

# TOTAL KINETIC ENERGIES IN $^{232}\text{Th}(n,F)$ AND $^{238}\text{U}(n,F)$

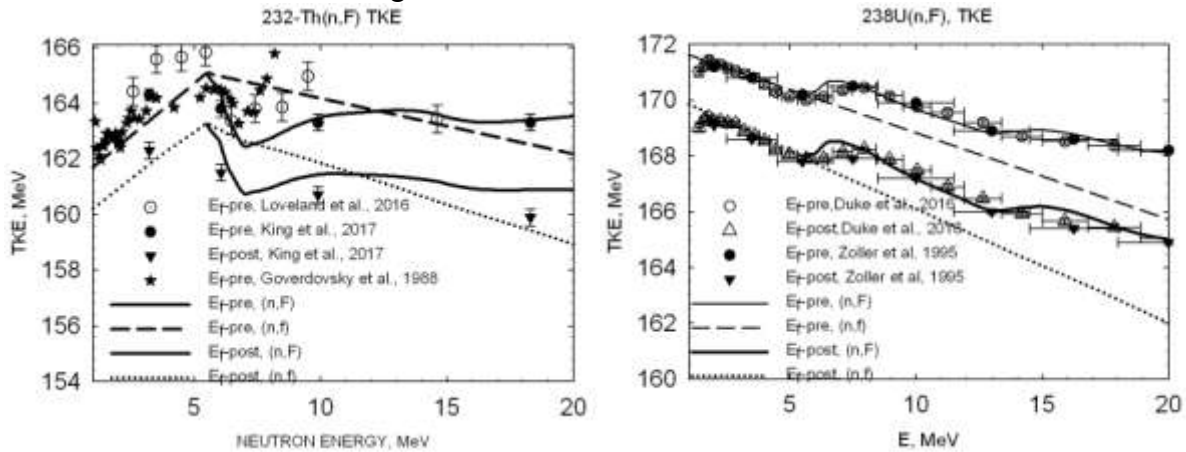
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Local minimum in TKE of fission fragments in  $^{232}\text{Th}(n,F)$  around  $^{232}\text{Th}(n,nf)$  threshold was first observed by Goverdovsky et al. in 1988. In 1995 Zoller et al. observed local maxima in TKE in  $^{238}\text{U}(n,F)$  reaction around  $^{238}\text{U}(n,xnf)$  thresholds. These variations are due to pre-fission (n,xnf) neutrons, which are pronounced in observed cross sections, prompt fission neutron spectra (PFNS) and mass distributions as well. Contribution of the (n,xnf) reaction to the  $\sigma_{n,F}$  of  $^{232}\text{Th}(n,F)$  around  $E_n \sim 7$  MeV is  $\sim 1.5$  higher than in case of  $^{238}\text{U}(n,F)$ , which is pronounced in TKE also. Partial contributions of (n,xnf) were fixed in [1–3] and seem to reproduce TKE variations. TKE values  $E_f^{pre}$  ( $E_f^{post}$ ) before (after) prompt neutron emission from fission fragments were calculated as

$$E_f^{pre}(E_n) = \sum_{x=0}^X E_{fx}^{pre}(E_{nx}) \cdot \sigma_{n,xnf} / \sigma_{n,F}, \quad E_{nx} = E_n + B_n - \sum_{x=0,1 \leq j \leq x}^X (\langle E_{n,xnf}^j \rangle + B_x).$$

TKE  $E_f^{post}$  were de-defined as  $E_f^{post} \approx E_f^{pre} (1 - \nu_{post} / (A - \nu_{pre}))$ ,  $\nu_p = \nu_{post} + \nu_{pre}$ . Components  $\nu_{post}$  and  $\nu_{pre}$  of  $\nu_p$  are defined via  $\nu_p$  and PFNS analysis at  $E_n \sim 2-20$  MeV. Assuming  $E_f^{pre}(E_n)$  for  $^{233-x}\text{Th}$ ,  $^{238-x}\text{U}$  are similar to those of  $^{233}\text{Th}$  and  $^{239}\text{U}$  (note TKE of  $^{232}\text{Th}(\gamma,F)$ ) we obtained TKE shown on the figure.



TKE shown on the figure is consistent with observed  $^{232}\text{Th}+n$  and  $^{238}\text{U}+n$  data on cross sections and PFNS. Straight lines are approximations of TKE values for non-emissive fission. The (n,xnf) neutrons influence on TKE values  $E_f^{pre}$  and  $E_f^{post}$  is much more pronounced in case of  $^{232}\text{Th}(n,F)$  reaction. That is due to the transition states structure of  $^{232}\text{Th}$  and competition of  $^{232}\text{Th}(n,n\gamma)$  and  $^{232}\text{Th}(n,nf)$  at  $E_n \leq 6.5$  MeV. For both  $^{232}\text{Th}(n,F)$  and  $^{238}\text{U}(n,F)$  (n,nf) neutrons influence on PFNS at  $E_n \sim 6.5$  MeV and  $E_n \sim 7$  MeV are quite different, which is due to the (n,2n) reaction neutrons competition to the (n,nf) reaction.

1. V.M. Maslov, Yu.V. Porodzinskij, M. Baba, et al., Phys. Rev., C69, 034607 (2004).
2. V.M. Maslov, Physics of Atomic Nuclei, vol. 71, No. 1, 9 (2008).
3. V.M. Maslov, Nucl. Phys. A743, 236 (2004).