

Neutron diffraction by a moving grating: theory and experiment

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Outline

- *Prehistory*
- *Moving grating as a non stationary phenomenon.
(kinematic theory)*
- *First observation of the energy splitting*
- *Application: focusing in time*
- *Application: gravity experiments*
- *Found problems*
- *Progress in theory I - improved kinematic approach*
- *Progress in theory II- dynamical theory of diffraction*
- *New experimental approach to the measurement of
diffraction spectra*
- *Conclusion*

СООБЩЕНИЯ
ОБЪЕДИНЕННОГО
ИНСТИТУТА
ЯДЕРНЫХ
ИССЛЕДОВАНИЙ
ДУБНА



P4 - 8851

И.М.Франк

О ВОЗМОЖНОЙ ПРИЧИНЕ АНОМАЛИИ
ВО ВРЕМЕНИ ХРАНЕНИЯ
УЛЬТРАХОЛОДНЫХ НЕЙТРОНОВ

1975

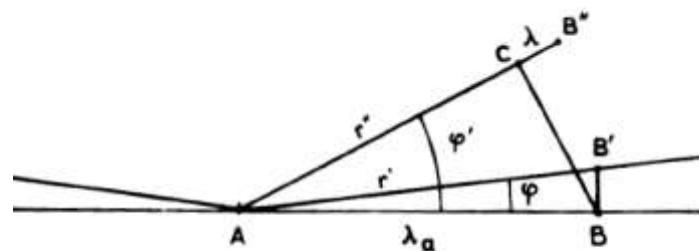


Рис. 2

**UCN diffraction by the Raleigh waves
results in changing of the normal
components of the k-vector**

I.M. Frank. On the possible issue of the anomaly in UCN storage

Neutron Diffraction by Surface Acoustic Waves

W. A. Hamilton, A. G. Klein, and G. I. Opat

School of Physics, University of Melbourne, Parkville, Victoria 3052, Australia

and

P. A. Timmins

Institut Laue-Langevin, 38042 Grenoble Cedex, France

(Received 3 March 1987)

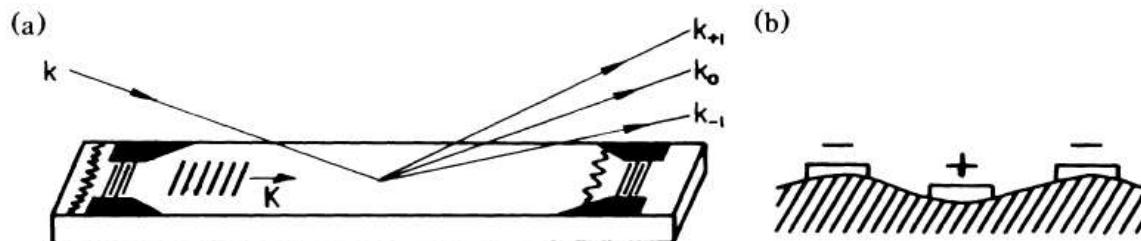
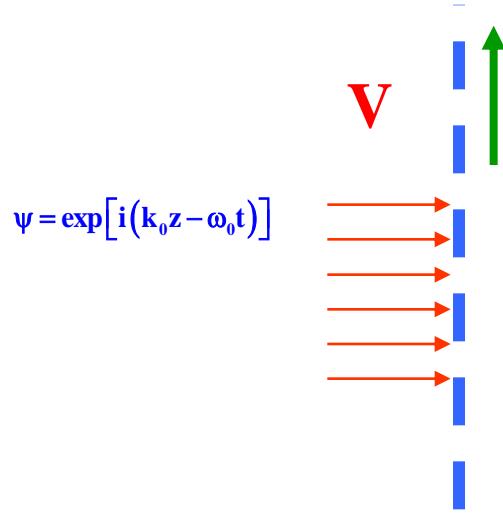
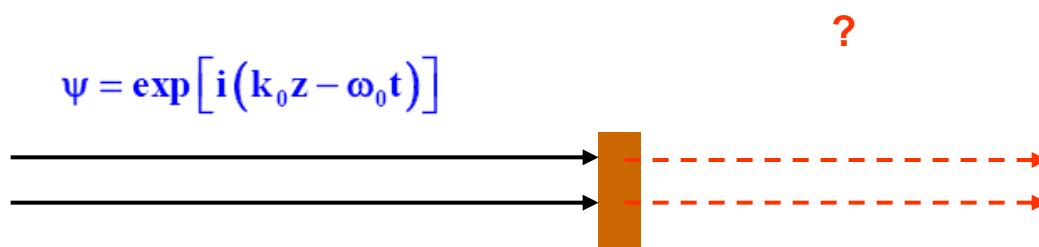


FIG. 1. (a) Schematic of neutron "grating" diffraction by surface acoustic waves. (SAW propagation sense $s = +1$.) (b) Instantaneous profile of the interdigitated electrode structure showing the applied voltage and the deformation of the piezoelectric surface.

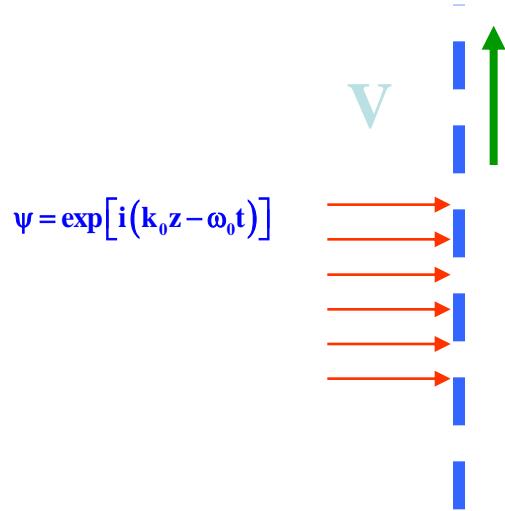
Matter wave diffraction by a moving grating. Formulation of the problem



V.G. Nosov and A.I.Frank , 1991



Neutron diffraction by a moving grating. Formulation of the problem

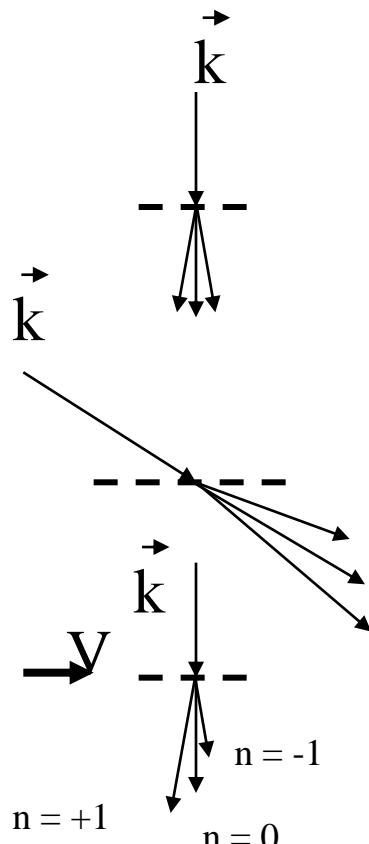


$$\psi = \frac{1}{2} \exp[i(kx - \omega t)] + \frac{1}{\pi} \sum_{s=-\infty}^{\infty} \frac{1}{2s-1} \exp[i(k_s x - \omega_s t)]$$

$$\omega_s = \omega + \frac{\pi(2s-1)}{T} \quad k_s = \left(\frac{2m\omega_s}{\hbar} \right)^{1/2}$$

Elementary (kinematic) theory.

A.Frank, V.Nosov, 1994



1. Solving the diffraction problem in a moving system.
2. Galilean transformation of the wave function.

$$\Psi(z, y, t) = \sum_j a_j \exp[i(\mathbf{k}_j z + \mathbf{q}_j y - \omega_j t)] \quad (k_0 L \ll 1)$$

$$a_j = \frac{1}{d} \int_0^L T(x) \exp(-iq_j x) dx \quad \mathbf{q}_j = \mathbf{j} \cdot \left(\frac{2\pi}{d} \right) = \mathbf{j} \mathbf{q}_0$$

$$\omega_j = \omega_0 + j\Omega \quad \mathbf{k}_j \cong \mathbf{k}_0 \left(1 + j \frac{\Omega}{\omega_0} \right)^{\frac{1}{2}} \quad j = 0, \pm 1, \pm 2, \dots$$

$$\Omega = \frac{2\pi}{T} = 2\pi f = 2\pi \left(\frac{\mathbf{v}}{\mathbf{d}} \right) \quad d - \text{space period of a grating}$$

Main peculiarity of the kinematic approach

1. Neglecting by the grating form (1D grating)
2. Amplitudes of the diffraction orders don't depend on the grating velocity

For the phase π - modulation

$$T(x) = \begin{cases} 1 & \text{if } 0 < y < d/2 \\ \exp(i\pi) & \text{if } d/2 < y < d \end{cases}$$

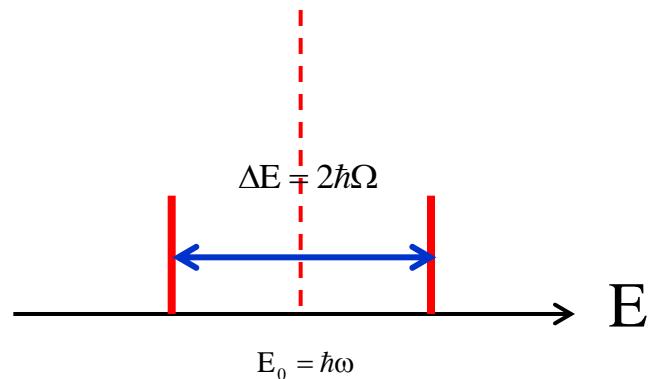
$$a_j = 0, \quad j = 2s$$

For zero and even orders

$$a_j = \frac{2}{i\pi j}, \quad j = 2s - 1 \quad \text{For odd orders}$$

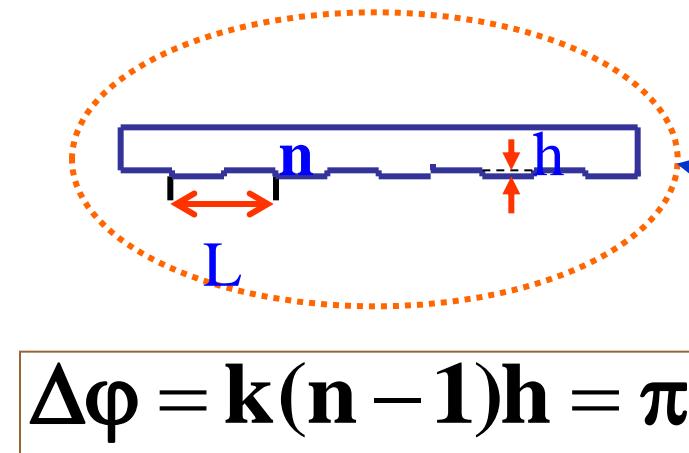
$$|a_{\pm 1}|^2 = \frac{4}{\pi^2} = 0.405$$

$$\Omega = 2\pi \left(\frac{V}{d} \right)$$

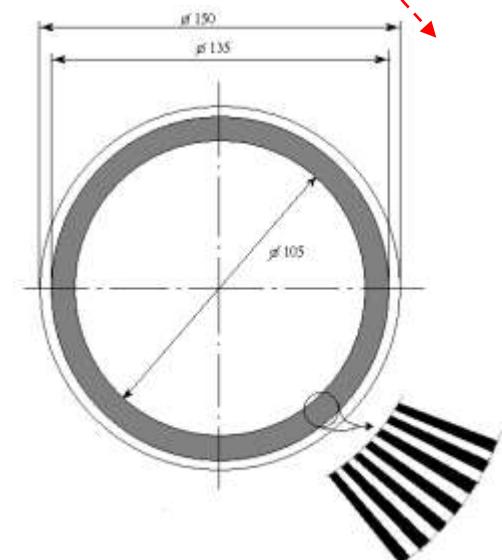
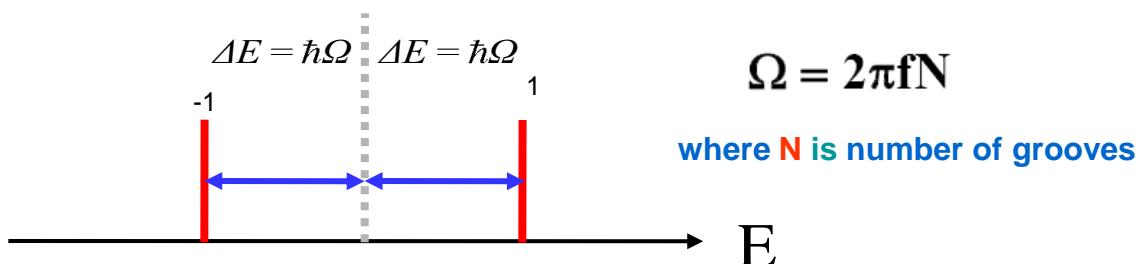
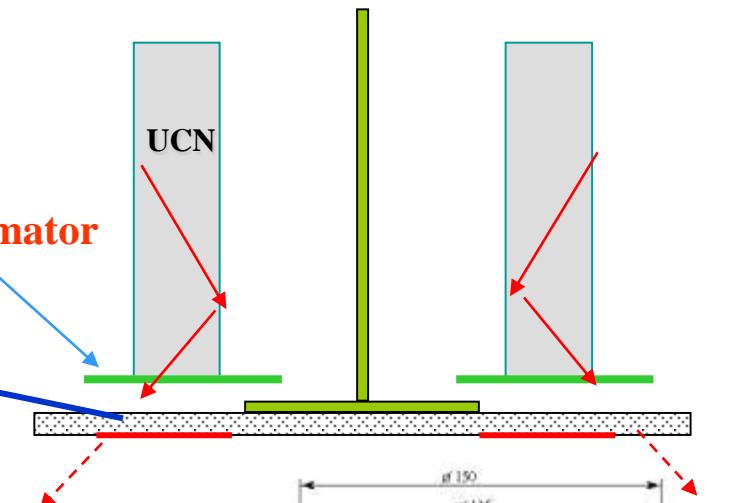


Experimental realization - rotating grating

Phase π -grating



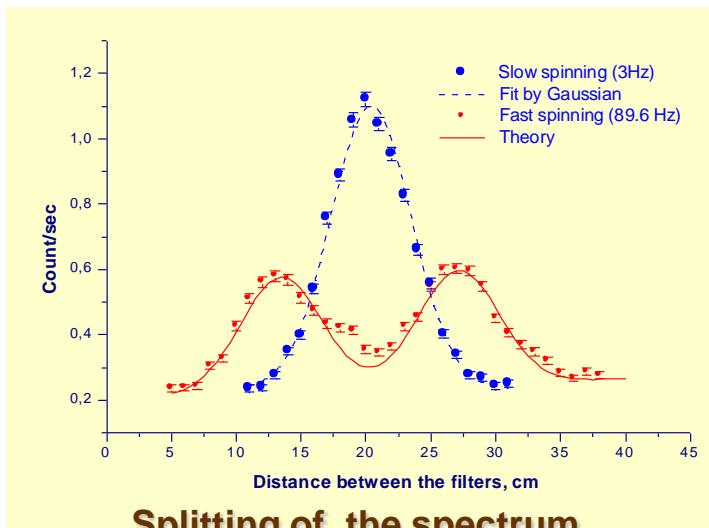
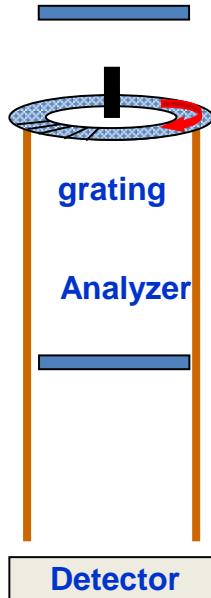
Monochromator



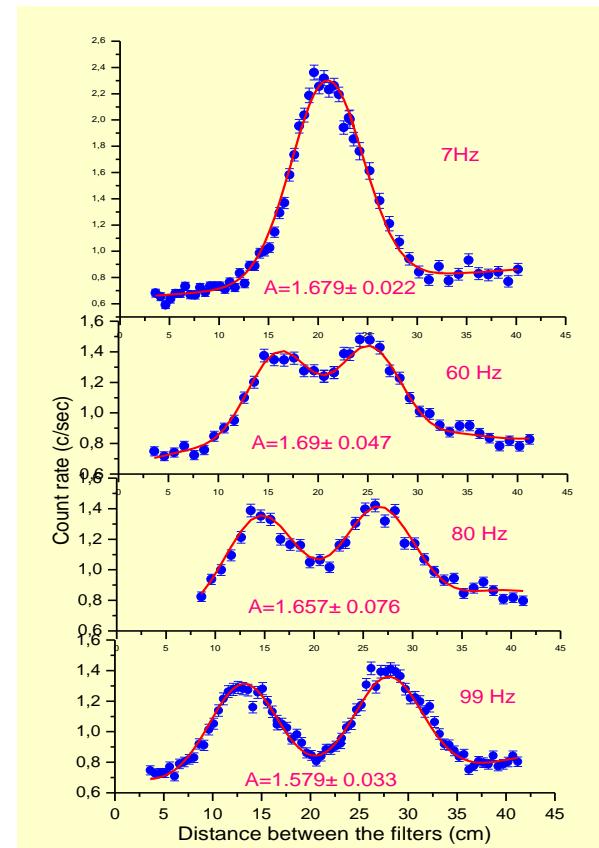
First experimental results

Angular period of grating 0.3325mrad (20μ at the middle diameter)

Monochromator



A.I.Frank et al. ILL annual report 2001
Phys.Lett.A 311 (2003) 6



$$|a_1|_{\text{th}}^2 = 0.405$$

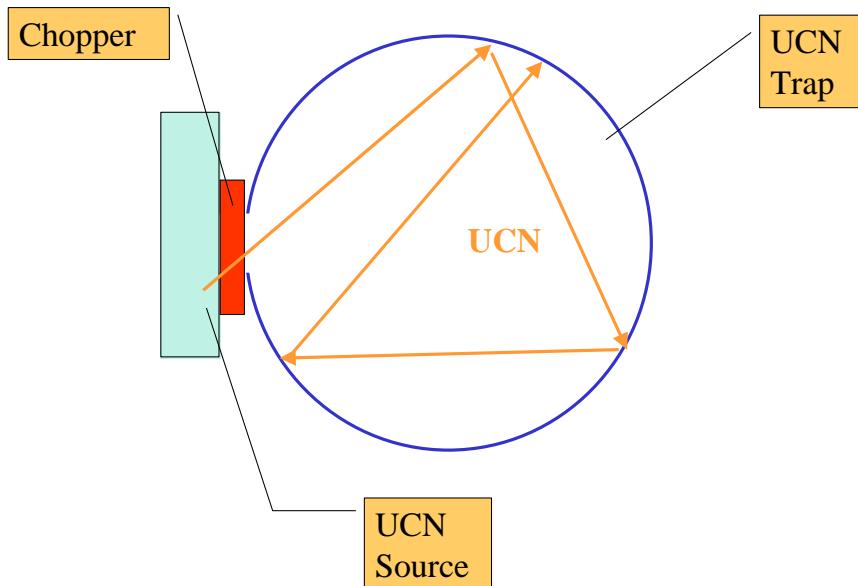
$$|a_1|_{\text{exp}}^2 = 0.383(8)$$

A.I.Frank et al. JETP Lett, 81 (2005) 427

Pulse source and UCN pumping in a trap



1915-1973

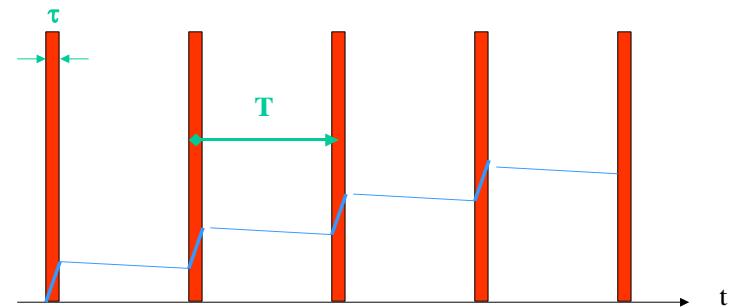


$$G = 1 + \frac{1 - \frac{\tau_1}{T}}{\frac{\tau_1}{T} + \frac{\Sigma \mu}{S}}$$

$G \rightarrow 10^2 \div 10^3$

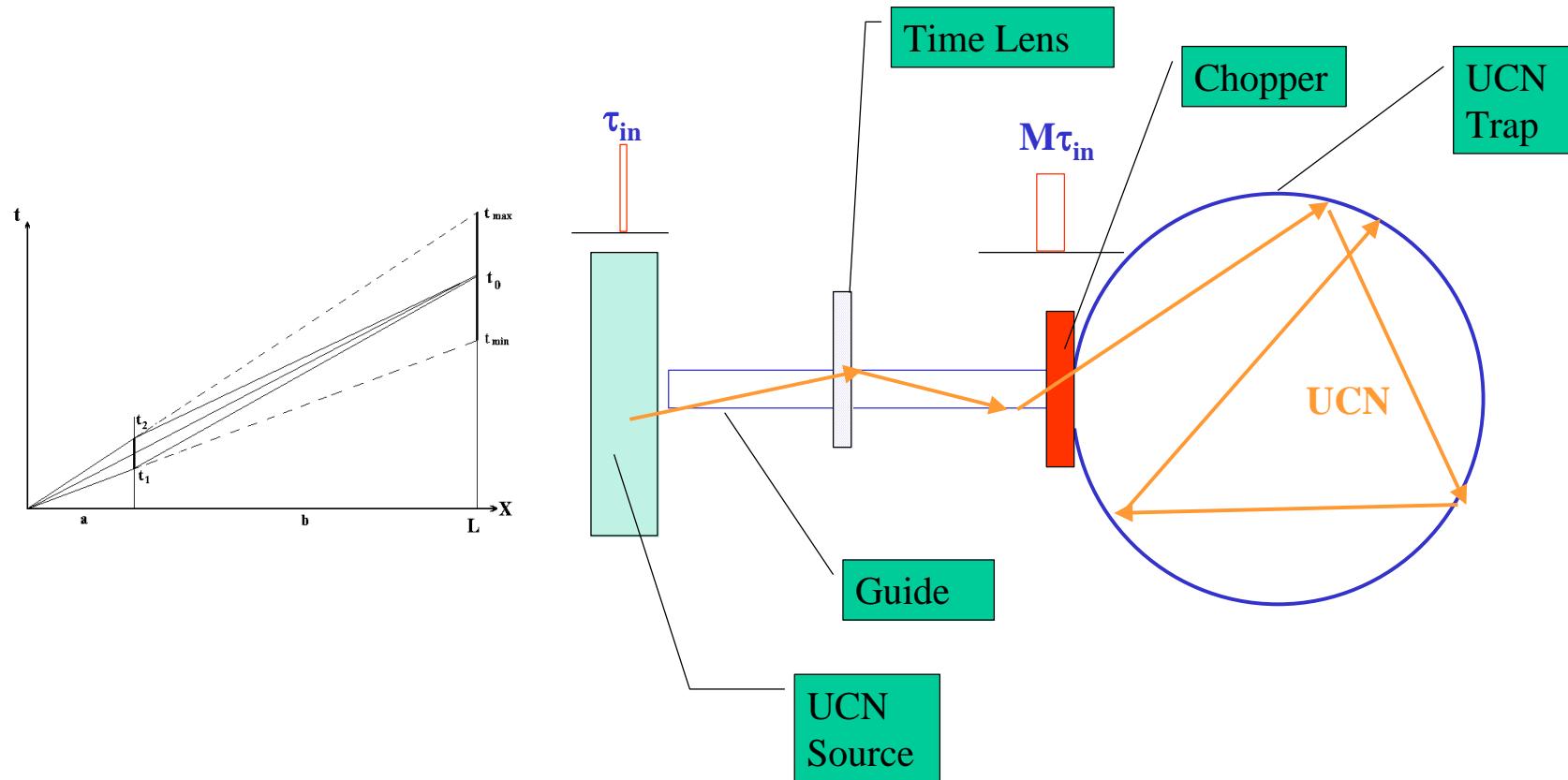
$\tau_1 > \tau$ - chopper opening time
S – active convertor area
 Σ -- area of the trap
 μ - probability of the UCN lost

F.L. Shapiro, 1972

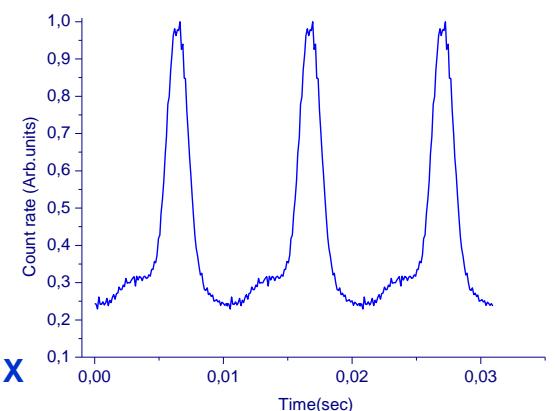
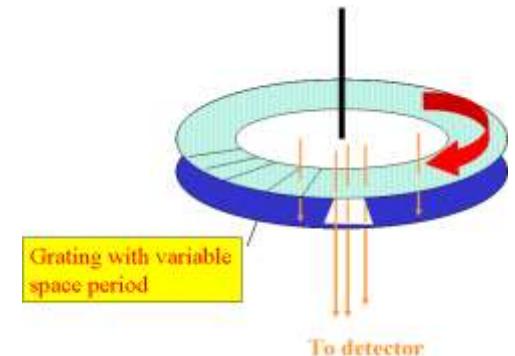
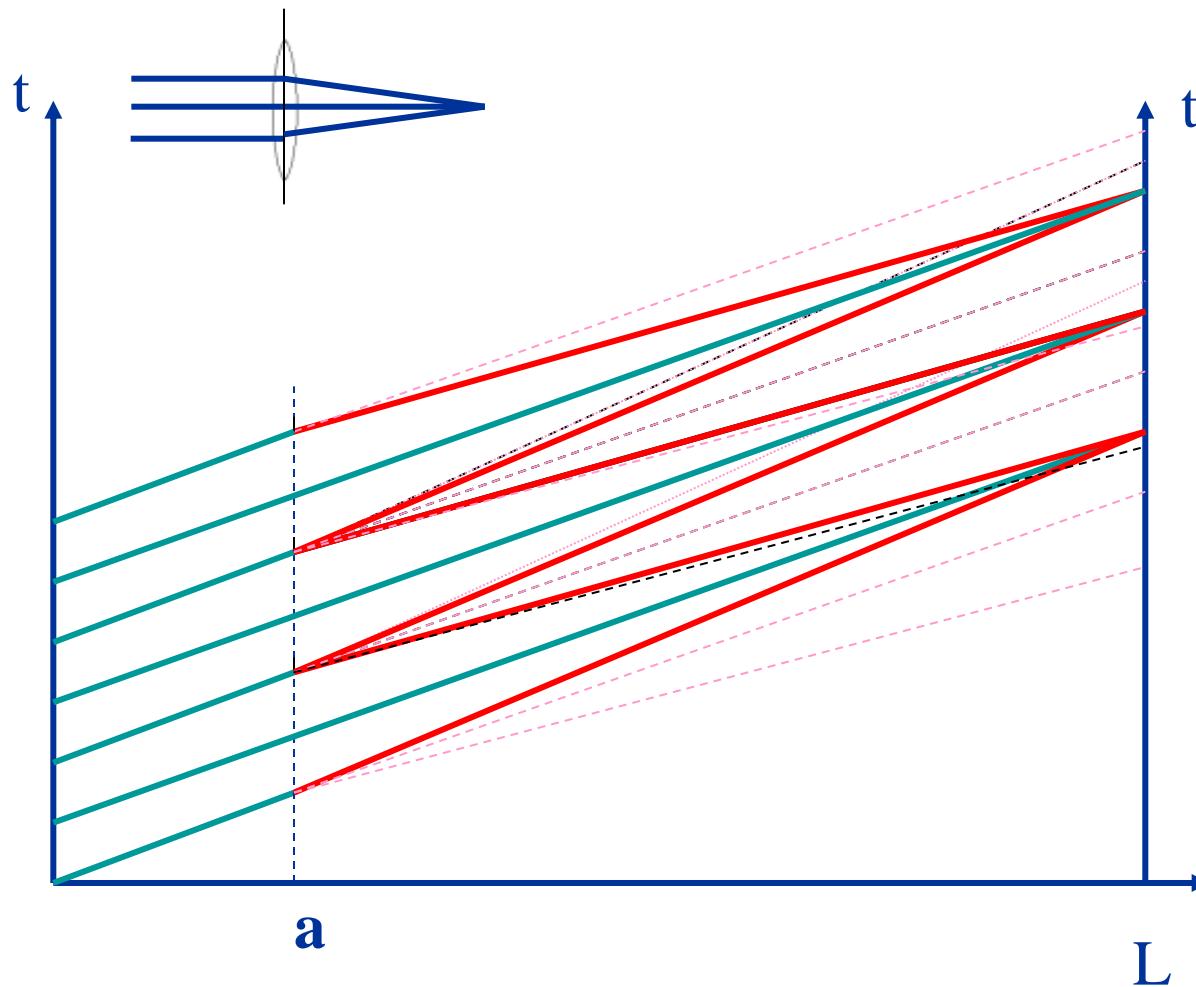


Pumping option of the pulsed source – time lens

A.I.Frank and R.Gähler. Phys. of Atomic Nuclei, 63, (2000) 545

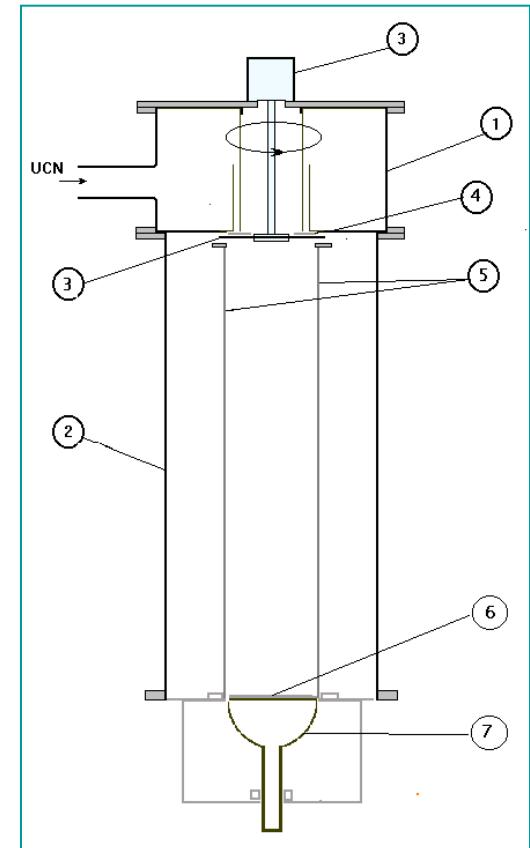
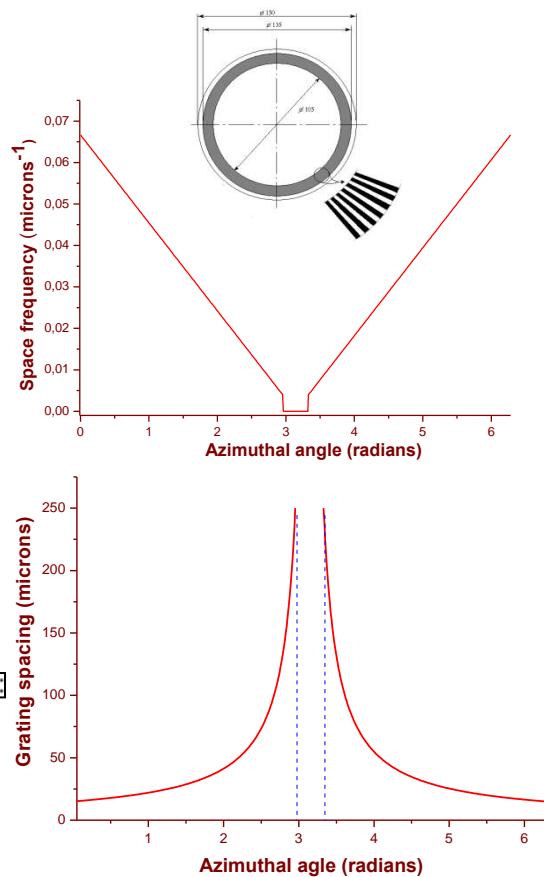
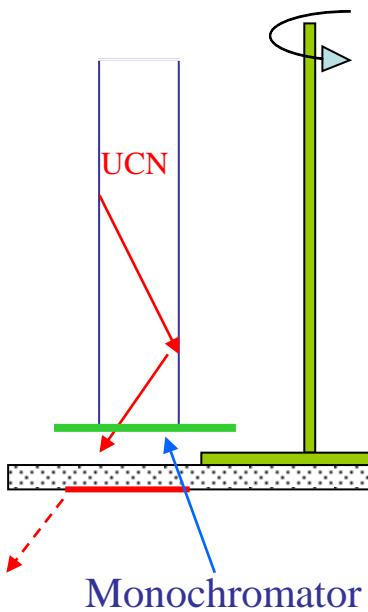


Neutron focusing in time

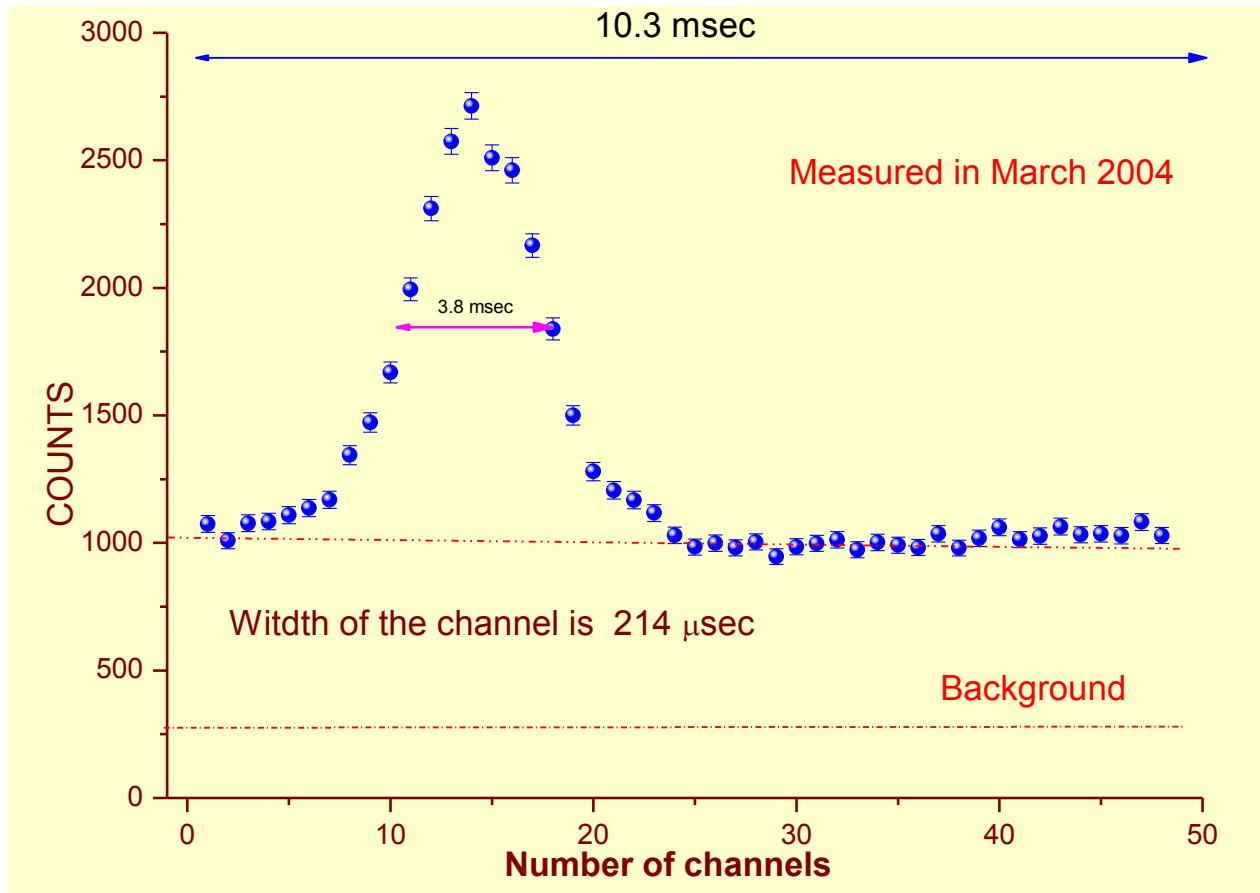


A.I.Frank and R.Gähler. Phys. of Atomic Nuclei, 63, (2000) 545

Rotating grating as a time lens

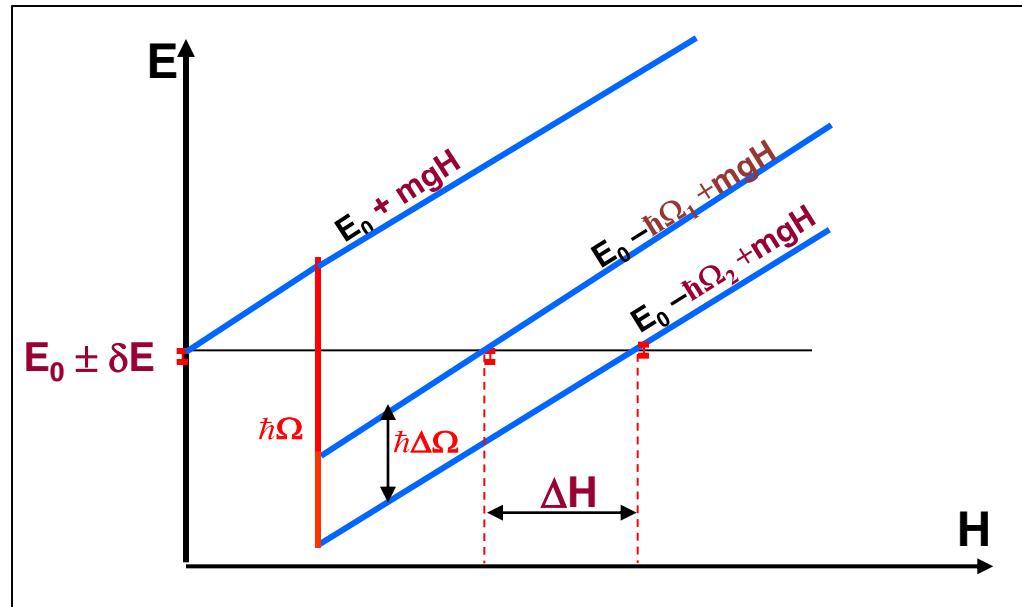


Time lens is working!



A. I. Frank, P. Geltenbort, G. V. Kulin et al. JETP Lett. 78, (2003) 188
S.N. Balashov, I.V. Bondarenko, A.I. Frank et al, Physica B, 350 (2004) 246

Test of the weak equivalence principle for neutrons (2006)



The idea was to compare the change of energy mgH with energy $\hbar\Omega$ transferred to neutron by a moving grating

Frank A.I., Masalovich S.V., Nosov V.G. (ISINN-12). E3-2004-169, 215, Dubna, (2004)

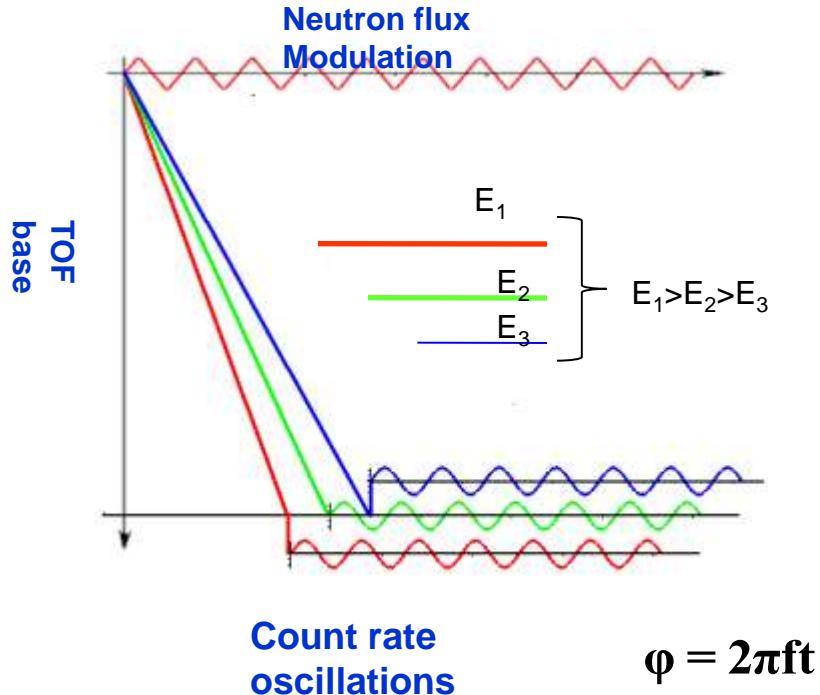
$$m_g g_n = \hbar \frac{\Delta\Omega}{\Delta H}$$

$$1 - \frac{m_g g_n}{m_n g_{loc}} = (1.8 \pm 2.1) \cdot 10^{-3}$$

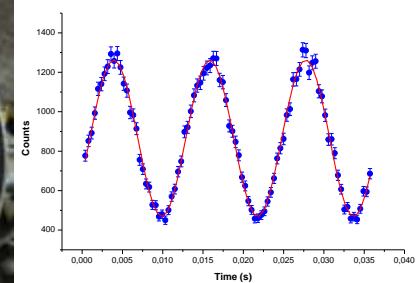
A.I. Frank, P. Geltenbort, M. Jentschel, et al. JETP Letters, 86, 225 (2007)

New experiment for the test of weak equivalence principle for neutron (started in 2010)

Comparing the energy $m_g g_n H$ with energy $\hbar\Omega$ as before.



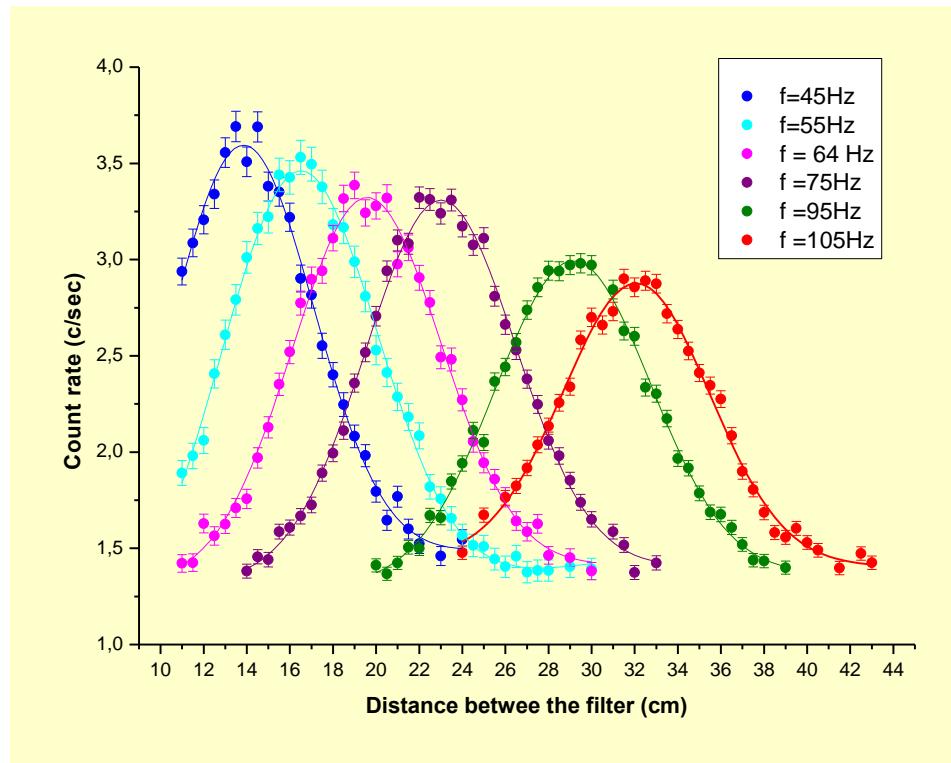
Combination of Neutron Interference
Filters with peculiar TOF
spectrometry



G.V. Kulin et al. NIM A 792 (2015) 38-46

Contradictions with kinematic theory (2006-2012)

1. Intensity of the first order DEPENDS on the velocity of the grating
2. The first order line is remarkably wider than initial spectrum
3. The presence of the zero diffraction order was detected

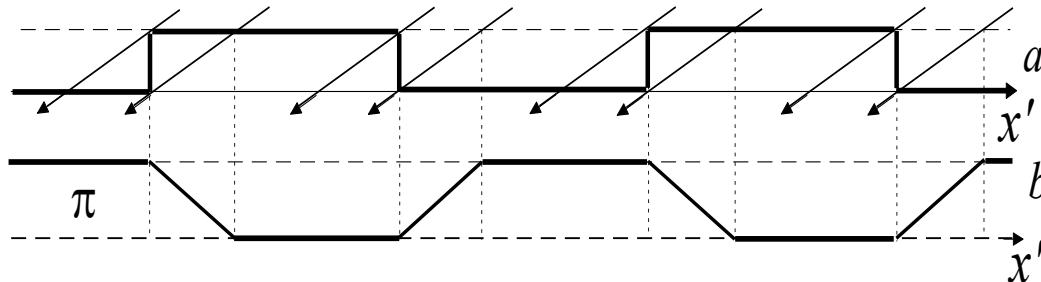


*Angular period of grating 0.0831mrad
(5 μ at the middle diameter)*

Scanning curves of the -1 diffraction order measured at different frequencies
of the grating rotation

*Next step in a theory was made in 2004
It was not taken into account at planning the experiment of 2010-12*

A.I. Frank et al. JINR Communication P3-2004-207



$$a_j = \frac{1}{d} \int_0^L T(x) \exp(-iq_j x) dx$$

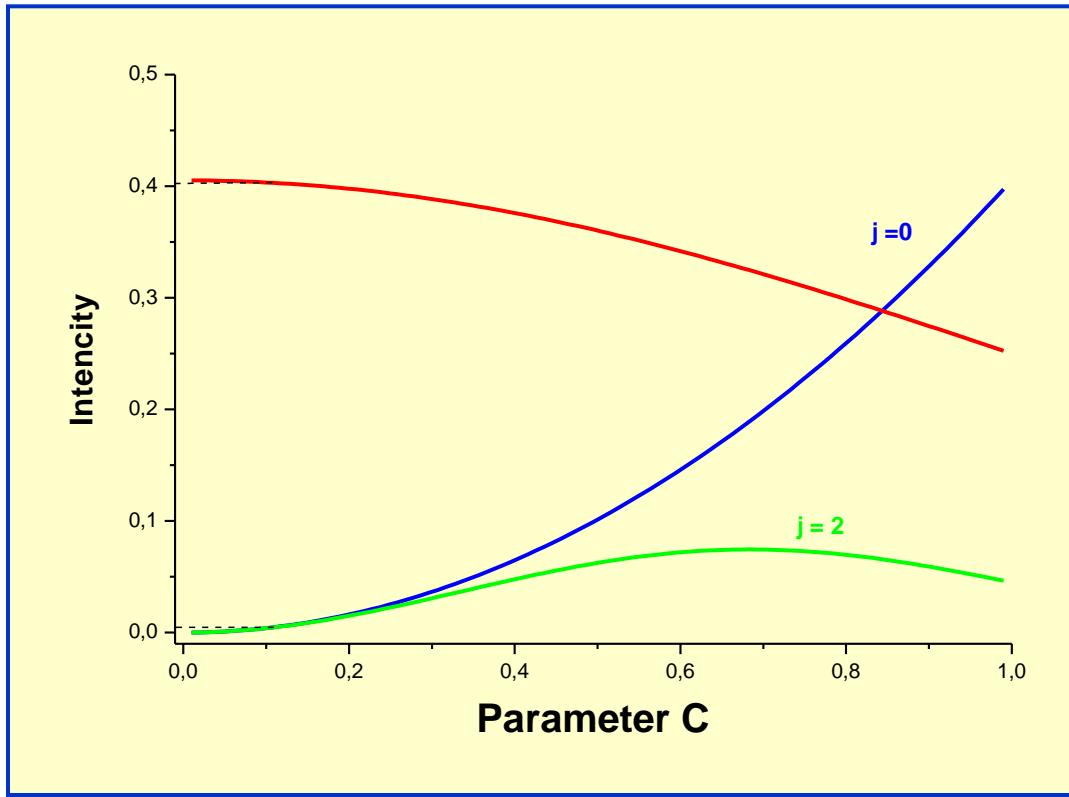
$$T(x) = \begin{cases} e^{i\frac{\pi}{cd}x} & \text{if } 0 < x < cd \\ e^{i\pi} & \text{if } cd < x < d \\ -e^{i\frac{\pi}{c}\left(1 - \frac{x}{d}\right)} & \text{if } d < x < d + cd \\ 1 & \text{if } d + cd < x < 2d \end{cases}$$

$$a_j = \begin{cases} cB_j & \text{if } j = 2s \\ -B/j & \text{if } j = 2s - 1 \end{cases}$$

$$B_j = \frac{1 + \exp(-i\pi j c)}{i\pi(1 - j^2 c^2)},$$

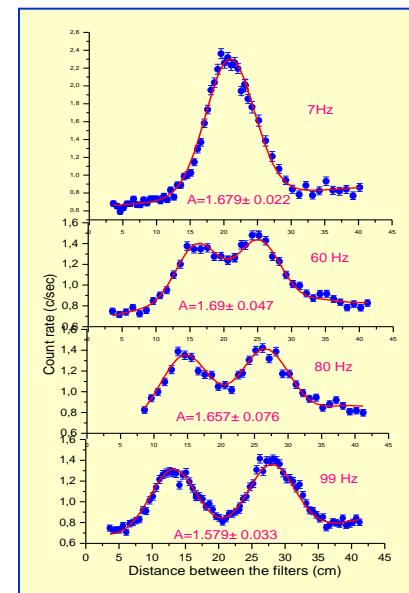
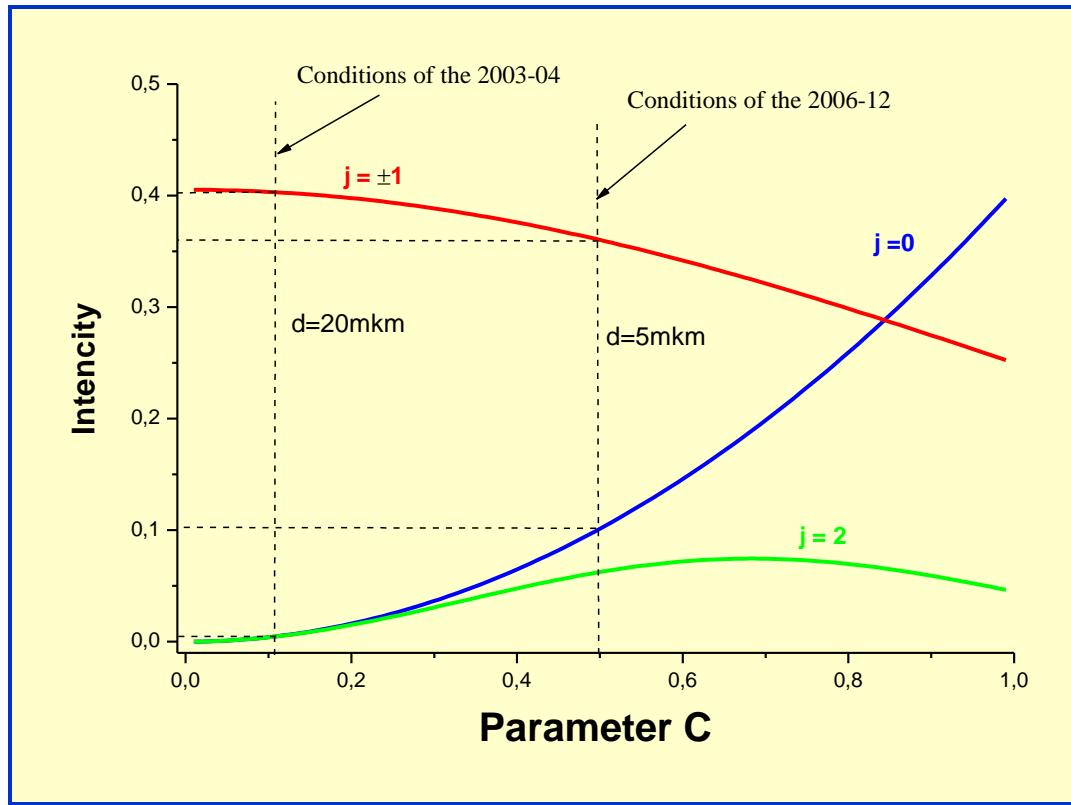
$$c = \frac{2hV}{dV_0}$$

Improved kinematic theory for the moving 3D grating



$$c = \frac{2hV}{dV_0}$$

Improved kinematic theory for the moving 3D grating



$$c = \frac{2hV}{dV_0}$$

1. Intensities of the diffraction waves depends on the grating velocity
2. All waves are independent
3. “Naive” combination of geometric and wave optics

New stage of the investigation of neutron diffraction by a moving grating.

1. Theory: Multiwave dynamic approach.

**V.A. Bushuev, A.I. Frank and V.G. Kulin. arXiv:1502.04751v1 [physics.optics]
(to be published in JETP)**

2. Experiment: TOF Fourier spectroscopy of UCN's

**Proposed in : V.G. Kulin, D.V. Kustov, A.I. Frank et al. JINR Communication,
P3-2014 -48**

G.Kulin (next presentation)

Multiwave dynamic theory

Moving reference

$$\Delta\psi(\mathbf{r}) + [k^2 - \chi(\mathbf{r})]\psi(\mathbf{r}) = 0$$

$$\chi(\mathbf{r}) = 4\pi N(\mathbf{r})b(\mathbf{r})$$

$$\chi(x) = \sum_{n=-\infty}^{\infty} \chi_n \exp(iq_n x)$$

$$(0 \leq z \leq h)$$

$$\chi_n = \frac{1}{d} \int_0^d \chi(x) \exp(-iq_n x) dx$$

$$q_n = n \cdot \left(\frac{2\pi}{d} \right) = nq_0$$

$$\Psi'(x', z) = \sum_{m=-\infty}^{\infty} \Psi_m(z) \exp[i(g_{mx}x' + g_{0z}z)]$$

Sum of the Bloch waves

$$g_{mx} = k_{0x} - k_v + q_m, \quad k_v = \frac{mV}{\hbar}, \quad g_{0z} = (k_{0z}^2 - \chi_0)^{1/2}$$

Multiwave dynamic theory

$$\frac{d^2\Psi_m}{dz^2} + 2iq_{0z} \frac{d\Psi_m}{dz} - \alpha_m \Psi_m - \sum_{n \neq 0} \chi_n \Psi_{m-n} = 0$$

where

$$\alpha_m = q_m [q_m - 2(k_v - k_{0x})]$$

Border conditions

$$\psi_0(z=0) = A_0, \quad \psi_{m \neq 0}(z=0) = 0$$

Multiwave dynamic theory

$$\cancel{\frac{d^2\Psi_m}{dz^2}} + 2iq_{0z} \frac{d\Psi_m}{dz} - \alpha_m \Psi_m - \sum_{n \neq 0} \chi_n \Psi_{m-n} = 0$$

$$|R_F|^2 \approx 1-2 \%$$

where

$$\alpha_m = q_m [q_m - 2(k_v - k_{0x})]$$

Border conditions

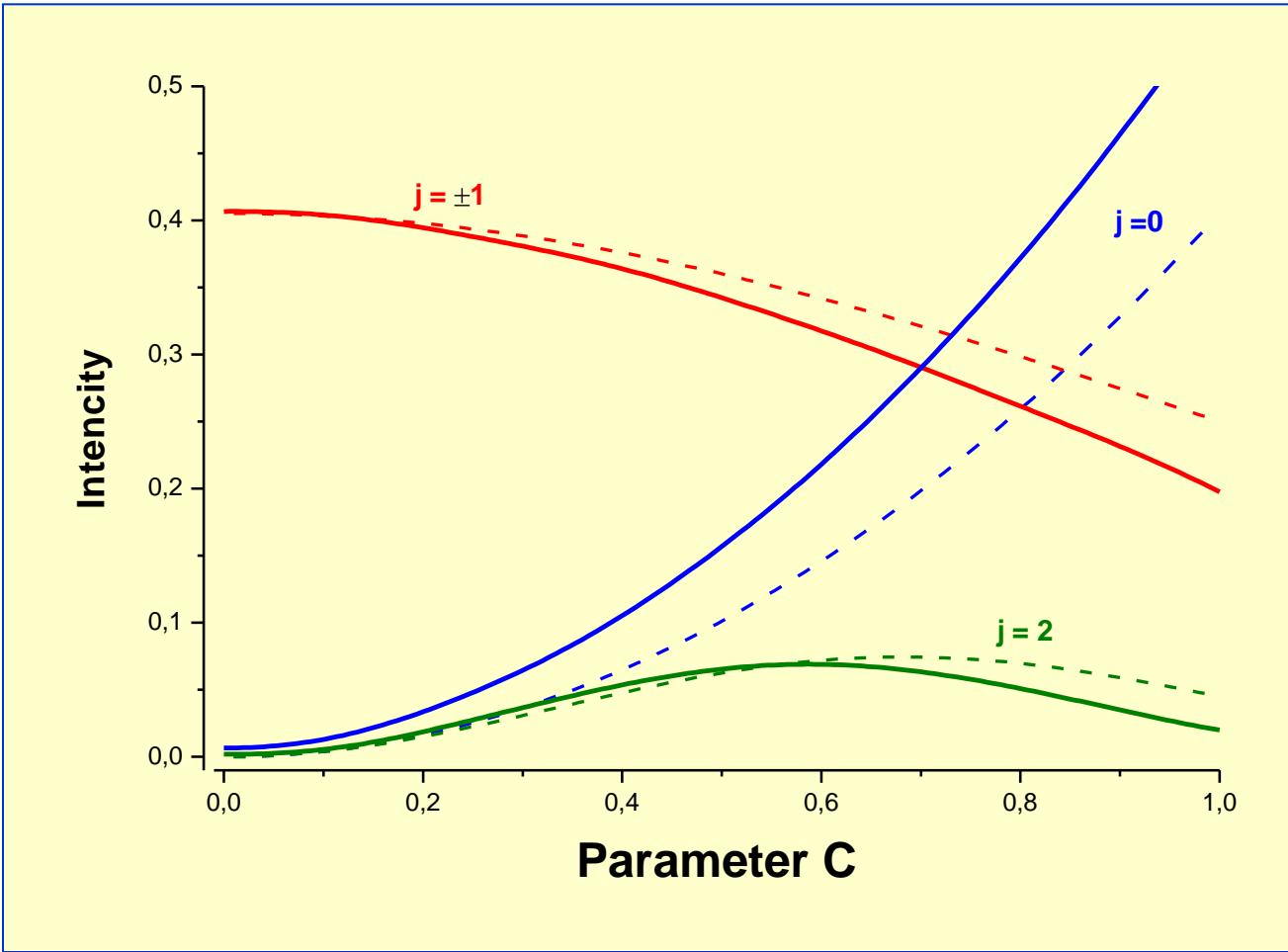
$$\psi_0(z=0) = A_0, \quad \psi_{m \neq 0}(z=0) = 0$$

$$\frac{d\Psi_m}{dz} = -i\gamma_m \Psi_m - i \sum_{n \neq 0} \beta_n \Psi_{m-n}$$

$$\gamma_m = \alpha_m / 2g_{0z}, \quad \beta_n = \chi_n / 2g_{0z}$$

$$q_n = n \cdot \left(\frac{2\pi}{d} \right) = nq_0$$

$$g_{0z} = (k_{0z}^2 - \chi_0)^{1/2}$$



$$c = \frac{2hV}{dV_0}$$

Solid lines –dynamic theory, dash lines - improved kinematics theory

TOF Fourier mode of the UCN spectrometer (November 2014)



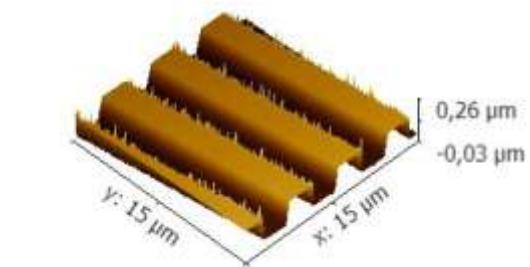
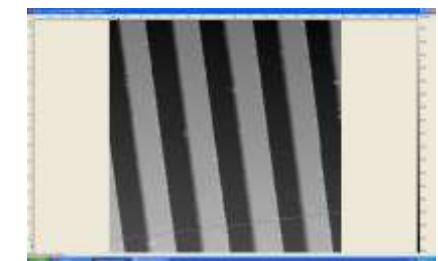
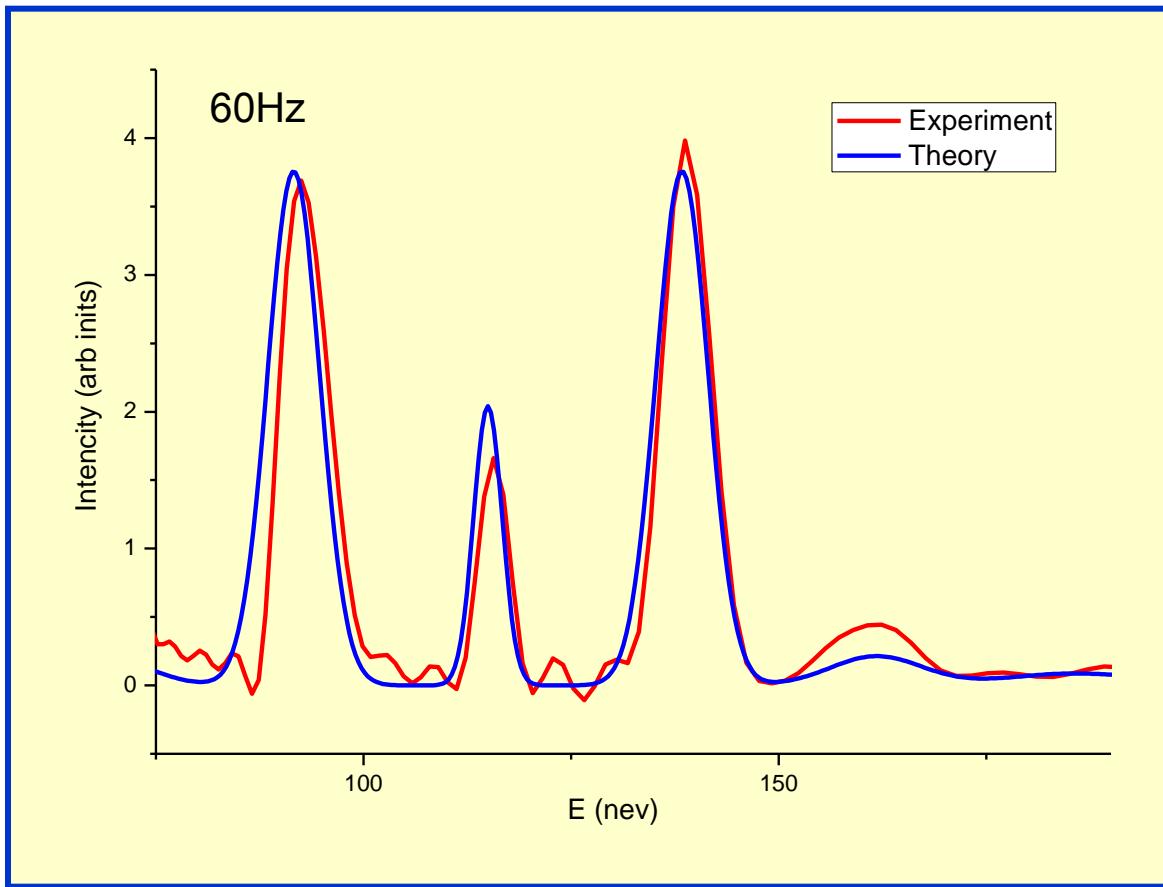
I(t) – time of flight spectra

$$I(t) = \int_0^{\infty} R(\omega) \sin[\omega t - \phi(\omega)] d\omega \text{ - Fourier expand}$$

$$Z(t) = \int_0^{\infty} I(t') \sin[\omega(t-t')] dt' \text{ Fourier harmonic (measured)}$$

G.Kulin (next presentation)

Recent results: TOF spectra and comparing with theory (December 2014)



*Non-destroyed
Atomic Force Microscope Inspection*

Angular period of grating 0.0665 mrad (4μ at the middle diameter)

Conclusion

1. *Diffraction grating which is moving across the direction of neutron propagation acts as a nonstationary device. That results in the appearance of line spectrum.*
2. *The phenomenon was firstly observed in 2001*
3. *It may be (and was) used for precise change of the neutron energy. That was demonstrated in the experiment for neutron focusing in time and gravity experiments with UCN.*
4. *Dynamical theory of UCN diffraction by a grating satisfactorily describes experimental data but the precision of these last is rather pure yet.*
5. *We plan to continue experimental investigations of the phenomena using recently tested TOF Fourier method of UCN spectrometry*

Many thanks to my co-authors



I.V. Bondarenko, S.V.Goryunov, G.V.Kulin, D.Kustov.



S.N. Balashov, S.V.Masalovich, V.G.Nosov S.N.Strepetov



V. A. Bushuev



P. Geltenbort, P. Høghøj, M. Jentschel

Thank you for your attention!