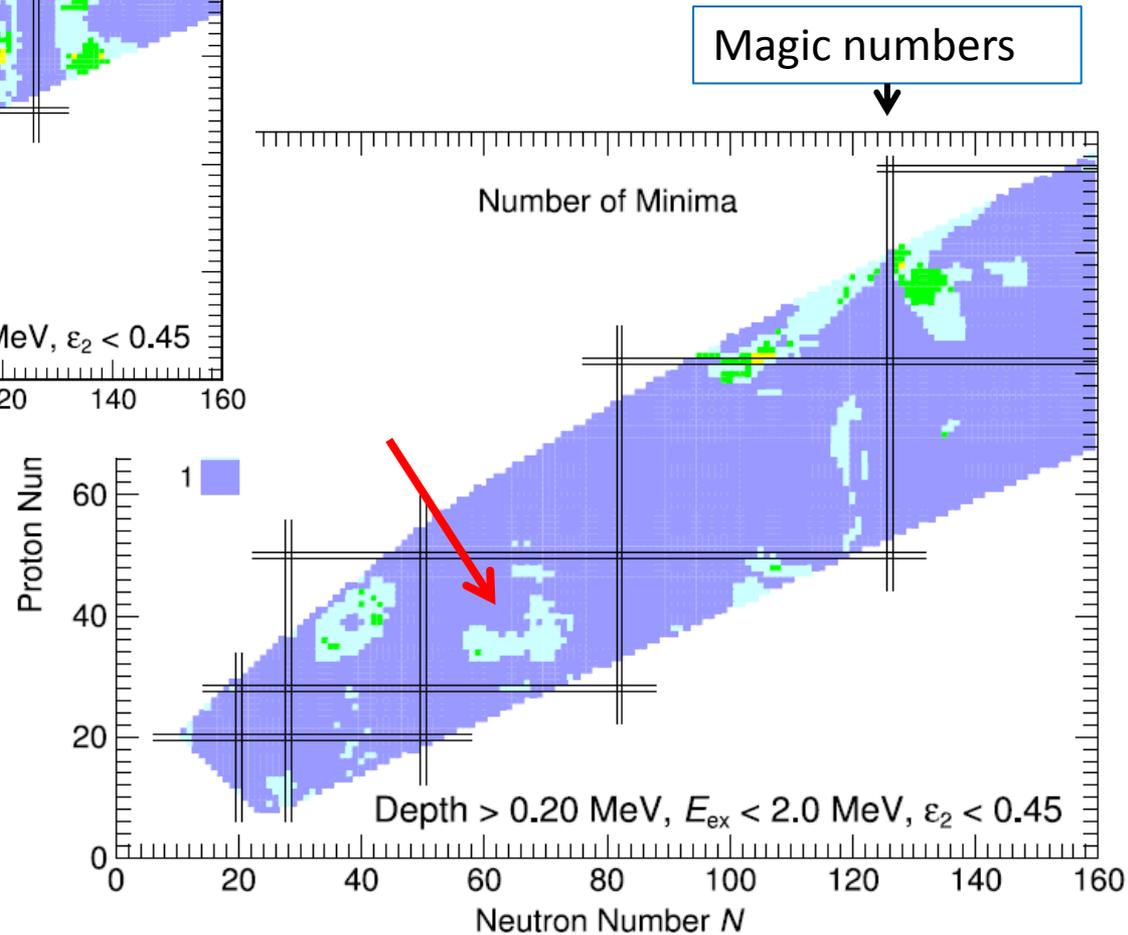


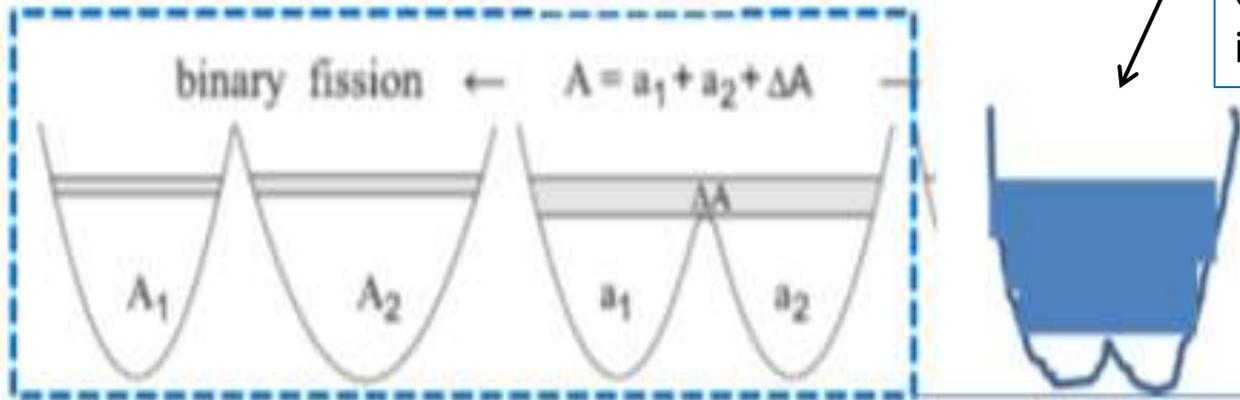
FIG. 5 (color). Number of minima found with deformation $\epsilon_2 < 0.45$. Only the ground-state and isomer minima that are deeper than 0.2 MeV and with energies relative to the ground state of less than 2.0 MeV are counted.

TCSM in description of fission & fusion

True ternary fission of superheavy nuclei

V. I. Zagrebaev,¹ A. V. Karpov,¹ and Walter Greiner²

→ fusion

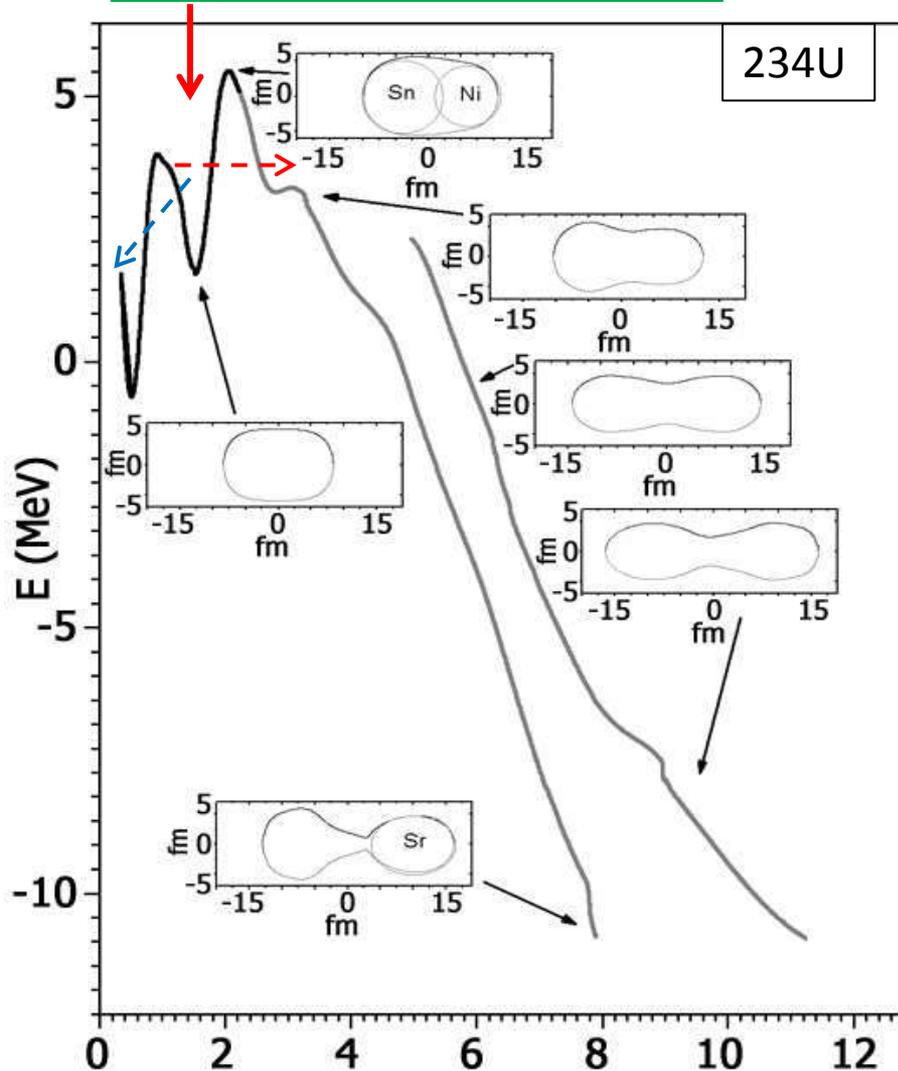


Our case of almost complete fusion of two nuclei constituting initial di-nuclear system

I'd like to stress once more:
 basing on our data we suppose **each FF** to be a **di-nuclear system** at least just after scission

Within the *two-center shell model* (TCSM) for a given nuclear configuration, we may determine the two deformed cores a_1 and a_2 surrounded with a certain number of shared nucleons: $A = a_1 + a_2 + \Delta A$. During binary fission, these valence nucleons gradually spread between the two cores with the formation of two final fragments, A_1 and A_2 .

Shape isomer state (SIS)
in the second potential well.
de-excitation is possible via
gamma channel



The bottoms of the fission valleys as a function of parameter Q (proportional to the quadrupole moment) for ^{234}U .

Treating of our case

Some kind of memory
about initial configuration:
2 distinct nuclei



Almost complete fusion.

Actually it is a low excited
state of the resultant a_1+a_2
nucleus.

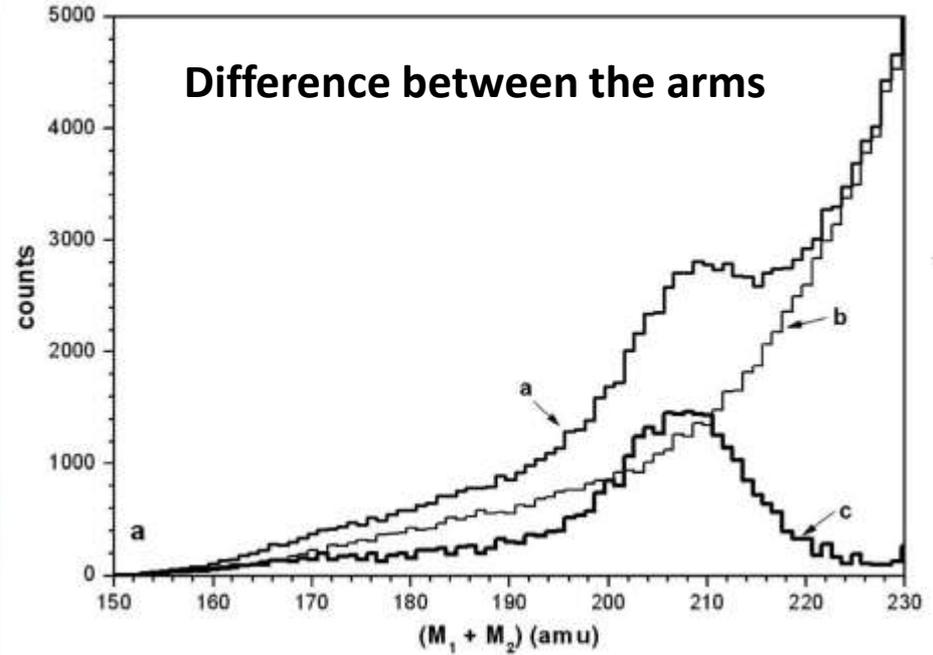
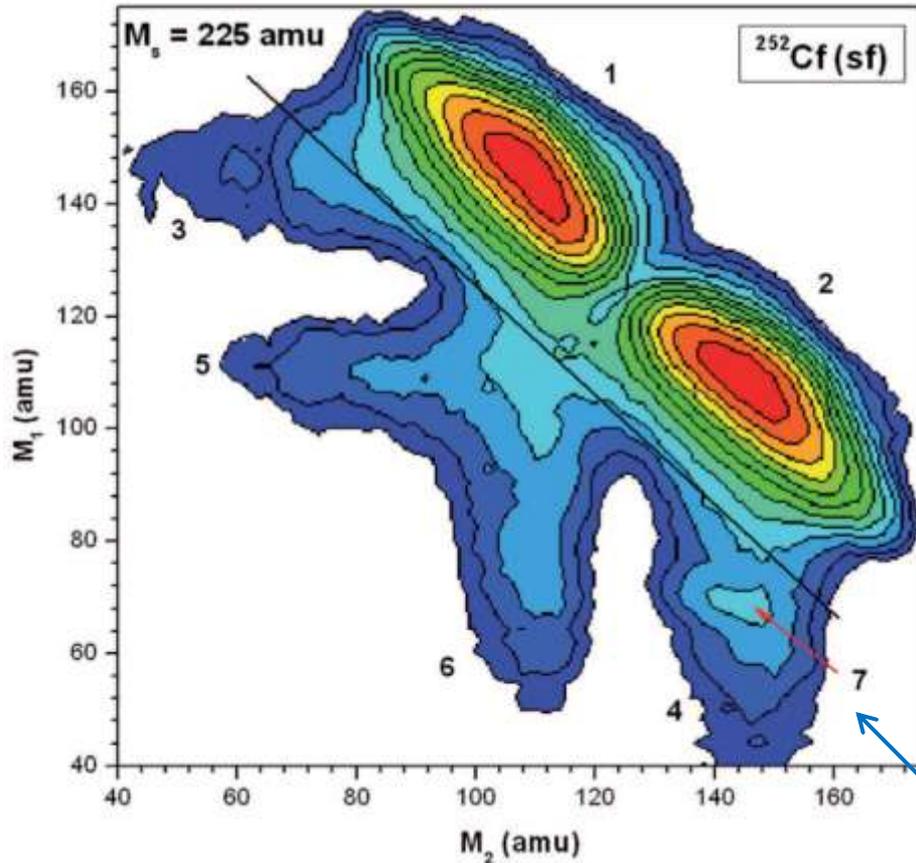
**The main part of the initial
 E^* is *already exhausted by
emitted neutrons.***

The fusion process is
stopped at the final stage
**reaching gamma-isomeric
state.** The residual E^* will
be carry away by gammas
through time to be
characteristic for this
channel of de-excitation.
(some μs presumably)

The last intriguing question to be discussed:

Whether all structures assigned to the CCT actually are resulted from the brake-up process discussed above?

Initial manifestation of the CCT

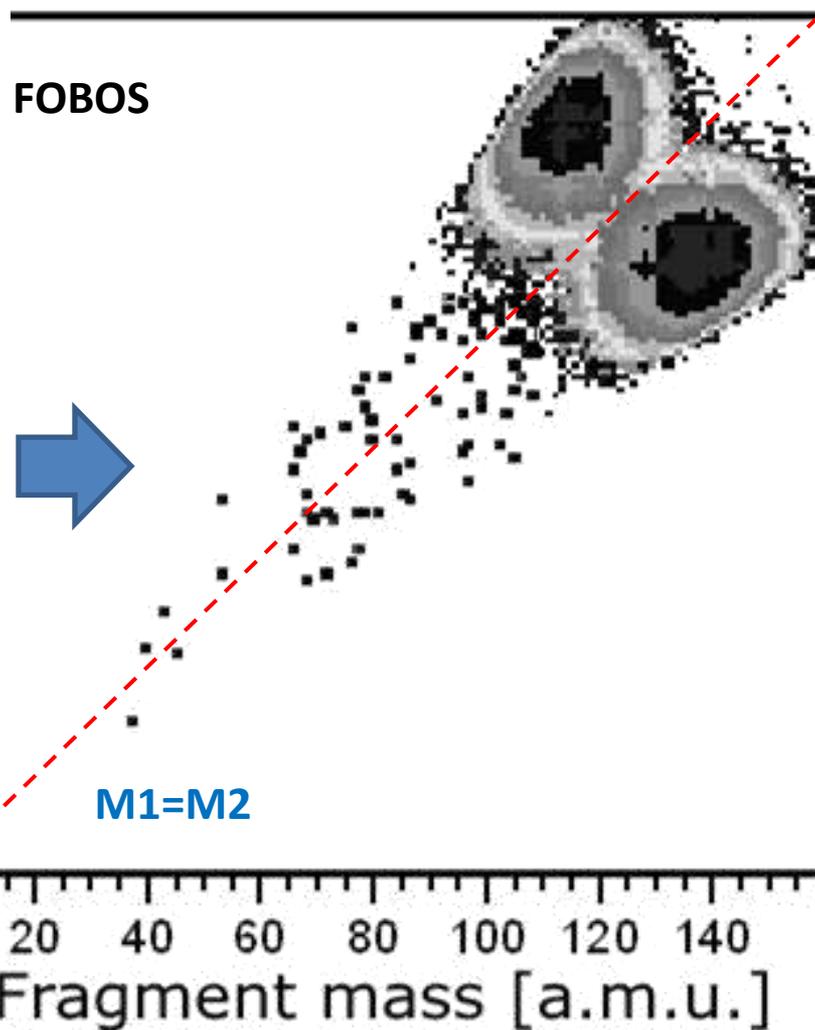
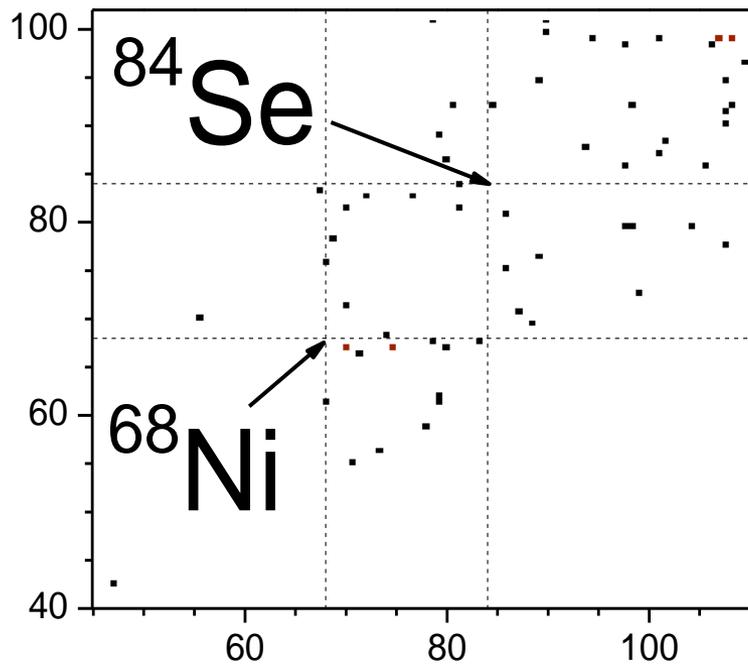
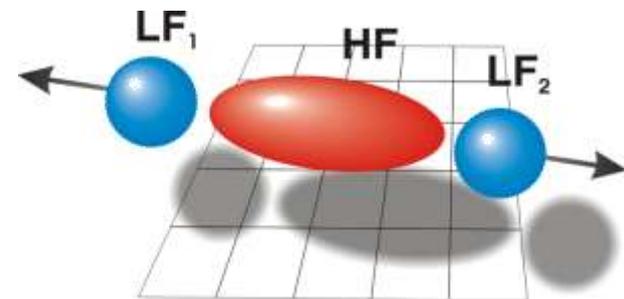


Effect is seen
in the arm from
the side of scattering
foil only.
Thus, it is due to
scattering in any cense

Yu. Pyatkov et al., Eur. Phys. J. A
45, 29–37 (2010)

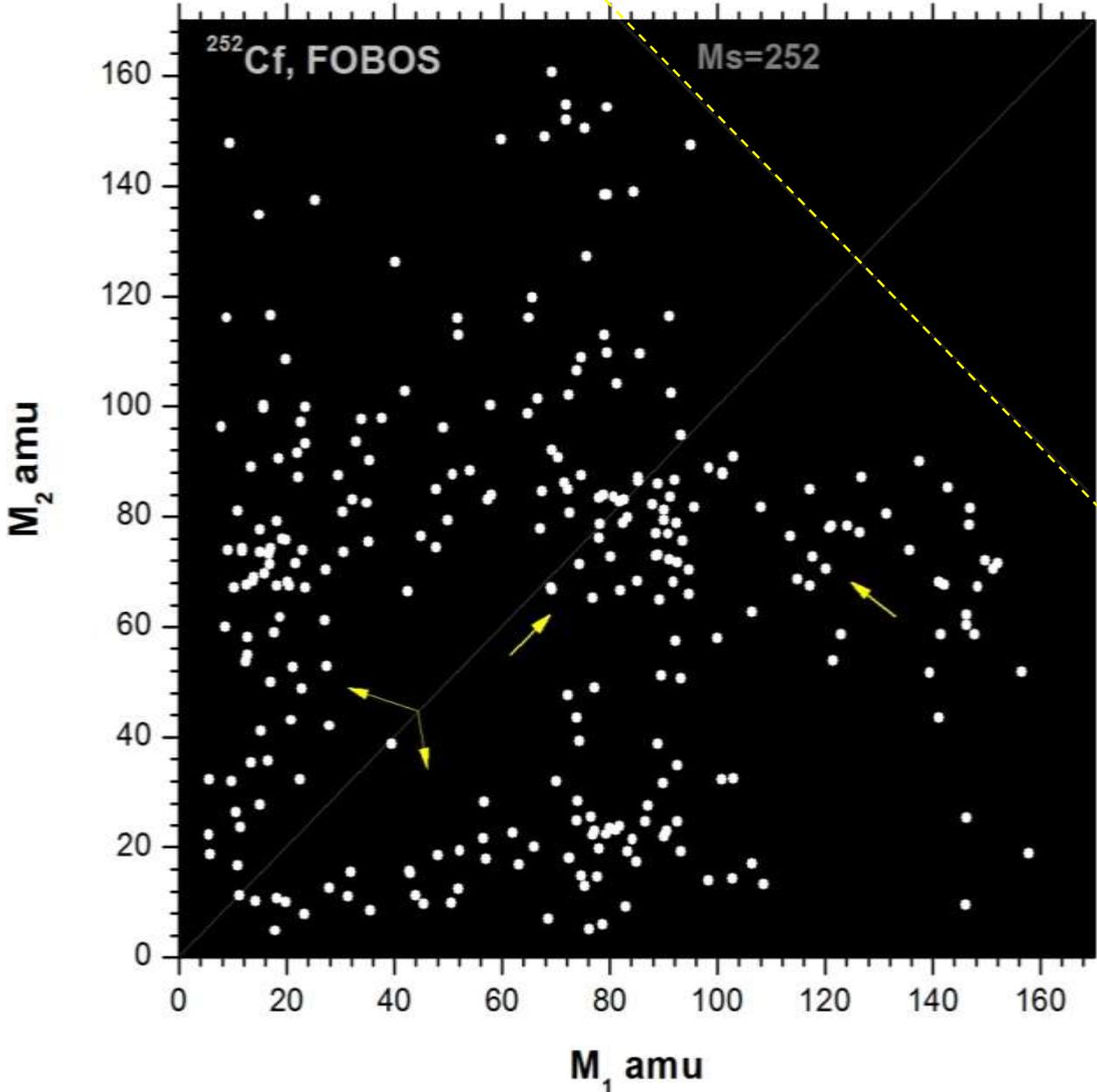
Structures symmetric to the arms

Individual modes / structures
with equal velocity/momentum window



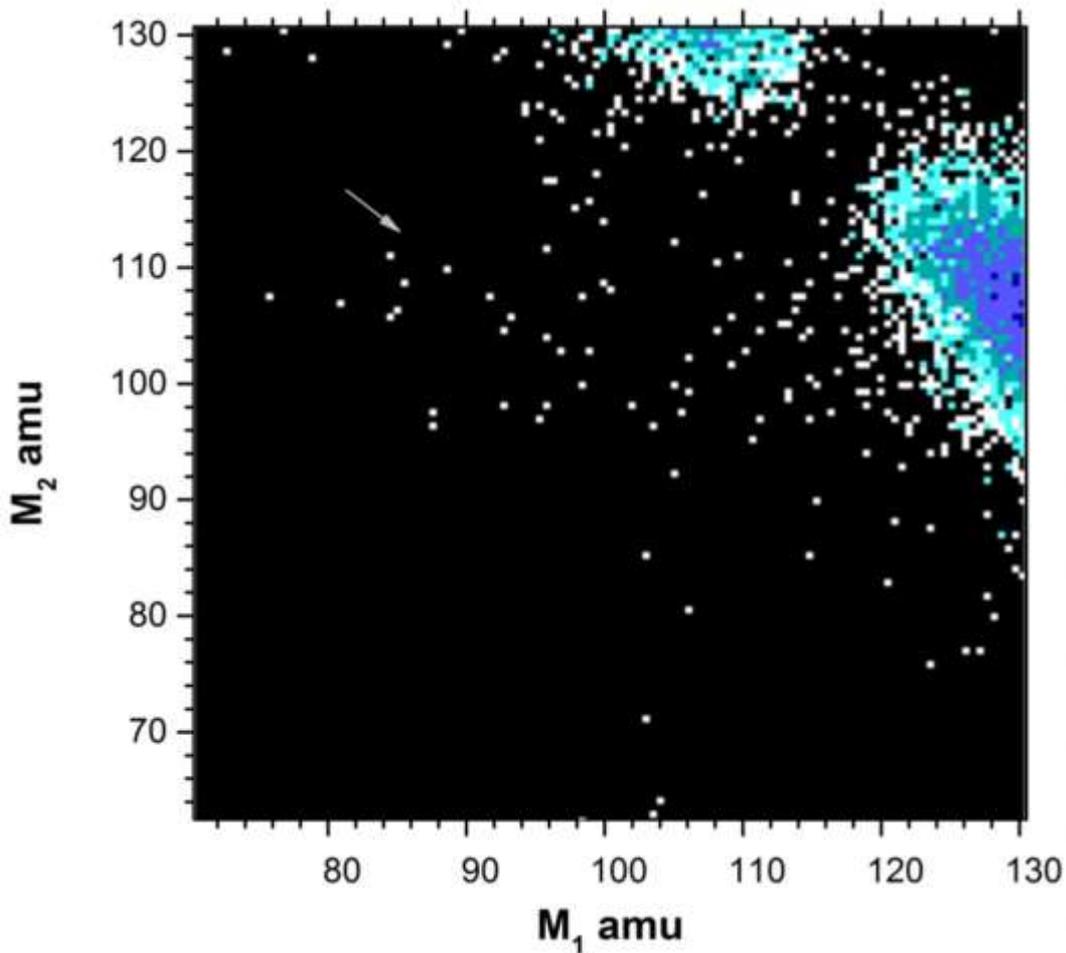
Vice versa:
structures symmetric
to the arms are independent
from the scattering

Structures symmetric to the arms



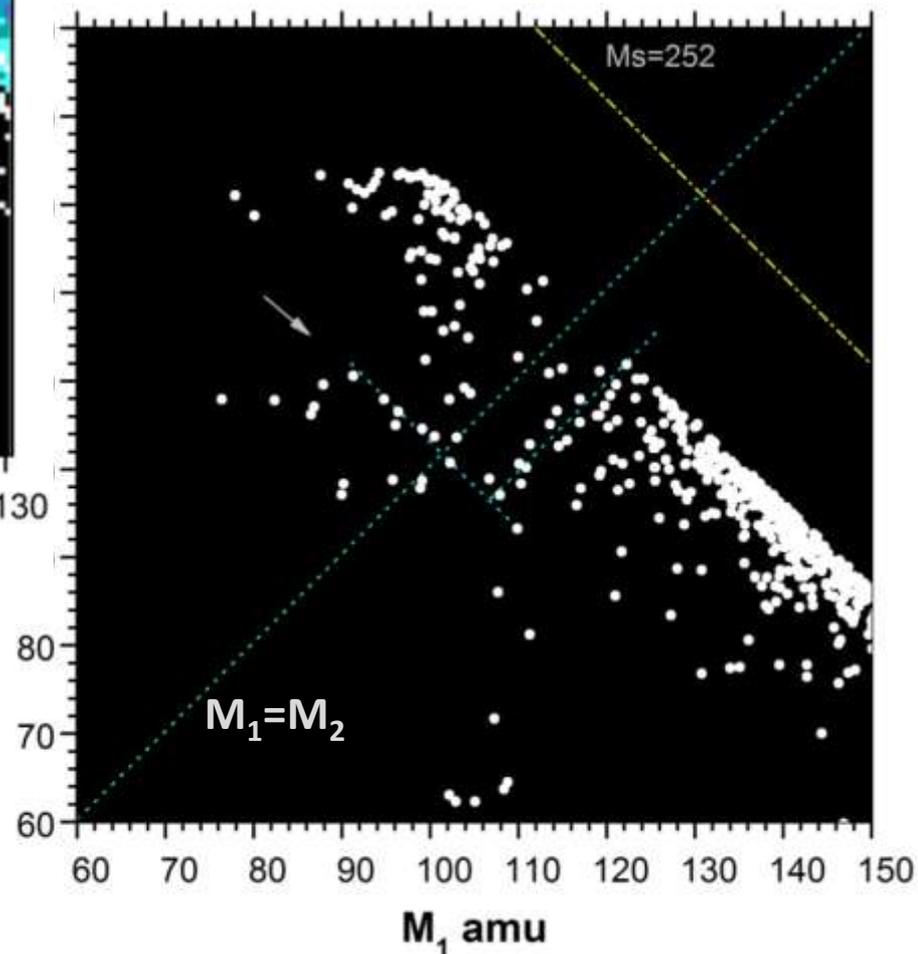
**FOBOS
n=1
&
momentum
box**

Structures symmetric to the arms



COMETA, Cu foil, $n=1$

Missing mass ~ 50 amu

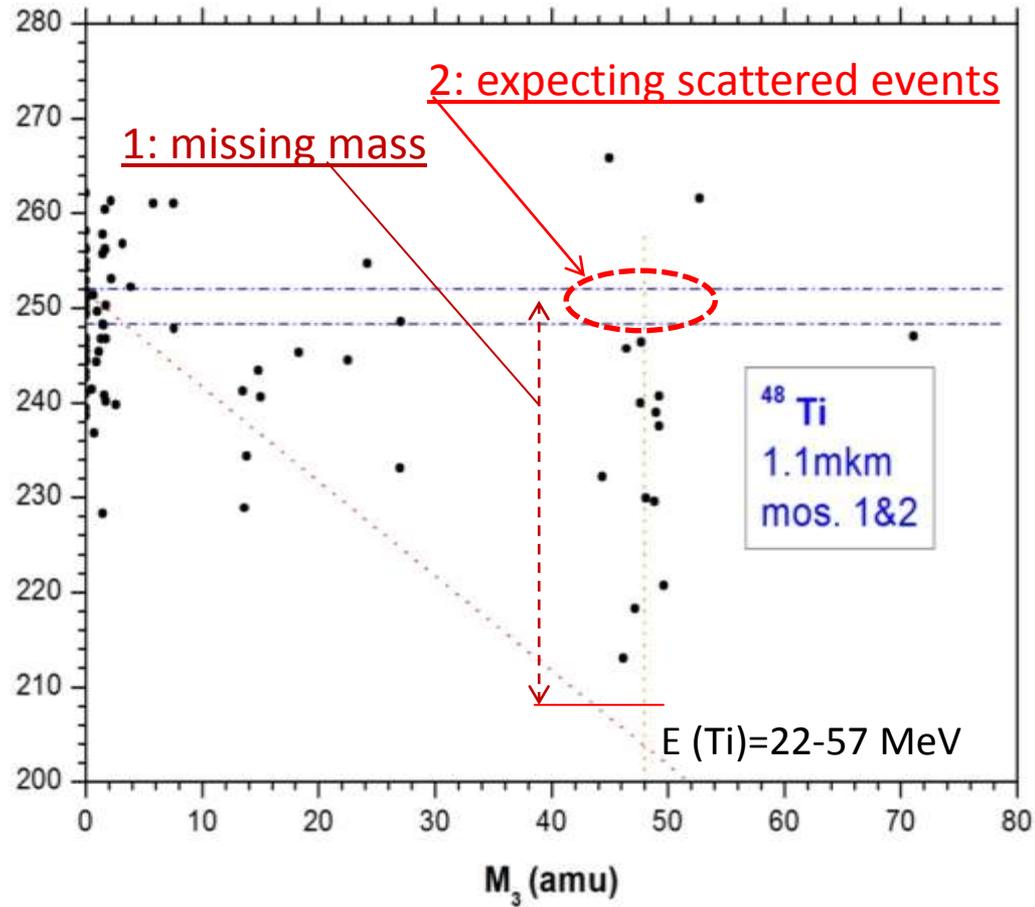
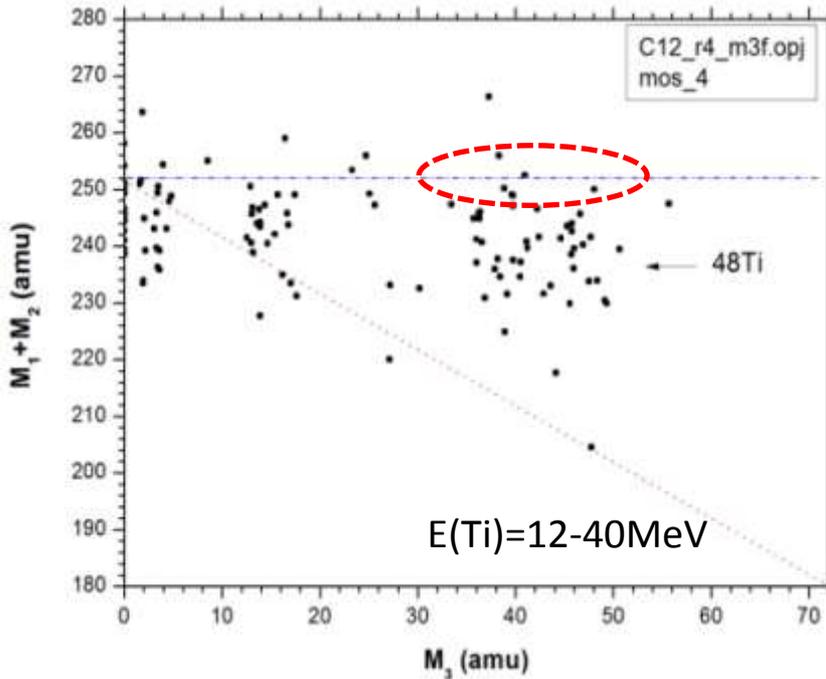
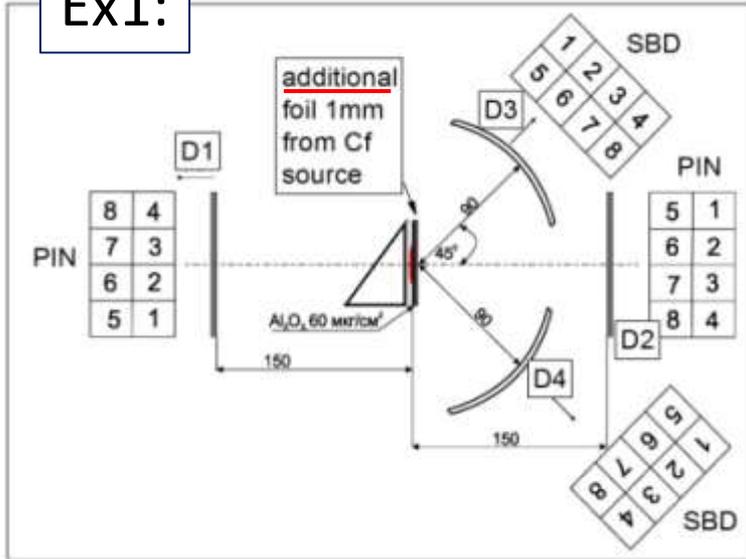


Conclusions

1. Now we have not find any contradictions between all our experimental data obtained and known features of conventional fission process.
2. Studying of rare multi-body decays flashes up fundamental properties of the main process of “conventional” binary fission.

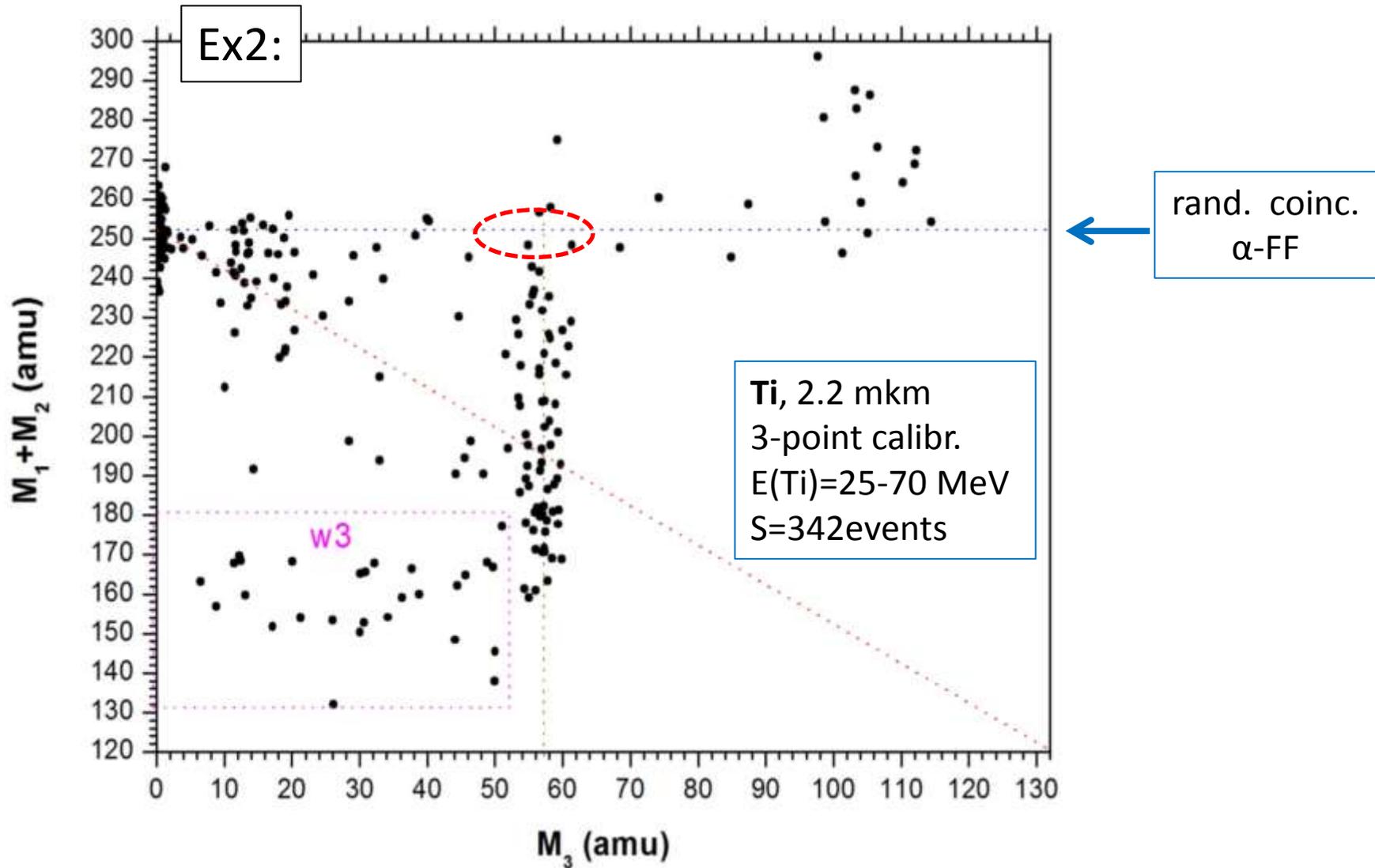
Experimental results_1

Ex1:

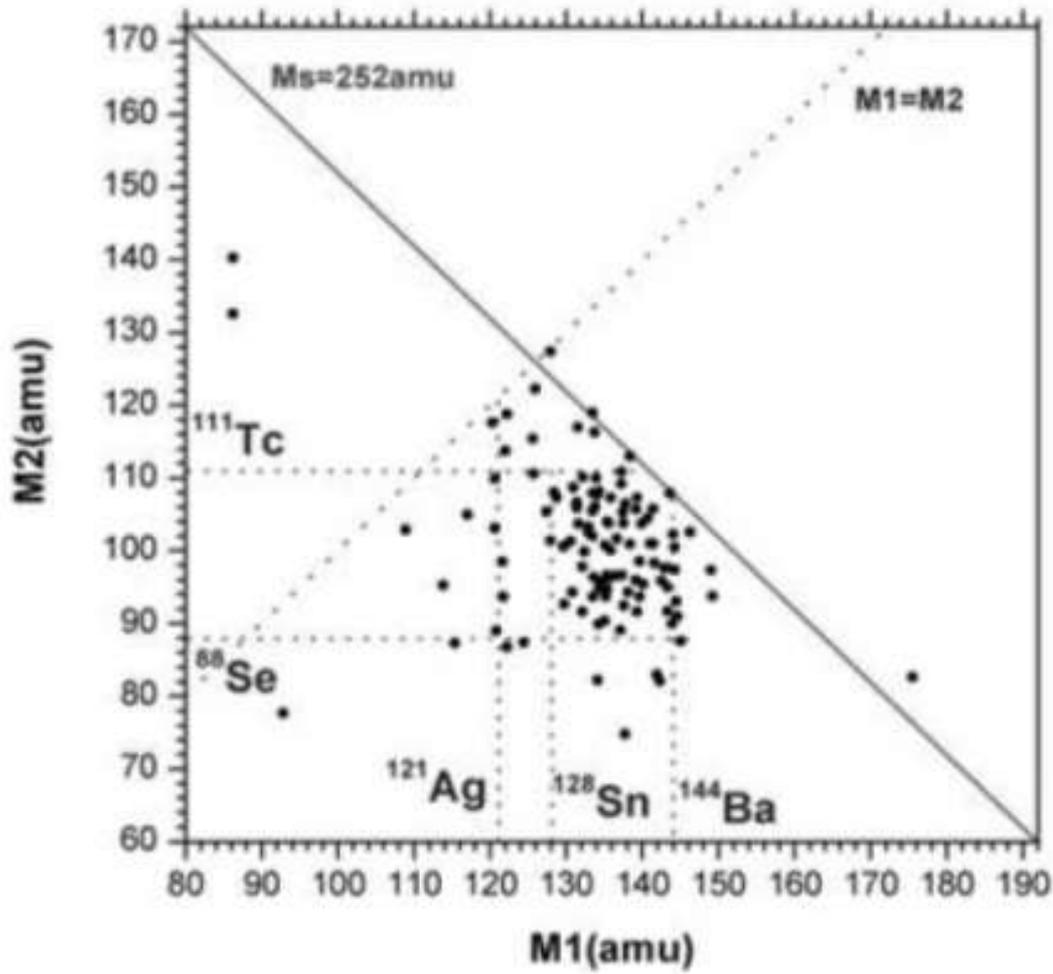


In each ternary event:
 $M_1 > M_2 > M_3$

Experimental results_2

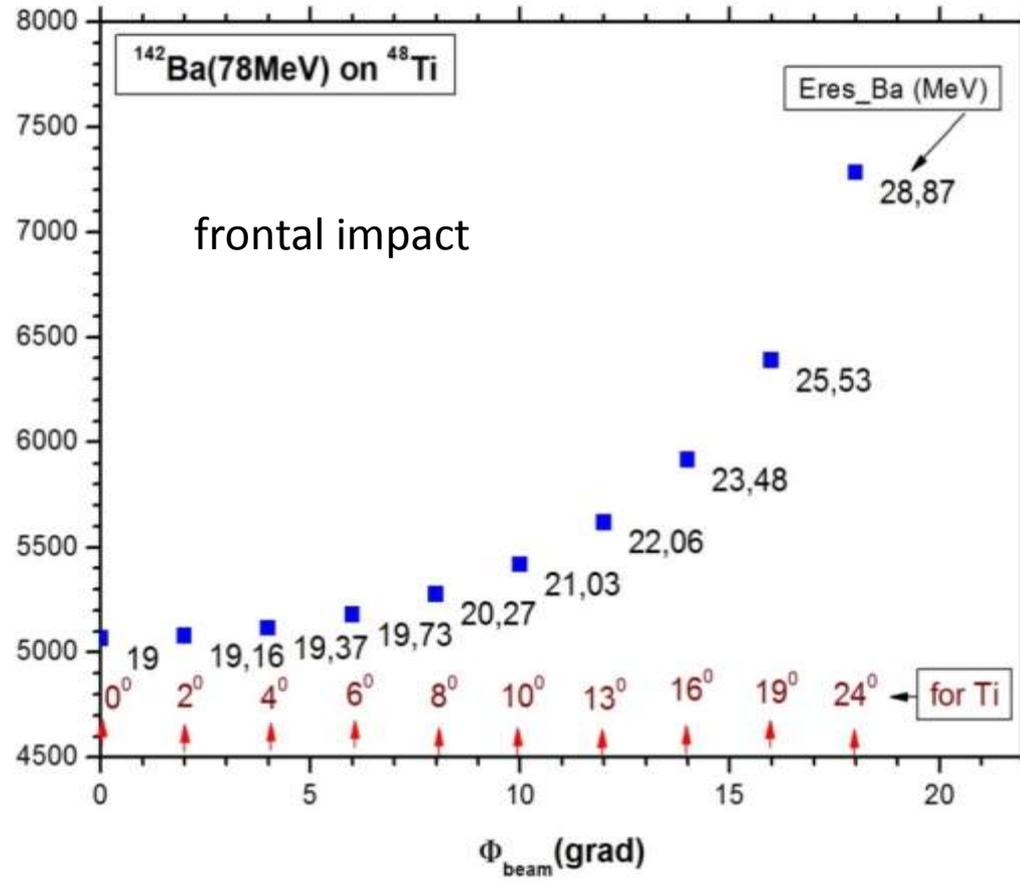
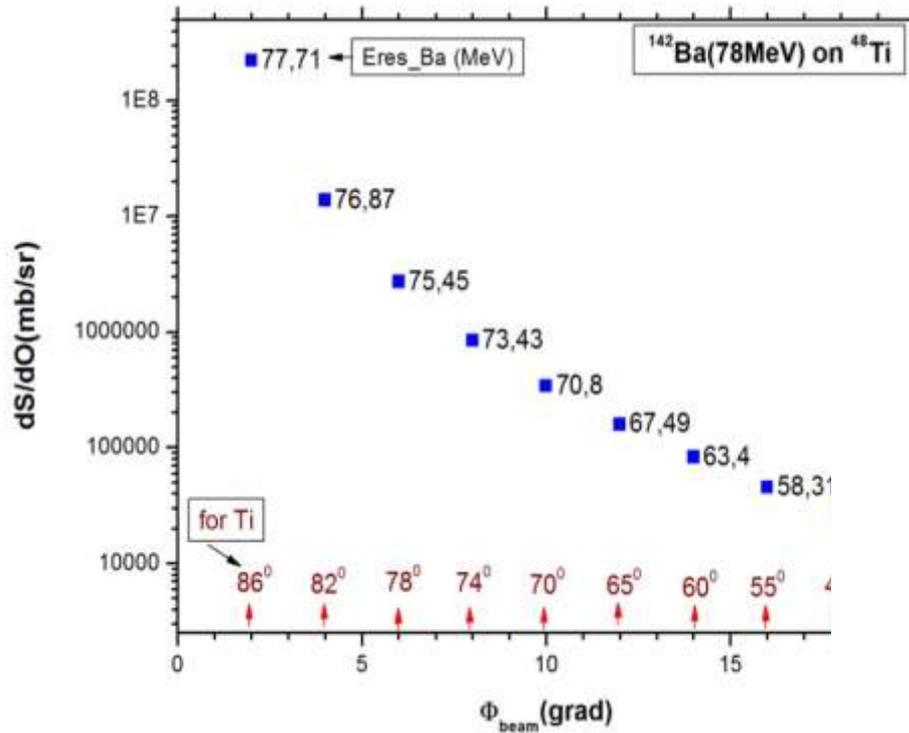


Previous experiments at the COMETA setup - only thin Al_2O_3 backing was in game



Similar structure of the M_1 - M_2 plot for ternary events

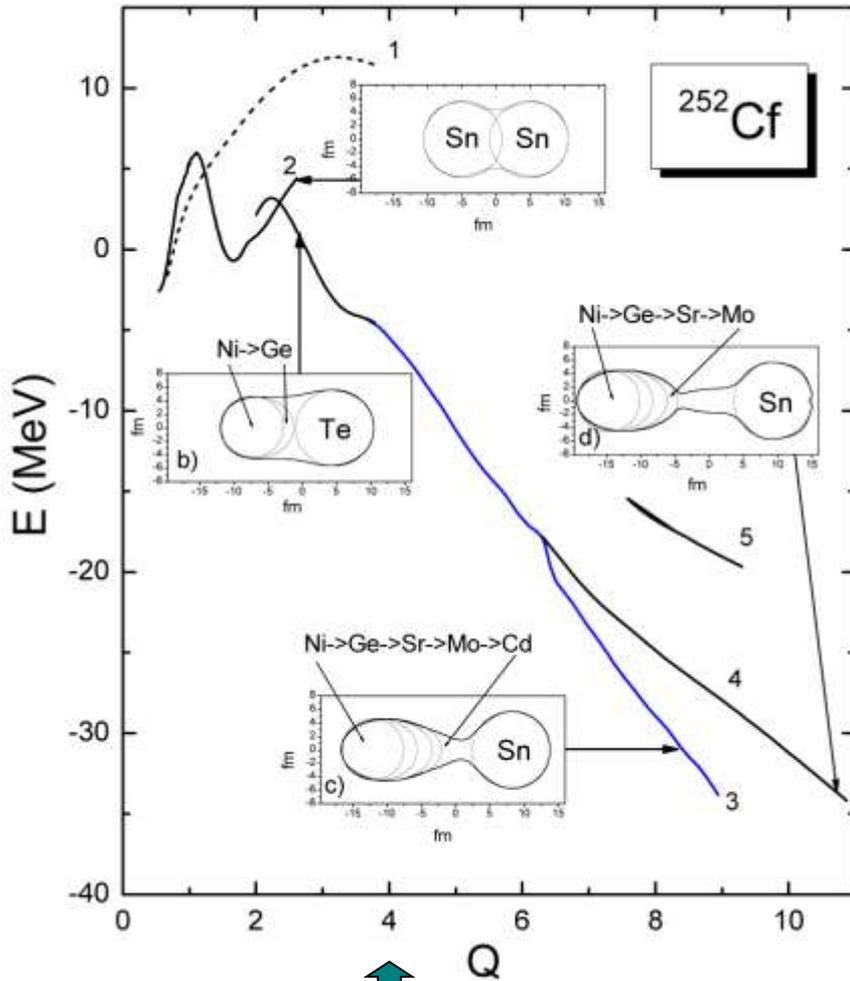
Rutherford scattering – H_FF/Ti



Both the experimental geometry and energies of the knocked-out ions give evidence that we deal with this branch of the scattering process



FF as a di-nuclear system – possible scenario of forming



Double- magic- cluster structure of the fissioning system:

V.V. Vladimirski, JETP (USSR) 5 (1957) 673

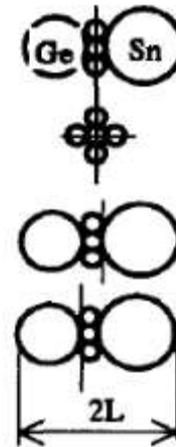
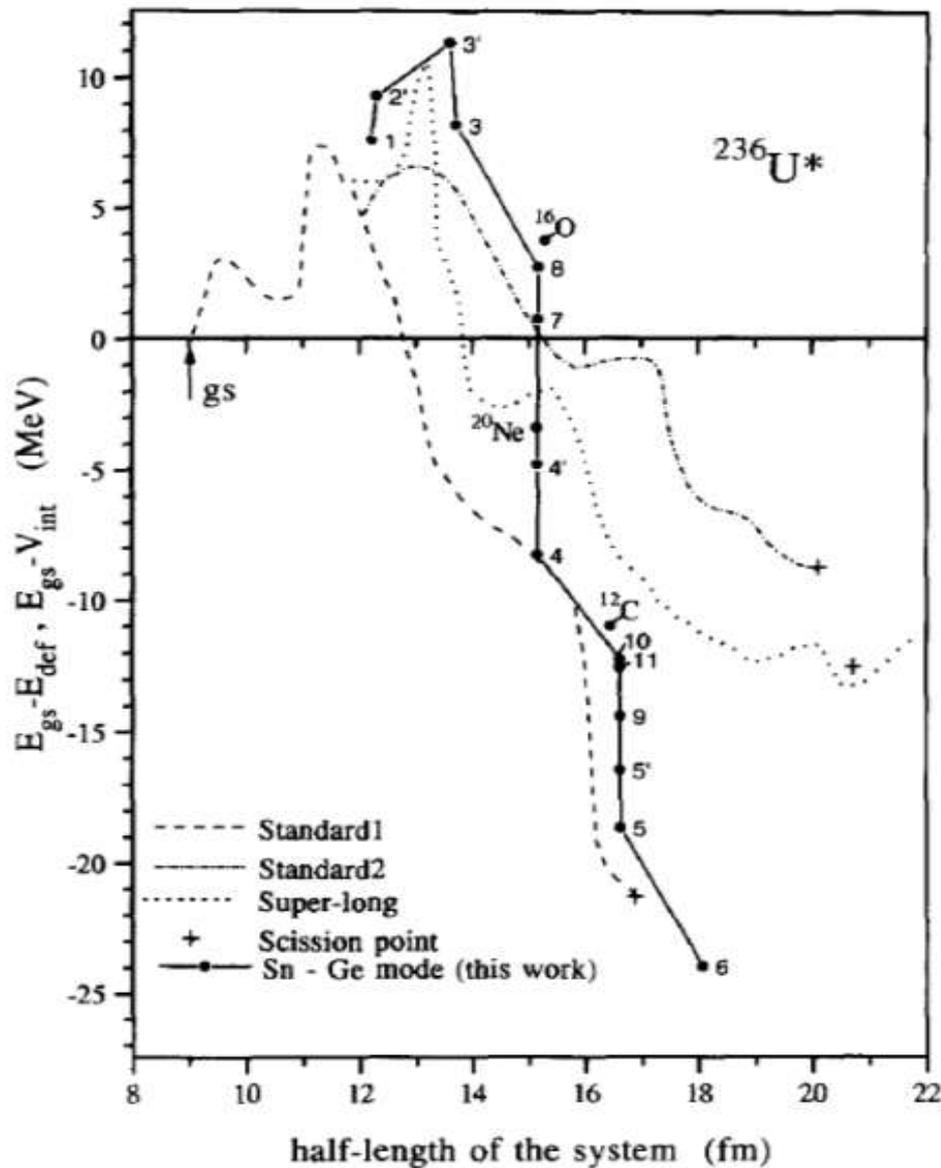
S.L. Whetstone, Phys. Rev. 114 (1959) 581

• • •

I. Tsekanovich,
H.-O. Denschlag, M. Davi,
Z. Büyükmumcu,
F. Gönnerwein, S. Oberstedt,
H.R. Faust
Nucl. Phys. A 688 (2001) 633

Yu.V. Pyatkov, V.V. Pashkevich, Yu.E. Penionzhkevich et al.,
Nucl. Phys. A 624 (1997) 140

FF as a di-nuclear system – possible scenario of forming



Initial configuration
Of the fission mode
Based on Sn & Ge clusters

Two magic clusters namely, light & heavy give rise to fission mode while the neck is also clusterised consisting of LCP

Yu.V. Pyatkov, G.G. Adamian,
N.V. Antonenko et al.,

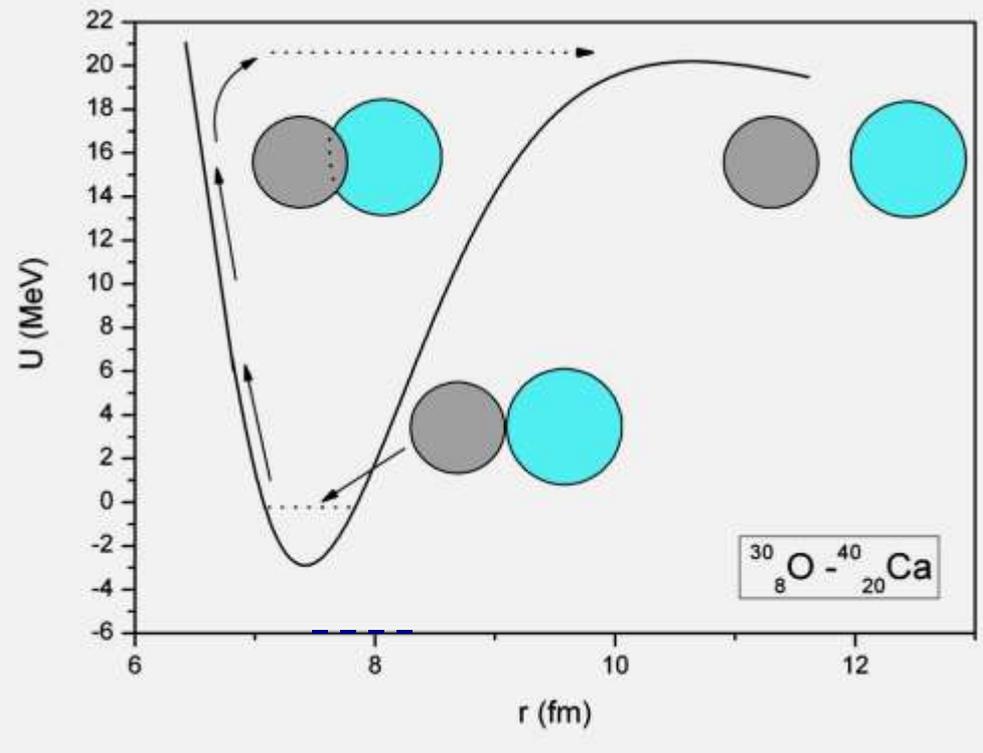
Nucl. Phys. A 611 (1996) 355

Discussion

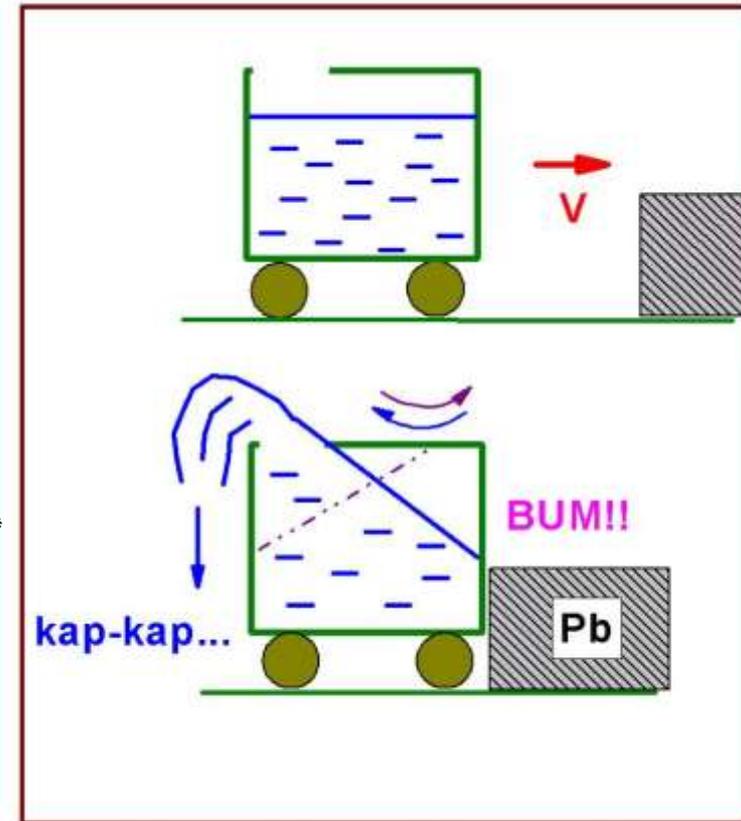
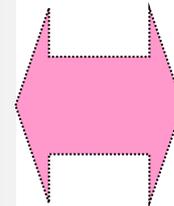
CCT mechanism in the light of the results presented above:

inelastic impact is not exclusive channel of producing ternary events!

A possible way of decaying of di-nuclear system



Different inertia of the partners in the frontal impact could be the reason of their scission.



naive illustration of an inertial effect likely to be decisive for decaying of a nuclear molecule

S. Ćwiok et al.,
 Phys. Lett. B 322 (1994) 304

SIS in
 heavy
 nuclei

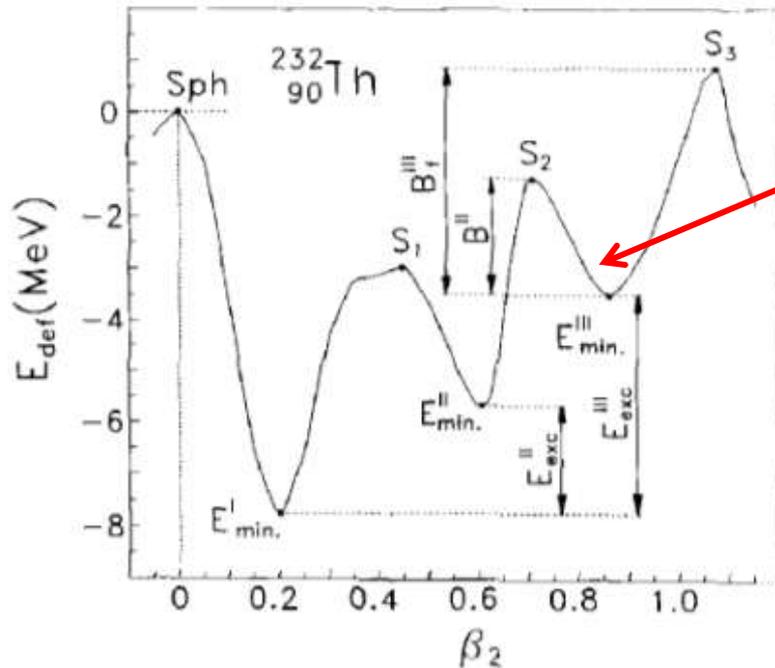
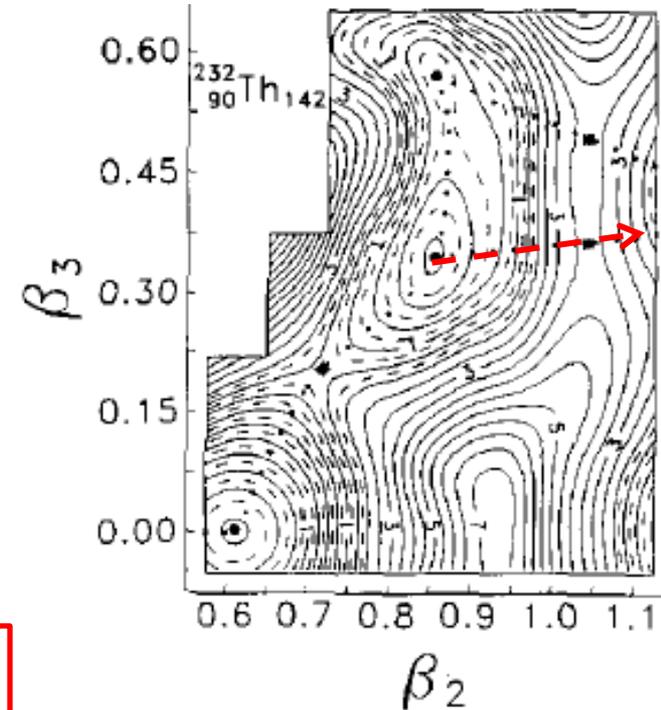
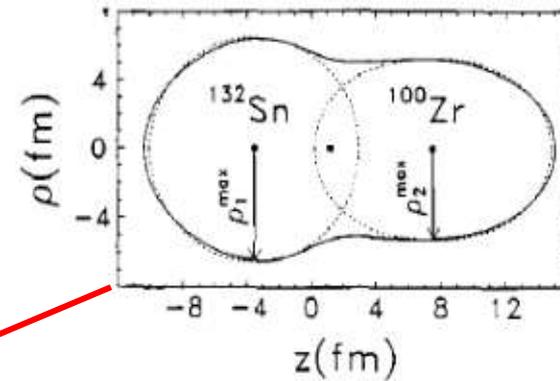


Fig. 1. Potential energy curve for ^{232}Th as a function of quadrupole deformation β_2 along the shorter static fission path of fig. 2.

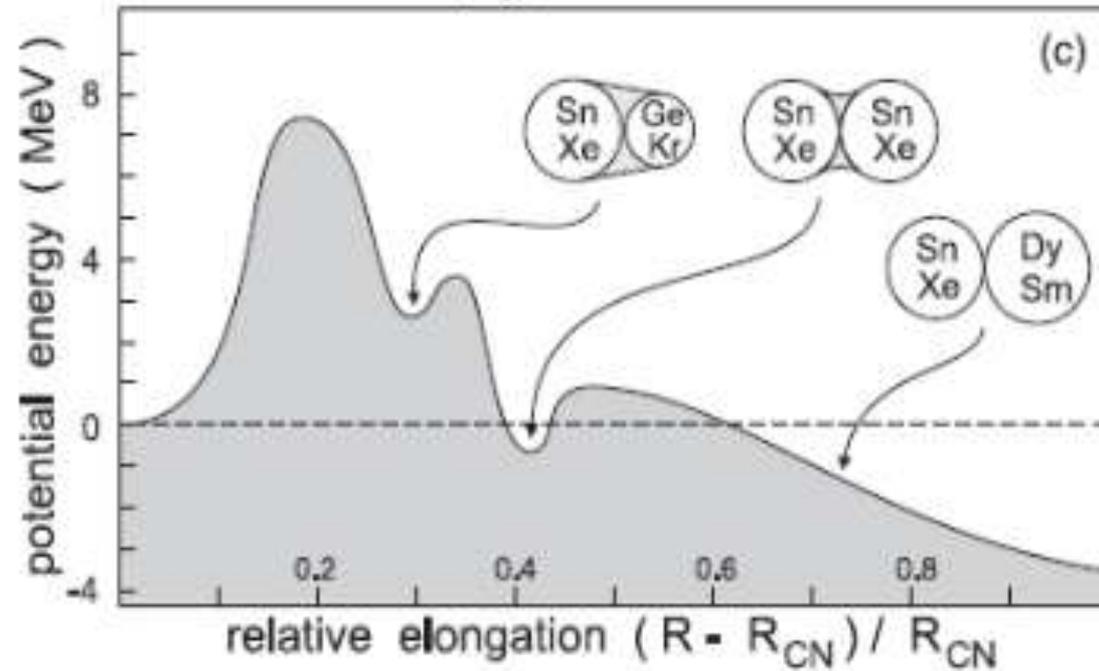


Fiss.
 path

density distribution at the third minimum looks like a di-nucleus consisting of a nearly-spherical heavier fragment (around doubly-magic ^{132}Sn) and a well-deformed lighter fragment (from the neutron-rich $A \sim 100$ region).

Fig. 2. The Woods-Saxon-Strutinsky total potential energy (relative to the spherical macroscopic energy) for ^{220}Rn , ^{222}Ra , ^{232}Th , and ^{234}U , as a function of β_2 and β_3 . At each (β_2, β_3) point the energy was minimized with respect to $\beta_4-\beta_7$. The distance between the solid contour lines is 0.5

SIS in super-heavy nuclei



Three-humped barrier calculated along the fission path of $^{296}_{116}\text{Lv}$ (Livermorium).

V. ZAGREBAEV, W. GREINER

Proc. Int. Symp. on Atomic Cluster Collisions (ISACC07), GSI Darmstadt, 2007, (Imperial College Press, London, 2008), Eds. J.-P. Connerade and A. V. Solov'yov, p. 23

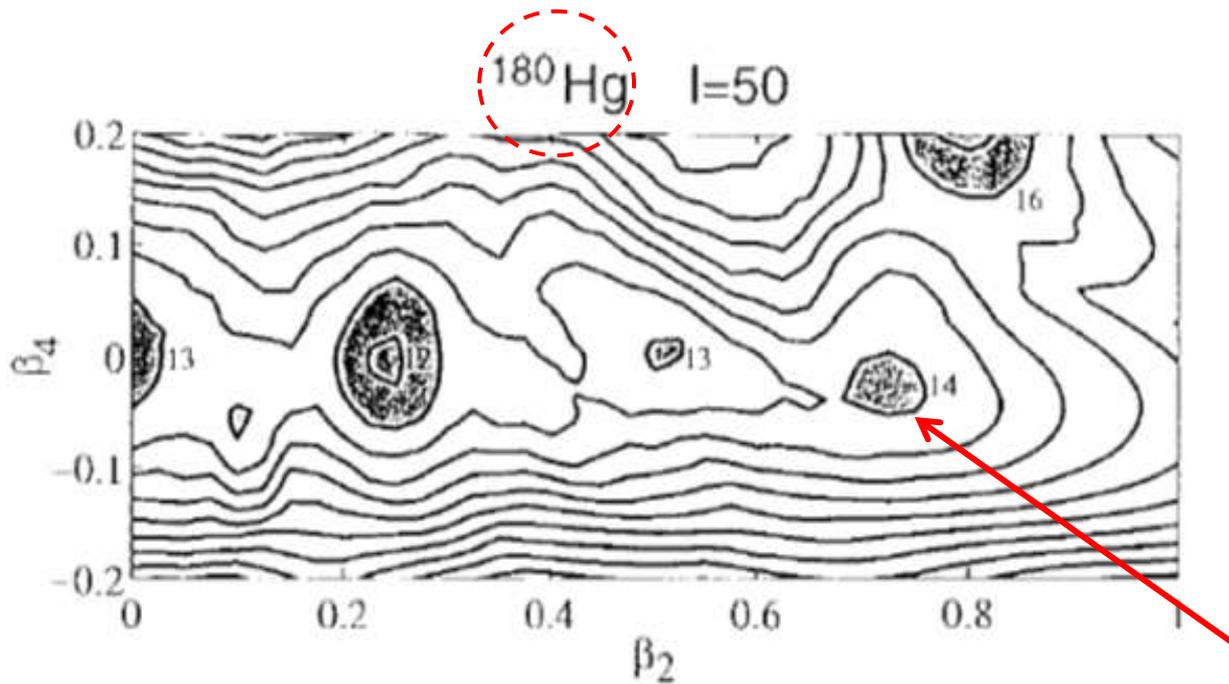
“These intermediate minima correspond to the shape isomer states.

From analysis of the driving potential we may definitely conclude that these isomeric states are nothing else but the two-cluster configurations with magic or semi-magic cores surrounded with a certain amount of shared nucleons.”

Shape isomers at high spin

SIS at
high spin

Sven Åberg et al., Z. Phys. A 358 (1997) 269



“Superdeformed (or hyperdeformed) nuclei with necking was calculated to exist, e.g. in ^{180}Hg . The exotic configuration was similar in shape as well as in single-particle structure to **two partly overlapping spherical 90Zr nuclei.**”

Fig. 5. Potential-energy surface valid for ^{180}Hg at $I=50$. The calculation has been performed within the cranked Nilsson-Strutinsky formalism using the Woods-Saxon potential. Local minima are shaded, and the line separation is 1 MeV. Notice the minimum at $\beta_2=0.75$, $\beta_4=0$ that corresponds to a necked-in superdeformed shape

The *cluster states of light nuclei* and the possible existence of the *necked-in shaped nuclei* were considered in:

Cseh, J., Scheid, W. J.: Phys. G 18 (1992) 1478

Cseh, J. et al.,: Phys. Rev.C 48 (1993) 1724

Sanders S.J. et al.,: Phys. Rep. 311 (1999) 487

Freer, M : Prog. Phys. 70 (2007) 2149

Sciani, W et al., : Phys. Rev. C80 (2009) 034319

Cseh, J. et al.,: Phys. Re v. C 80 (2009) 034320

Beck, C. et al.,: Phys. Rev. C 80 (2009) 034604

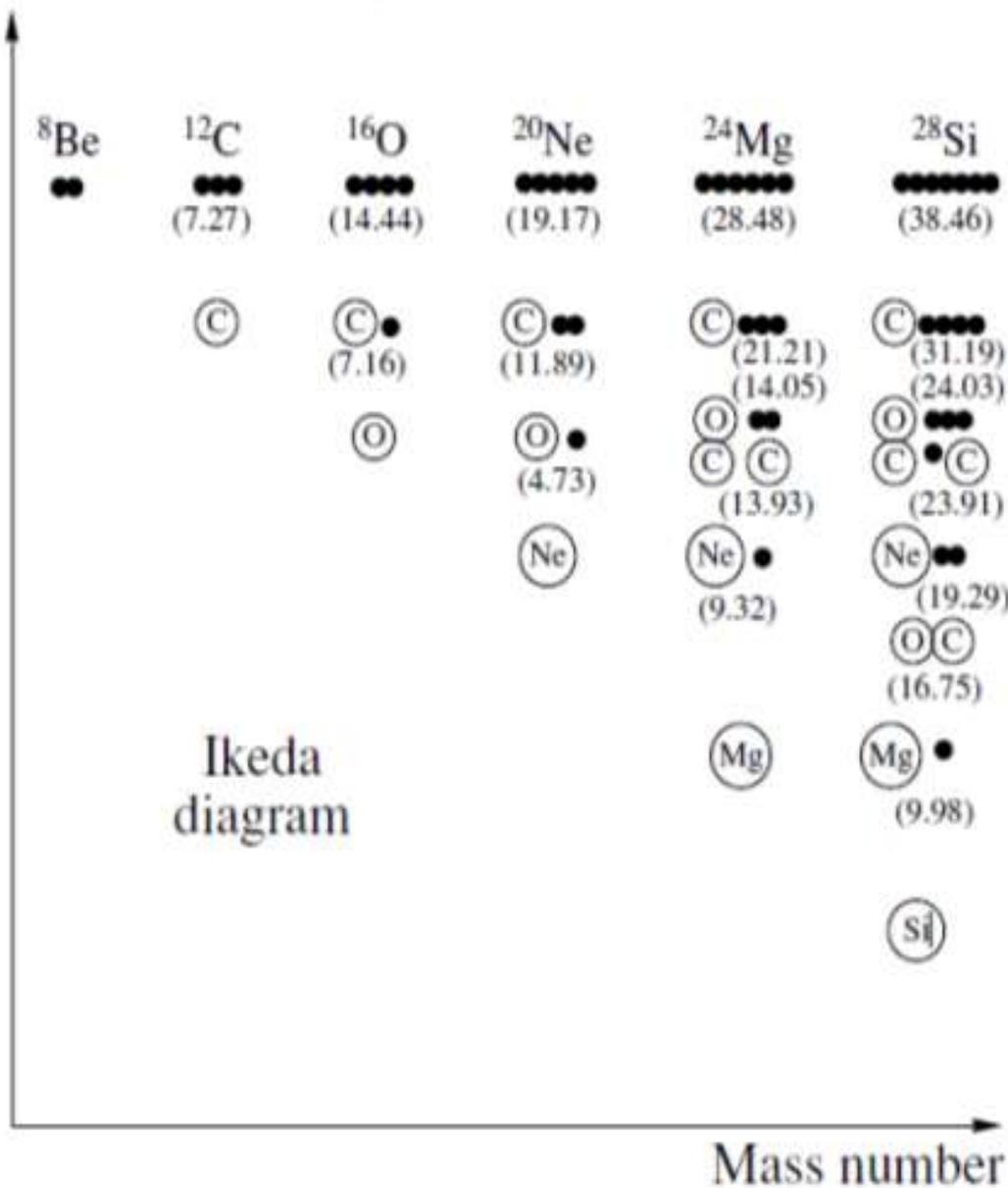
Von Oertzen W. et al.,; Phys. Rev. C 78 (2008) 044615

We likely deal with shape-isomers in the new mass rang typical for the fission fragments of the conventional binary fission

Conclusions.

- 1. New mechanism of ternary decay based presumably on the Rutherford break-up of the fragment in the shape-isomeric state is observed.**
- 2. Break-up is only one of the different ways leading to the CCT.**
- 3. The results obtained let us to suppose that the bulk of the fragments from the conventional binary fission are born in the shape-isomeric states.**
- 4. The conclusions above can be regarded as the preliminary ones till further estimation of the life times of the shape isomers under discussion will be obtained.**

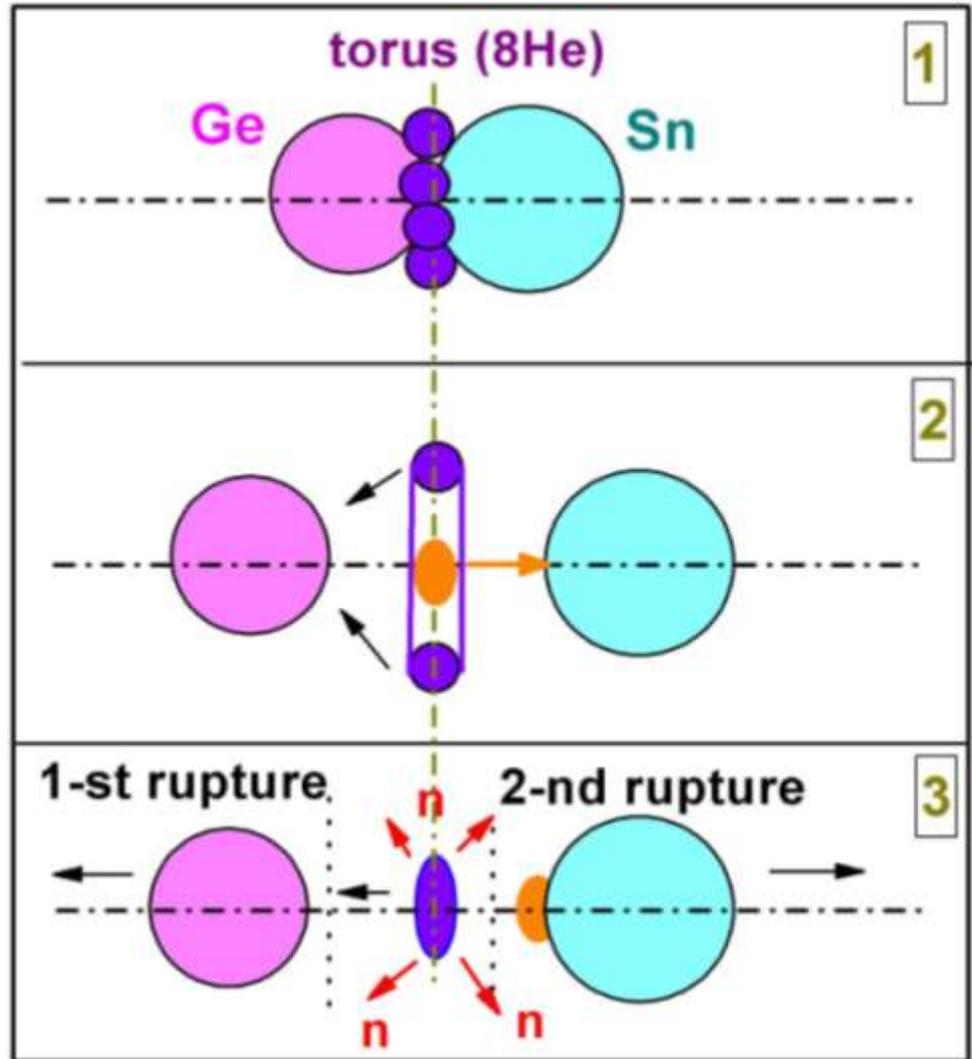
Excitation energy



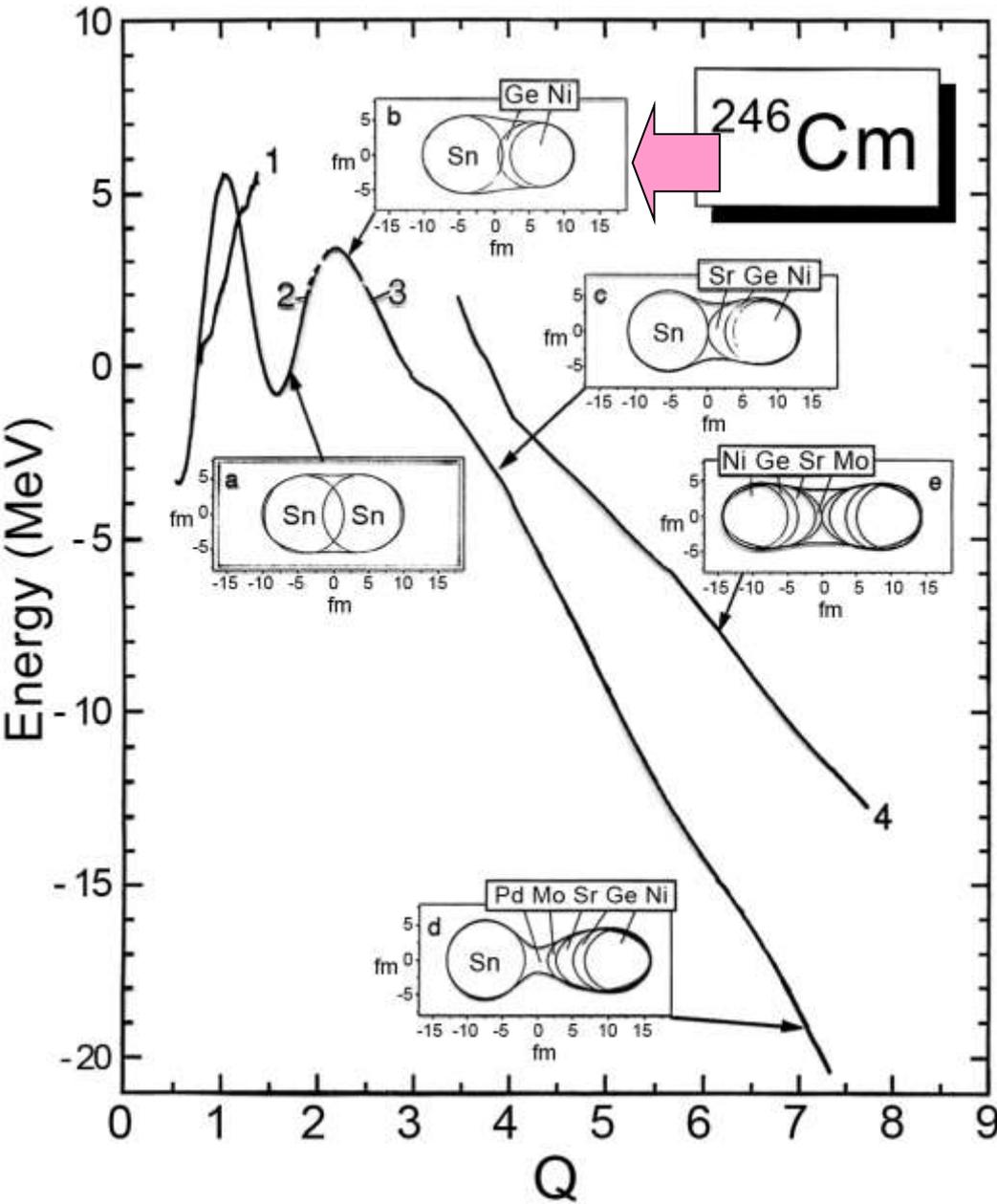
Scenario of the collinear cluster tripartition

Discovery of fission (shape) isomers: FLNR (JINR) 1961
 ^{242}Am , $\tau = 0.014\text{sec}$

More than 30 fission isomers of heavy nuclei, namely, isotopes of U, Np, Pu, Am, Cm, Bk are known including short lived in the ns range. (Flerov, Polikanov)



Calculated Fission Valleys (^{246}Cm)



Valley of the mass-asymmetrical shapes

Valley of the mass-symmetrical shapes

V. V. Pashkevich et al.

Known calculations

Table 1: TKE for some configurations of the fissioning system $^{236}\text{U}^*$

| lsh+l | Scission configuration | TKE MeV | lsh+l | Scission configuration | TKE MeV | lsh+l | Scission configuration | TKE MeV |
|-------|------------------------|---------|-------|------------------------|---------|-------|------------------------|---------|
| 1 | | 196 | 5c | | 141 | 10a | | 151 |
| | | | 5d | | 138 | 10b | | 149 |
| | | | 5d' | | 130 | 10c | | 137 |
| | | | 5b' | | 135 | 10d | | 132 |
| | | | 5c' | | 138 | 11a | | 153 |
| | | | 5d'' | | 143 | 11b | | 139 |
| | | | 6a | | 135 | 11c | | 136 |
| | | | 6b | | 132 | 11d | | 130 |
| | | | 6c | | 126 | 12a | | 206 |
| | | | 6d | | 122 | | | |
| 6e | | 121 | | | | | | |
| 6f | | 117 | | | | | | |
| 2a | | 209 | 7a | | 169 | 12b | | 184 |
| | | | 7b | | 168 | 12c | | 183 |
| | | | 7c | | 150 | 12d | | 171 |
| 3a | | 193 | 8a | | 172 | 13a | | 193 |
| | | | 8b | | 157 | | | |
| | | | 8c | | 149 | | | |
| 8d | | 147 | | | | | | |
| 4a | | 166 | 9a | | 149 | 3'a | | 182 |
| | | | 9b | | 147 | 14a | | 164 |
| | | | 9c | | 142 | | | |
| 5a | | 148 | 9d | | 132 | 14b | | 145 |
| | | | 5b | | 145 | | | |

α -cluster configurations analyzed in:

Yu.V. Pyatkov, G.G. Adamian,
N.V. Antonenko, V.G. Tishchenko
Nucl. Phys. A 611 (1996) 355

Thus pre-formation of light clusters in the neck region and just forming of two and three-neck shapes are energetically possible. At the same time nothing is known from theory about probability of the decay modes discussed here.

Known calculations

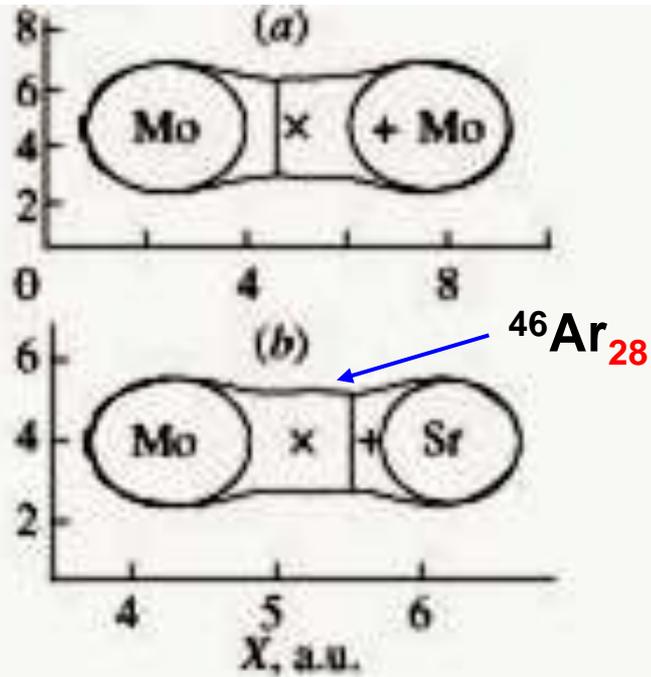
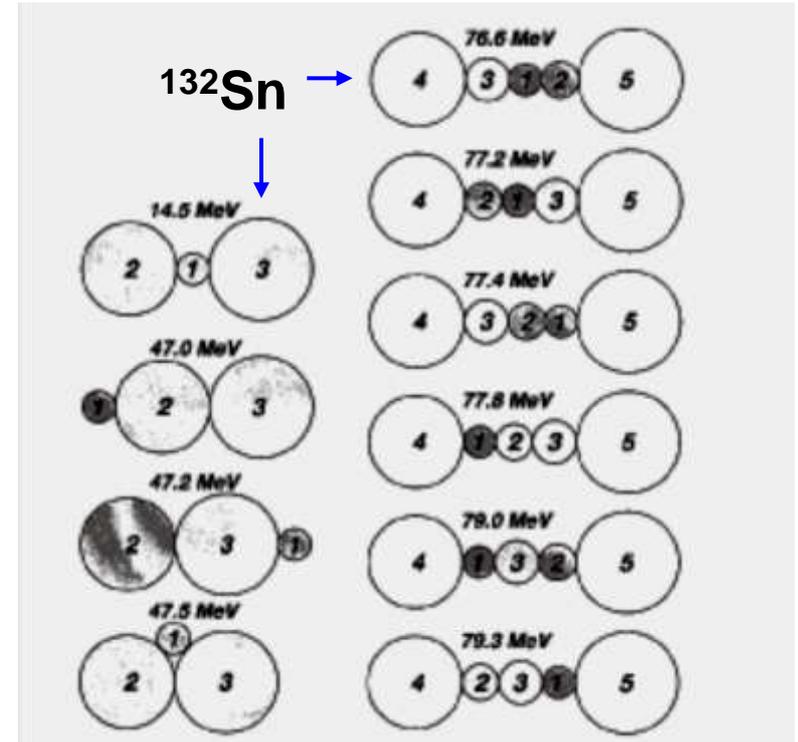


Fig. 7. The shape of the nucleus at the bottom of the "symmetric" valley ($Q_2 = 7.52$ a.u., $\eta = 0.074$) (a); the same system at the point $Q_2 = 7.52$ a.u., $\eta = -0.208$ (b).

Yu.V. Pyatkov, V.V. Pashkevich, A.V. Unzhakova et al., *Physics of Atomic Nuclei* 66 (2003) 1631



Aligned and compact configurations for α -accompanied and $\alpha+^6\text{He}+^{10}\text{Be}$ accompanied cold fission of ^{252}Cf
 D.N. Poenaru et al., *Phys. Rev. C* 59 (1999) 3457