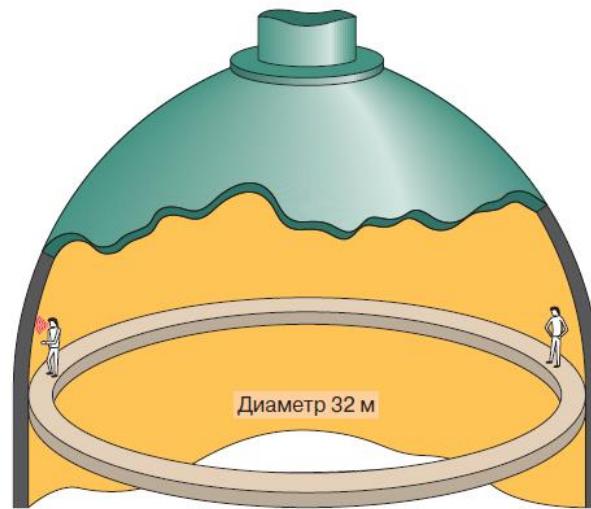


Neutron whispering gallery resonator

ISINN- 23 Dubna 2015



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Important references

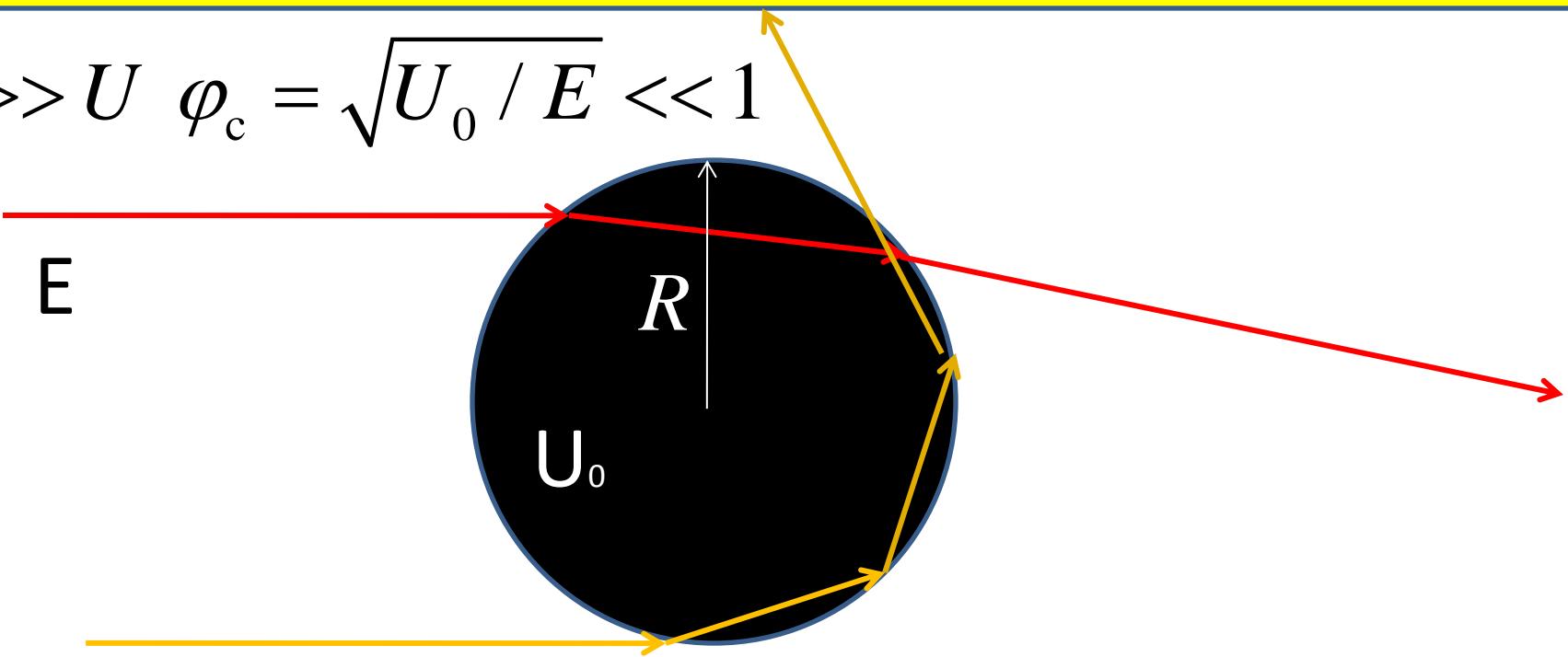
- V.V. Nesvizhevsky, A. Voronin, R.Cubitt and K.V. Protasov (2010) “**Neutron whispering gallery**” **Nature Physics** 6:114-117
- R. Cubitt, V.V. Nesvizhevsky, A.K. Petukhov, A.Yu. Voronin, G. Pignol, K.V. Protasov, and P. Gurshijants (2009). “**Methods of observation of the centrifugal quantum states of neutrons.**” **NIM A** 611: 322-325
- V. V. Nesvizhevsky, A. K. Petukhov, K. V. Protasov, and A. Yu. Voronin “**Centrifugal quantum states of neutrons**” **Phys. Rev. A** 78, 033616 (2008)
- K. V. Protasov and A. Y. Voronin
“**Regge Poles in Neutron Scattering by a Cylinder**”
Adv. High En. Physics 124592, (2014)

Plan of the talk

- Whispering gallery effect in neutron scattering on cylinder
- Neutron whispering gallery resonator and trap
- Short-range forces and gravitational effect

Neutron scattering on a cylinder

$$E \gg U \quad \varphi_c = \sqrt{U_0 / E} \ll 1$$

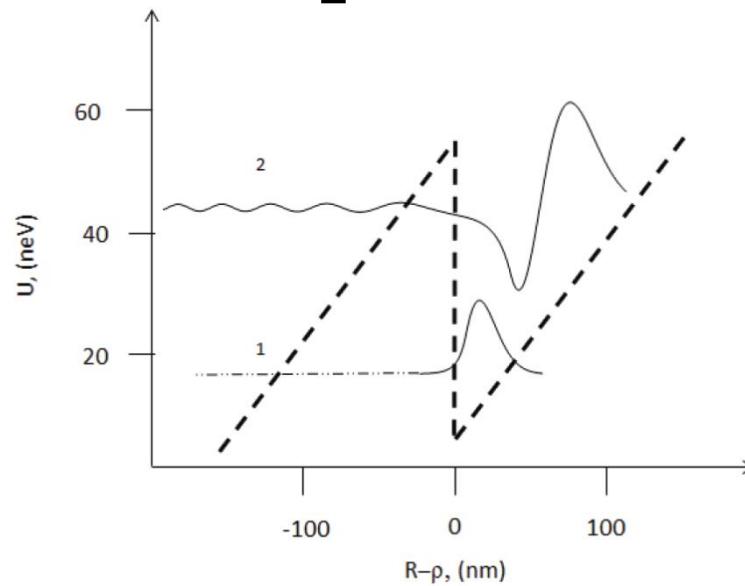


$$\varphi \gg \varphi_c \quad \frac{mv_{\perp}^2}{2} \approx U_0 \quad mvR / \hbar \gg 1$$

Whispering gallery states

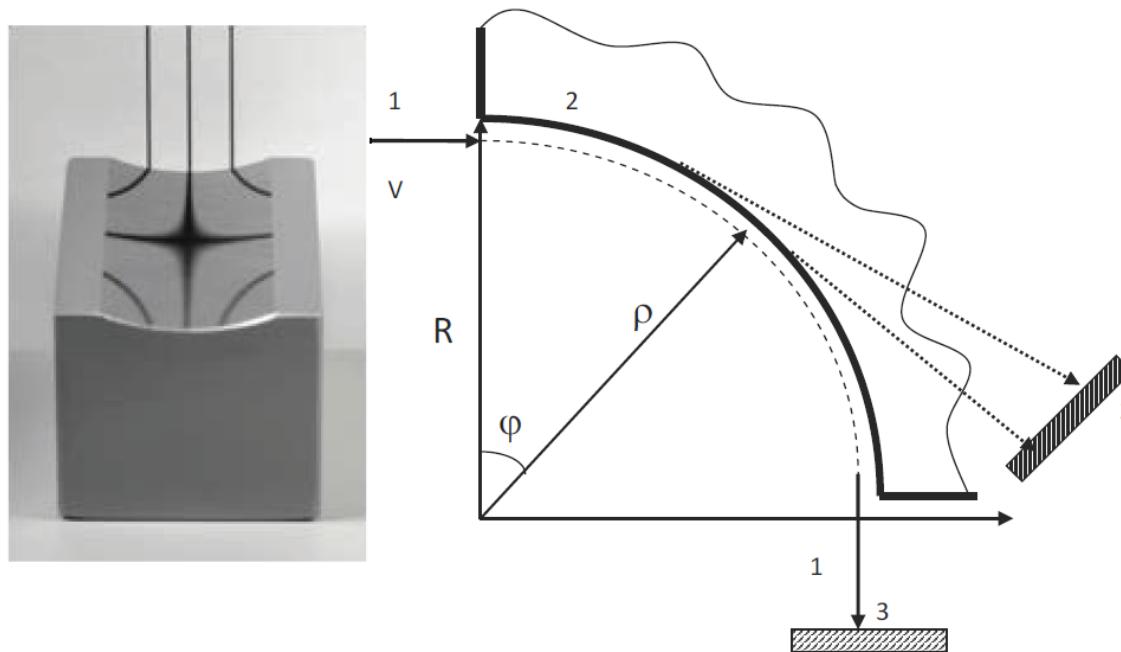
$$\left[-\frac{1}{2M} \frac{\partial^2}{\partial z^2} - U_0 \Theta(z) - \frac{Mv^2}{R} z - \epsilon_\mu \right] \chi_\mu(z) = 0$$

$$\epsilon_\mu = E - \frac{(\mu^2 - 1/4)}{MR^2}$$



$$l_0 = \left(\frac{\hbar^2 R}{2M^2 v^2} \right)^{1/3} \approx 40 \text{ nm} \quad \epsilon_0 = \left(\frac{\hbar^2 M v^4}{2R^2} \right)^{1/3} \approx 15 \text{ neV} \quad g_{eff} = \frac{v^2}{R} \sim 10^7 g$$

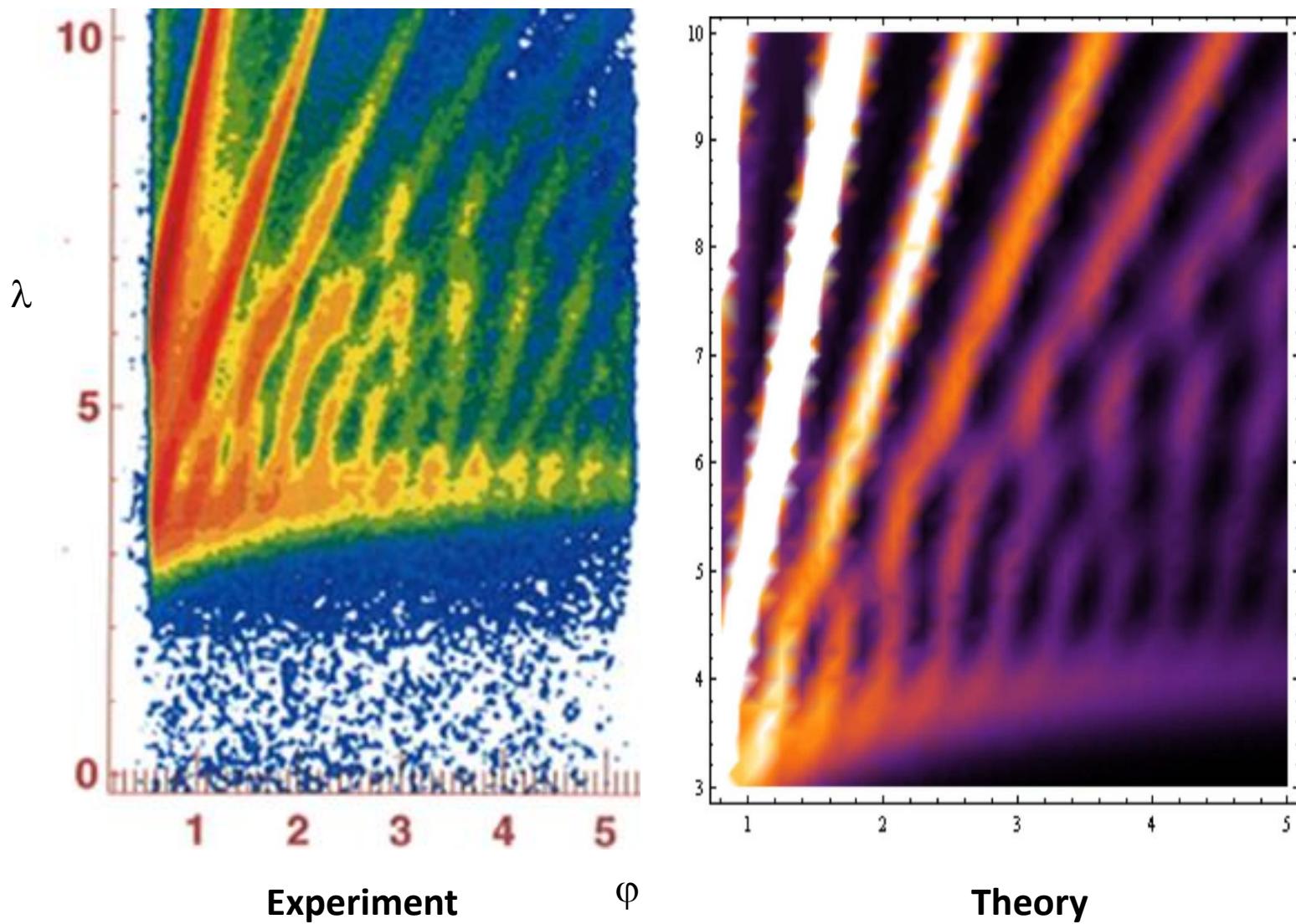
First observation of neutron WG effect



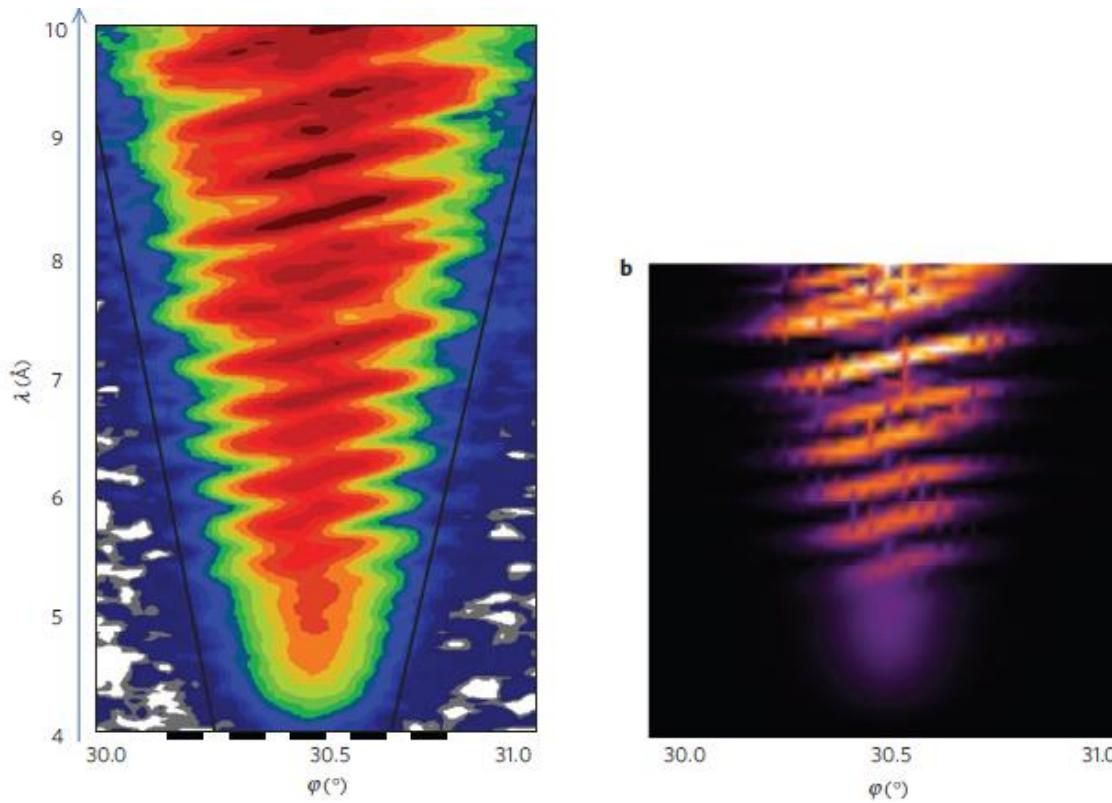
1- neutron classical trajectories, 2- cylindrical mirror, 3-detector

V.V. Nesvizhevsky, A. Voronin, R.Cubitt and K.V. Protasov (2010) “Neutron whispering gallery”
Nature Physics 6:114-117

Neutron rainbow



Interference of deep WG states

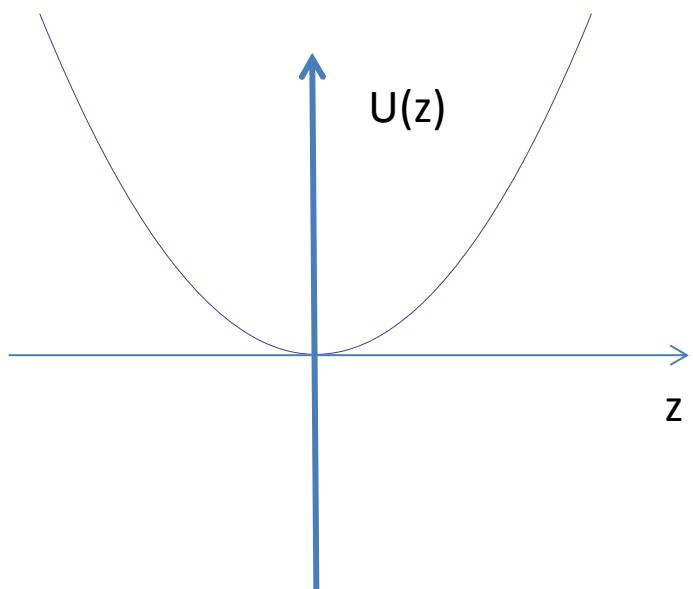
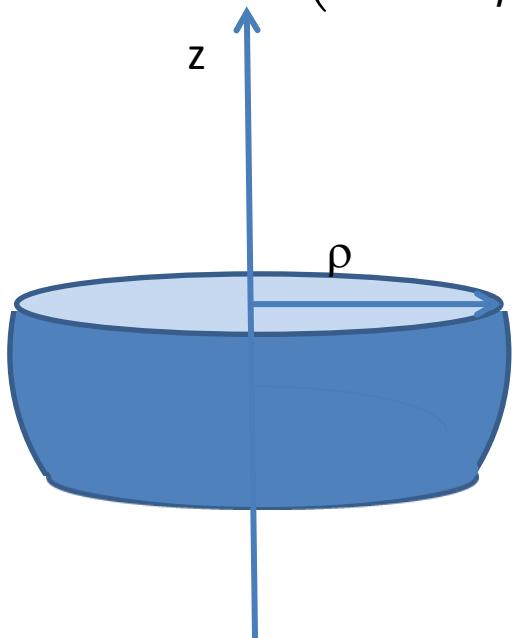


High sensitivity of interference pattern on details of neutron-surface interaction

WG resonator

$$\rho(z^2) = \sqrt{R^2 - z^2}$$

$$\left(\hat{T}_z + \frac{\hbar^2}{2m\rho(z)^2} \left(\mu^2 - \frac{1}{4} \right) + \varepsilon_n(z) - E \right) f(z) = 0$$



A 3D neutron trap is produced by superposition of centrifugal and optical potential

Spherical WG resonator

$$\left[\hat{T}_\rho + \hat{T}_\theta + \hat{T}_\phi - E \right] \chi(\rho) Y_{l\mu}(\theta) \exp(i\mu\varphi) = 0;$$

$$\left[-\frac{\hbar^2 d^2}{2mdx^2} - \frac{\hbar^2 l(l+1)}{mR^3} x - \varepsilon_n \right] \chi(x) = 0;$$

$$x = R - \rho; \quad \varepsilon_n = \frac{\hbar^2 l^{4/3}}{2^{1/3} m R^2} \lambda_n; \quad Ai(-\lambda_n) = 0;$$

radial

$$\left[-\frac{\hbar^2 d^2}{2mdz^2} - \frac{\hbar^2 \mu^2}{2mR^4} z^2 - \hbar\omega(n + \frac{1}{2}) \right] \psi(z) = 0;$$

vertical

$$z = R \sin(\theta); \quad |\theta - \pi/2| \ll 1; \quad \omega = \frac{\hbar\mu}{mR^2} \approx \frac{v_\tau}{R}$$

Spherical WG energies

$$\Psi(x, z, \varphi) = Ai(x - \lambda_n) H_k(z) \exp(i\mu\varphi);$$

$$E_{n,k,\mu} = \frac{\hbar^2 \mu^2}{2mR^2} + \frac{\hbar^2 \mu(k + \frac{1}{2})}{mR^2} + \frac{\hbar^2 (\mu + k)^{4/3}}{2^{1/3} m R^2} \lambda_n$$

$\mu \gg 1$ – WG modes

$k \sim 1; n \sim 1$ – full quantum regime

$k \gg 1; n \sim 1$ – radial quantum regime

$k \sim 1; n \gg 1$ – vertical quantum regime

Short-range forces and gravity effects

Gravity

$\Delta z_0 = gR^2 / v_\tau^2$ – shift of zero level (of order of mm for UCN)

$\Delta\omega(g)$ – account of z^4 terms

Short-range forces

$\Delta\varepsilon_n$ – account of hypothetical Yukawa-type terms $\exp(-\lambda x) / x$

Observation- interference of WG states

Conclusions

- WG states- quantum tool for studying fundamental neutron-matter interactions
- WG effect- long time localization in quantum state near material surface
- WG resonator- “both-sided” open 3D neutron trap for extra-long keeping in WG quantum state
- Promising tool for studying short-range forces and gravity