

# **Some features of the data processing in the time-of-flight mass-spectrometry of heavy ions**

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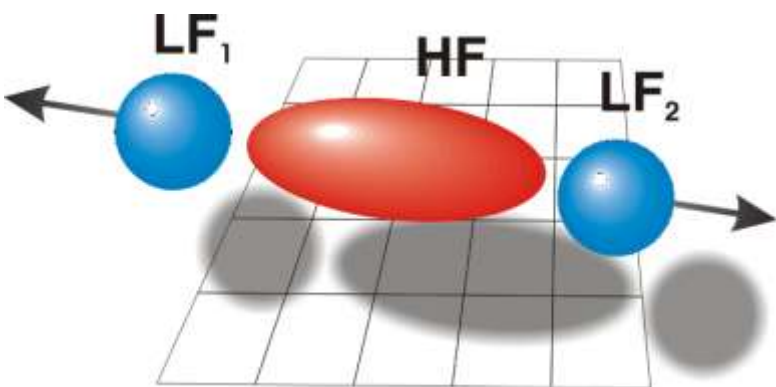
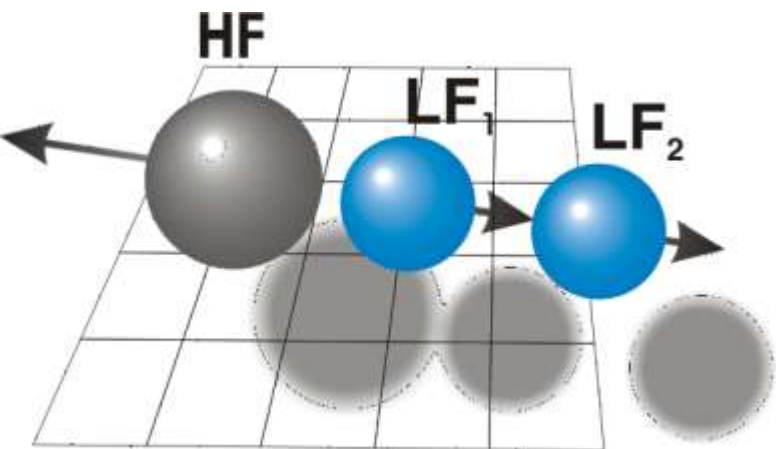
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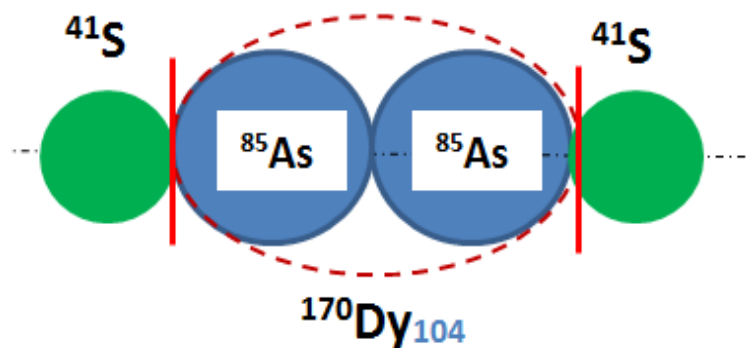
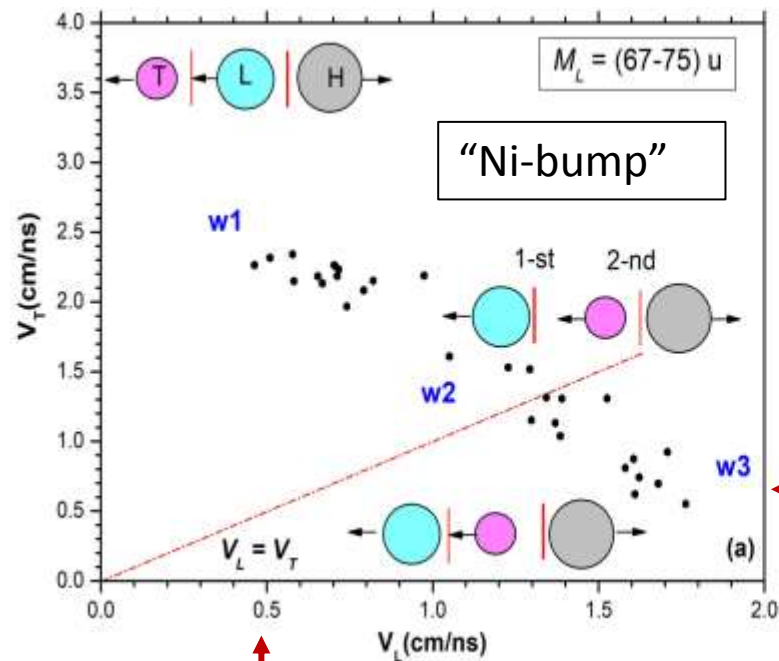
# **Setting of the physical problem**

# Collinear Cluster Tri-partition ( Multi-Cluster Decay )



Vide range of heavy ions masses, energies and velocities

Actually find



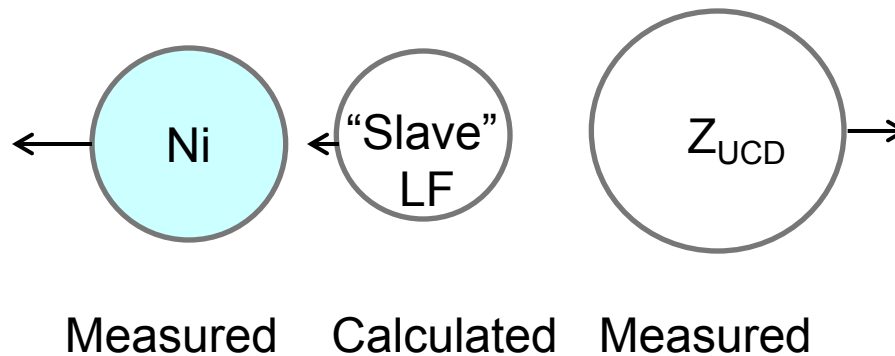
## Method of analysis of the missing mass data

$$Z(\text{HF}) = Z_{\text{UCD}}(M(\text{HF}))$$

$$M(\text{LF}) = 252 - M(\text{HF})$$

$$Z(\text{LF}) = 98 - Z(\text{LF})$$

$V(\text{LF})$  - momentum conservation law

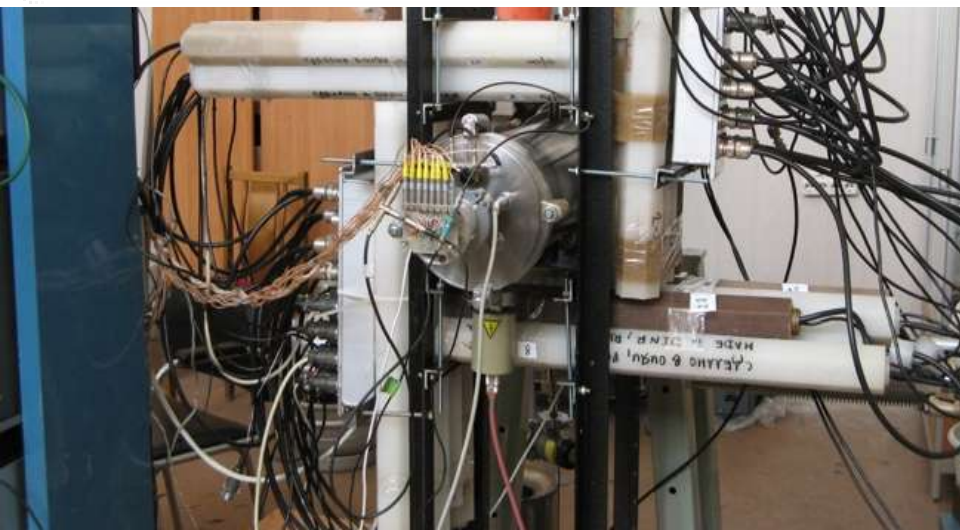
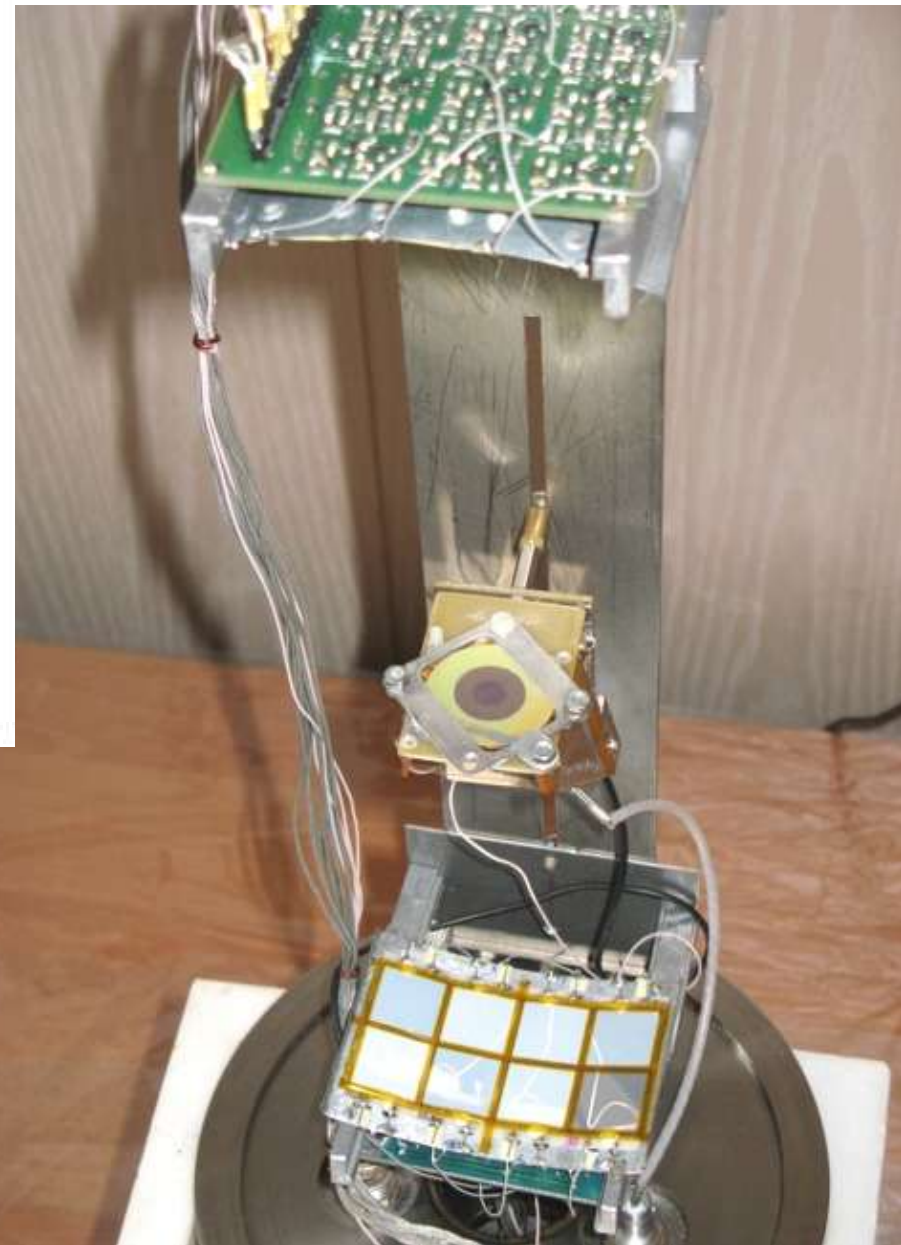
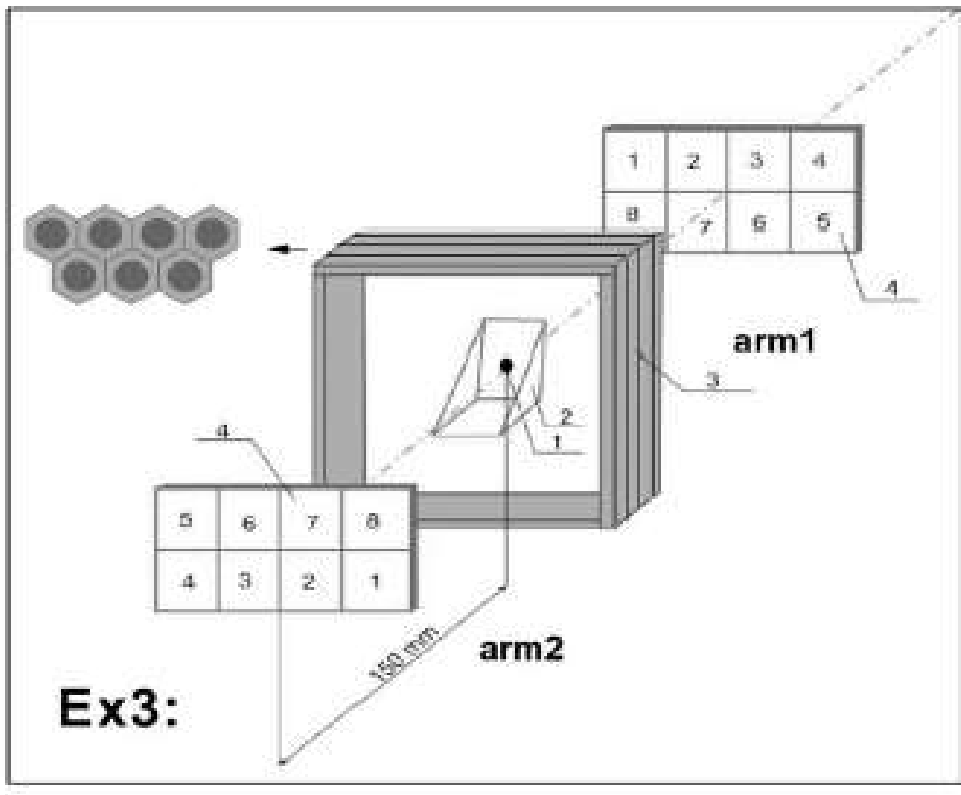


Comparing velocities with the fission ones  
one could put forward the ideas on the decay scenarios



# **Experimental setups**

# COMETA= COrrrelation Mosaic E-T Array





# COMETA-2 setup,

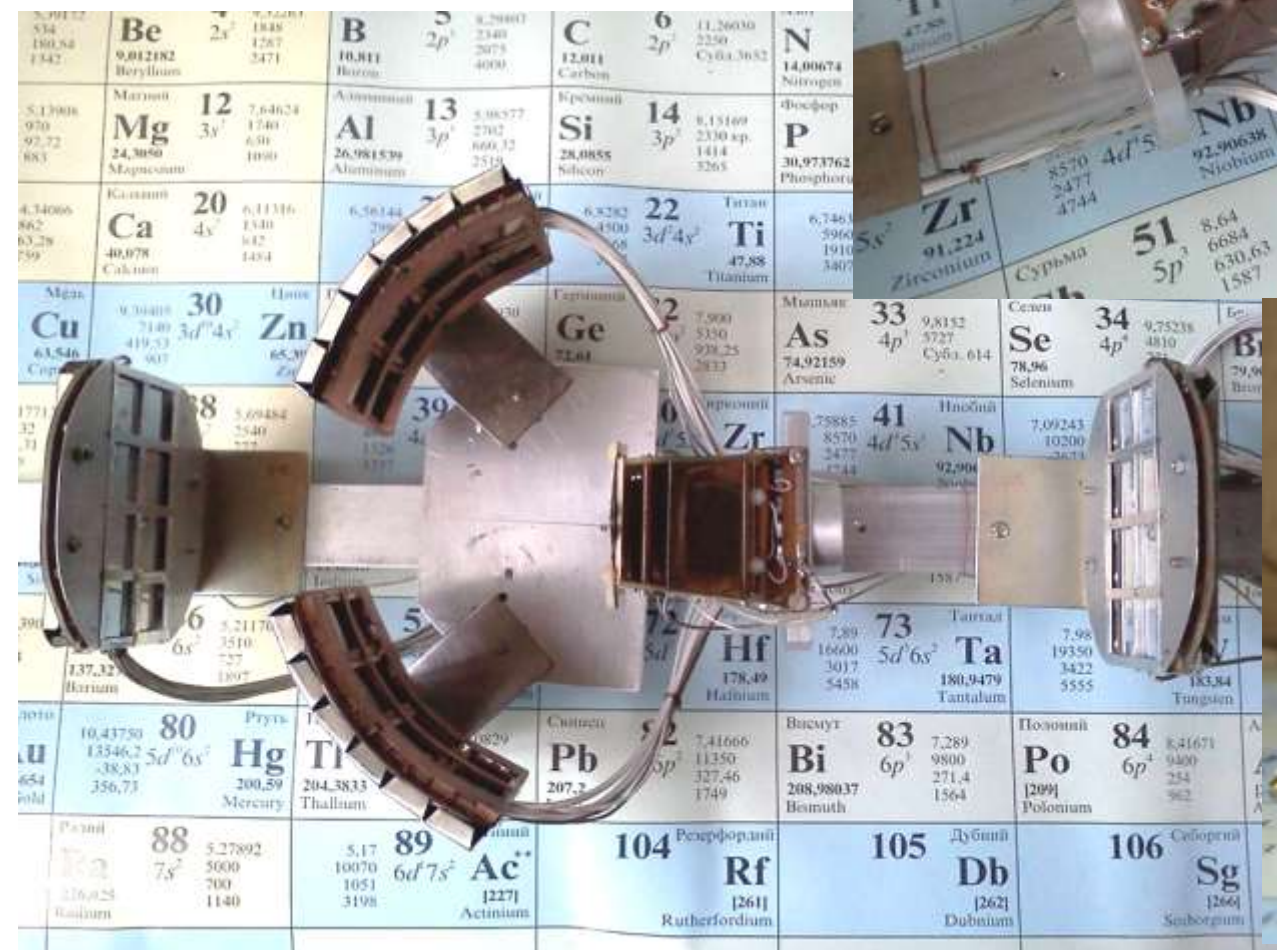
PIN-diodes 120° & MCP

neutron belt 28<sup>3</sup>He counters

Summer 2011



23/05/2012

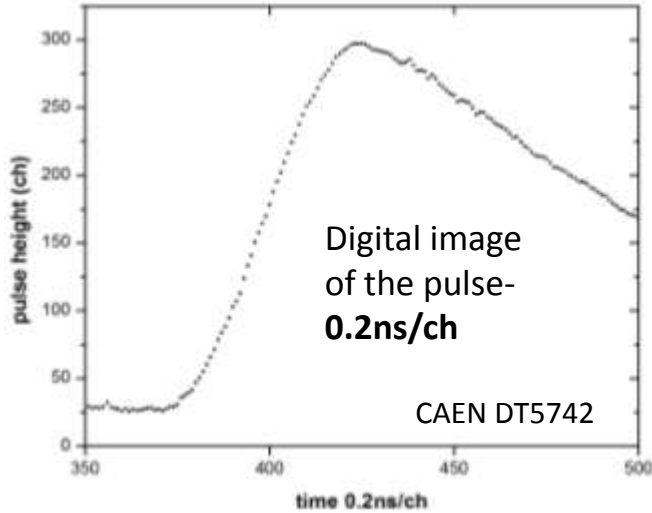




# **Improvement of the data processing algorithms**

**Feature №1**

# Our experimental approach

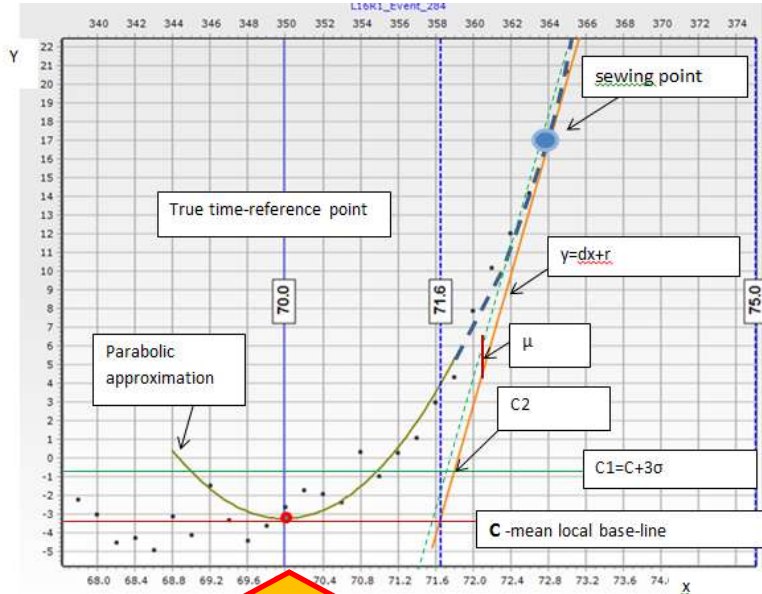


Digital image  
of the pulse-  
0.2ns/ch

CAEN DT5742

time 0.2ns/ch

L1OK1\_event\_204



True time reference point

PHD:

$$E = E_{det} + R(M, E), \quad (1)$$

$$R(M, E) = \frac{\lambda \cdot E}{1 + \varphi \cdot \frac{E}{M^2}} + \alpha \cdot ME + \beta \cdot E, \quad (2)$$

$$E = \frac{M \cdot V^2}{1.9297} \quad (3)$$

$$\longrightarrow G(\{\lambda, \varphi, \alpha, \beta\}, M, V) = 0$$

Combining equation (1), (2) and (3), we obtain:

$$G = \frac{MV^2}{k} - [E_{det} + \frac{\lambda \cdot \frac{MV^2}{k}}{1 + \varphi \cdot \frac{MV^2}{k}} + \alpha \cdot \frac{M^2 V^2}{k} + \beta \cdot \frac{MV^2}{k}] = 0,$$

where  $k = 1.9297$ .

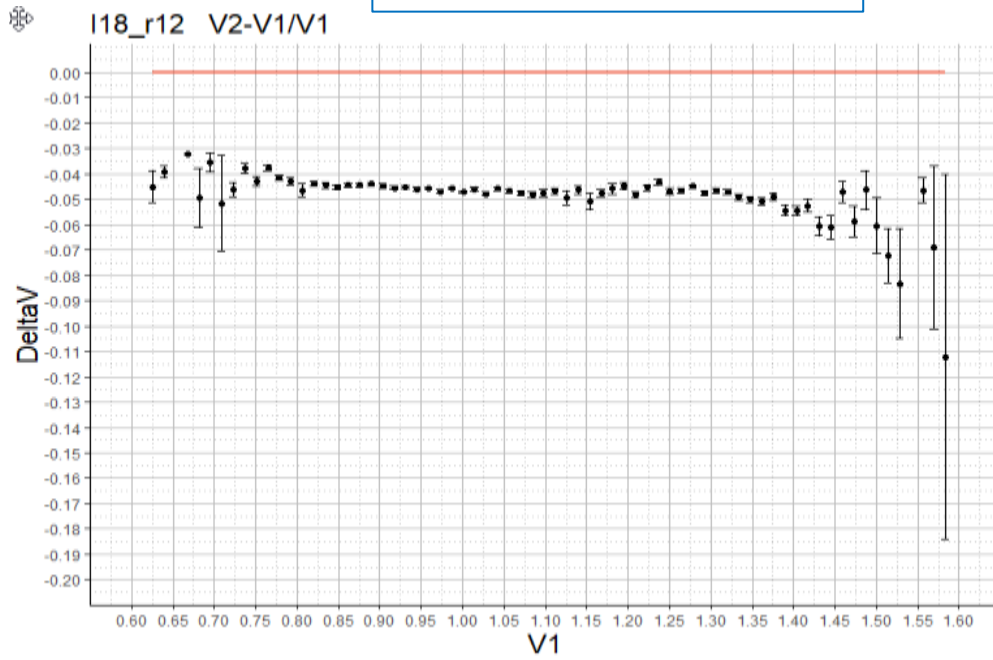
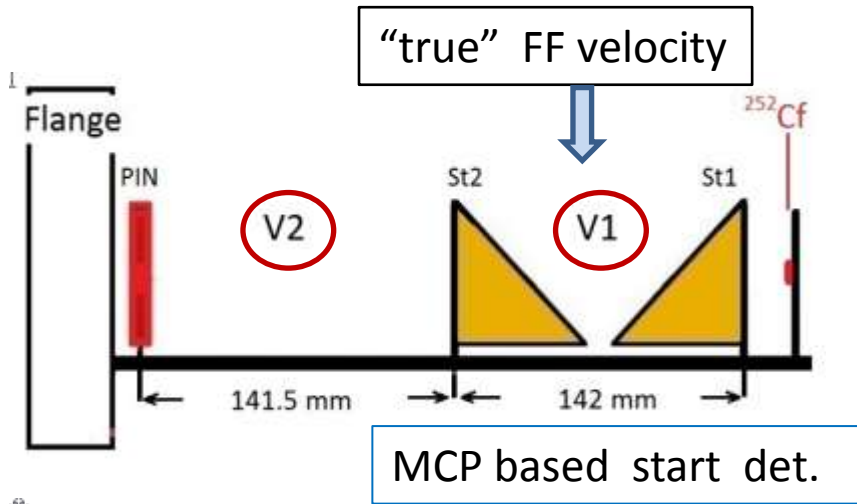
$$\min F = [(\langle ML_T \rangle - \langle ML \rangle)^2 + (\langle MH_T \rangle - \langle MH \rangle)^2] + \mu \sum_{M_{TE}} \frac{(Y(M_{TE}) - Y_T(M_{TE}))^2}{Y(M_{TE})}$$

PD:

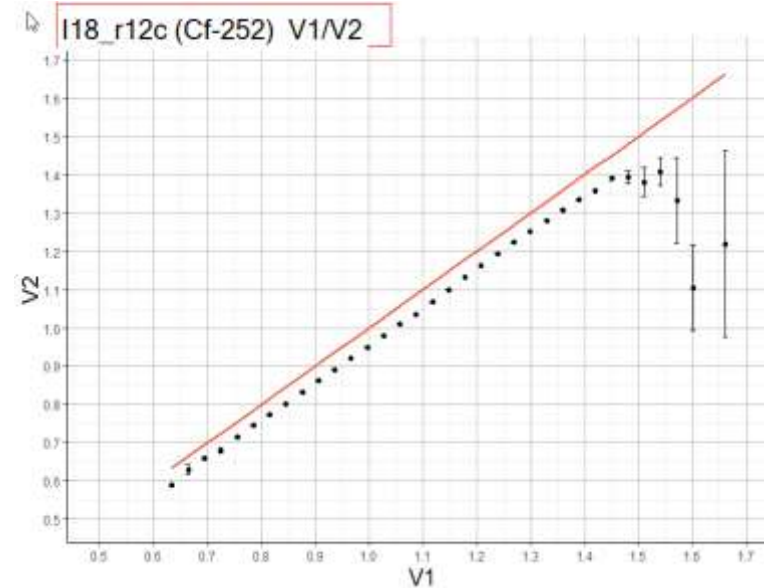
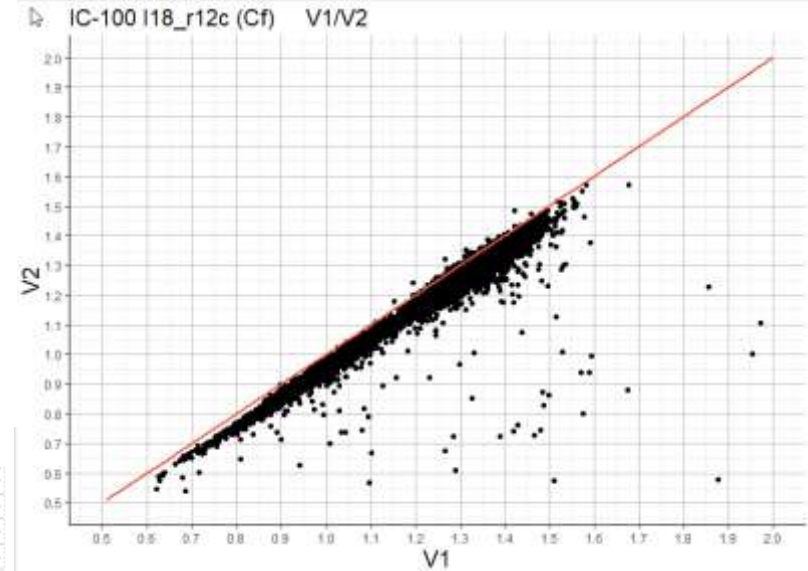
$$\Delta t_p = \gamma \frac{M^{1/6} E^{1/2}}{}$$

A new off-line method of time-pickoff  
"sewing-parabola"

# Validating of the Sewing-parabola method



Comparison of V1&V2:  
good agreement within ~0.5%



## **Feature №2**



## Specific approach to the energy spectrometry with PIN diodes

Peculiarity of the CCT experiments : short intervals of the fragments following while the double-hit regime of their detection is extremely desirable .

**Providing  
the double-hit  
detection  
mode:  
two fragments in  
the  
same detector  
simultaneously**



formation of short pulses  
From PIN diode in the sub-  
microsecond range, but  
normally milliseconds  
formations are adopted  
(?!)

*Sorry for some formulas:*

In the classical approach  $Q \sim E = \int_{-\infty}^{\infty} \frac{i(t)}{dt}$  (1)

i.e. the area under the current graph is calculated  $i(t)$

Let us show that with an arbitrary response  $h(t)$  of the electronic circuit connected to the detector, the area of the output signal S is also proportional to the charge Q.

Then the response of the external circuit  $U(t)$  to the current  $i(t)$ :

$$U(t) = \int_0^t i(r)h(t-r)dr,$$

And

$$S = \int_0^{\infty} U(t)dt = \int_0^{\infty} \left\{ \int_0^t i(r)h(t-r)dr \right\} dt \quad (2)$$

According to the Fumini theorem:

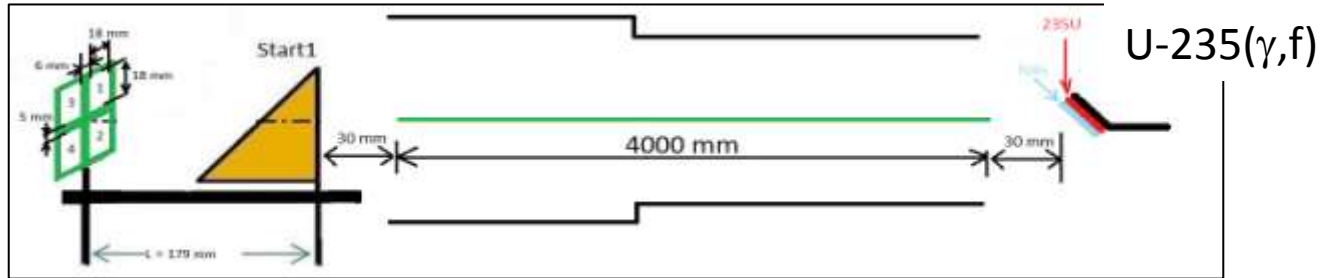
$$S = \int_0^{\infty} U(t)dt = \int_0^{\infty} \left\{ \int_0^t i(r)h(t-r)dr \right\} dt = \int_0^{\infty} i(r)dr \int_r^{\infty} h(t-r)dt$$

After replacing the variables in the inner integral, we have:

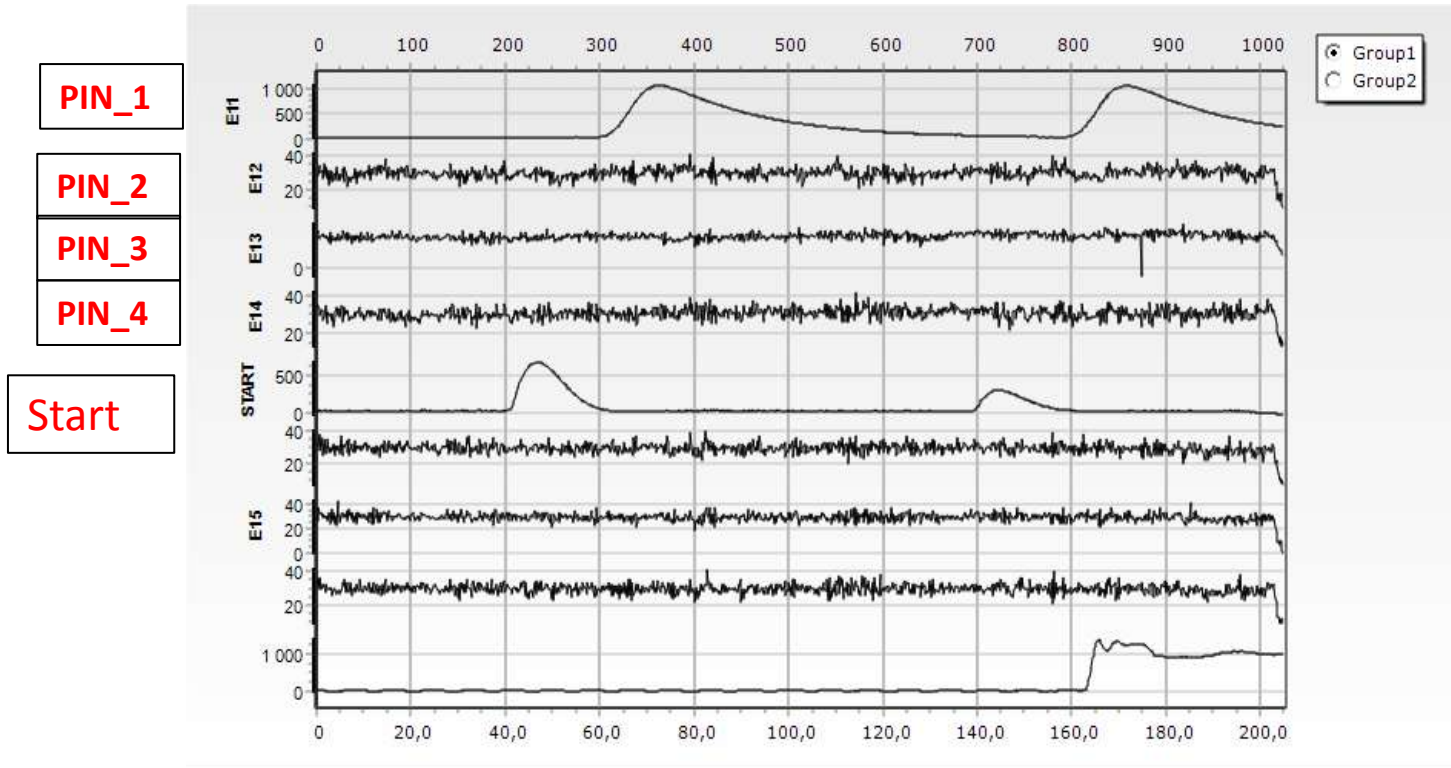
$$S = \int_0^{\infty} i(r)dr \int_r^{\infty} h(t-r)dt = \int_0^{\infty} i(r)dr \int_0^{\infty} h(z)dz = C \int_0^{\infty} i(r)dr = C \cdot Q, \quad (3)$$

where  $C = \text{const}$  is the area of the response function, and  $Q$  is the charge created by the fragment in the detector.

## Hot example of the double-hit success



VEGA (V-E Guide based Array) setup at the MT-25 microtrone in FLNR



**First simultaneous registration of two CCT partners in one spectrometer arm – direct demonstration of the CCT**

## **Feature №3**

## Modernization of the PHD parametrization used

Parametrization utilised earlier:

[Mulgin et al., Two-parametric method for silicon detector calibration in heavy ion and fission fragment spectrometry// Nuclear instruments and Methods in Physics Research A. 388, ( 1997), 254-259]

$$E = E_{det} + R(M,E), \quad (1)$$

$$R(M,E) = \frac{\lambda \cdot E}{1 + \varphi \cdot \frac{E}{M^2}} + \alpha \cdot ME + \beta \cdot E, \quad (2)$$

$$E = \frac{M \cdot V^2}{1.9297} \quad (3)$$

$$\longrightarrow G(\{\lambda, \varphi, \alpha, \beta\}, M, V) = 0$$

Combining equation (1), (2) and (3), we obtain:

$$G = \frac{MV^2}{k} - \left[ E_{det} + \frac{\lambda \cdot \frac{MV^2}{k}}{1 + \varphi \cdot \frac{V^2}{Mk}} + \alpha \cdot \frac{M^2 V^2}{k} + \beta \cdot \frac{MV^2}{k} \right] = 0,$$

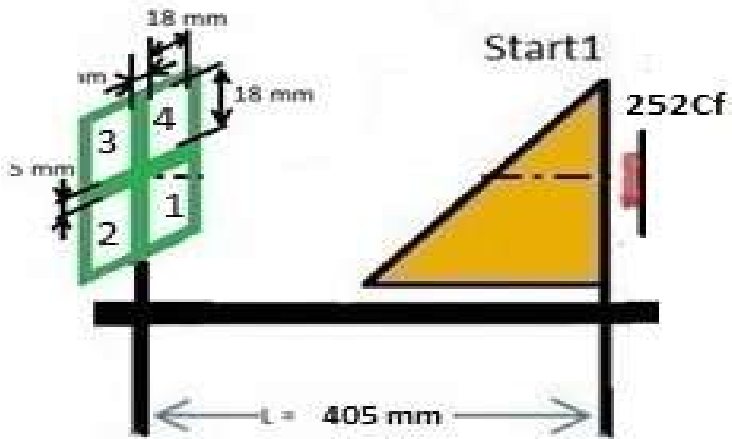
where  $k = 1.9297$ .

$$\min F = [(\langle ML_T \rangle - \langle ML \rangle)^2 + (\langle MH_T \rangle - \langle MH \rangle)^2] + \mu \sum_{M_{TE}} \frac{(Y(M_{TE}) - Y_T(M_{TE}))^2}{Y(M_{TE})}$$

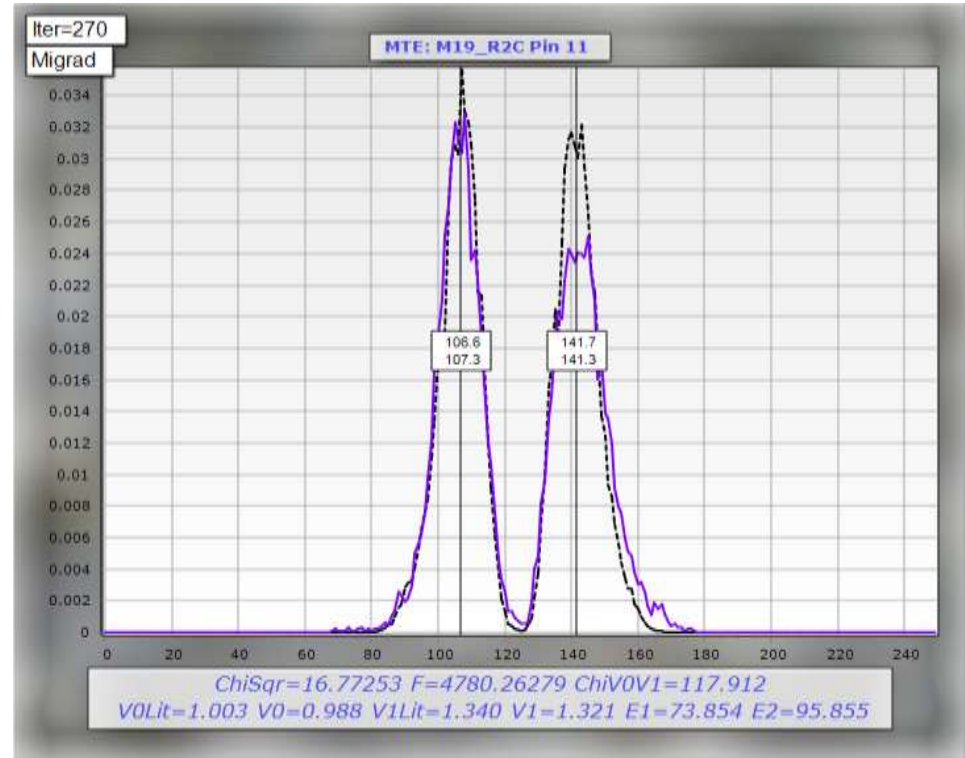
Extended parametrization:  $R(M,E) = \frac{\lambda \cdot E}{1 + \varphi \cdot \frac{E}{M^2}} + \alpha \cdot ME + \beta \cdot E, \longrightarrow \alpha M^{(1+d)} E + \beta M^f E$

Motivation: guaranteed unbiased position of the mass lines

# Testing of the extended PHD parametrization



VEGA detection block



Pin	VI			VH			ML			MH			D	F
	Vel <sub>exp</sub>	Vel <sub>lit</sub>	$\Delta V$	Vel <sub>exp</sub>	Vel <sub>lit</sub>	$\Delta V$	Mass <sub>exp</sub>	Mass <sub>lit</sub>	$\Delta M$	Mass <sub>exp</sub>	Mass <sub>lit</sub>	$\Delta M$		
1	1.340	1.321	0.019	1.003	0.988	0.015	107.25	106.6	0.65	141.28	141.7	-0.42	-0.35	0.47
2	1.340	1.33	0.01	1.003	0.994	0.009	107.25	106	1.25	141.28	142	-0.72	-0.53	-0.55
3	1.340	1.32	0.02	1.003	0.983	0.02	107.25	106.6	0.65	141.28	141.9	-0.62	-0.69	0.49
	1.340	1.316	0.024	1.003	0.983	0.02	107.25	106.4	0.85	141.28	140.8	0.48	-0.19	-0.15

**Conclusion:** additional parameters d & f are significantly different from the previously adopted zero values



# **Conclusion:**

Original hardware and software complexes created in our group proves to be an effective instrument for investigation of the multibody decays. Recent modernization of the data processing algorithms aimed at providing better robustness of the data processing have demonstrated positive results in the latest experiments.

**Thank you for attention!**