

# Super-asymmetric Fission

F. Gönnenwein<sup>1)</sup>, E. Chernysheva<sup>2)</sup>, J.M. Itkis<sup>2)</sup>, M.G. Itkis<sup>2)</sup>, G. Knyasheva<sup>2)</sup> and E.Kozulin<sup>2)</sup>

1) University of Tübingen/Germany, 2) JINR, Laboratory of Nuclear Problems, Dubna/Russia

Nuclear masses are described by the Liquid Drop Model (LDM):

$$M(A,Z) = a_v A + a_s A^{2/3} + a_c Z^2/A^{1/3} + a_i (N-Z)^2/A - \delta(A)$$

Volume, Surface, Coulomb, Symmetry, Pairing

Macroscopic LDM describes average masses.

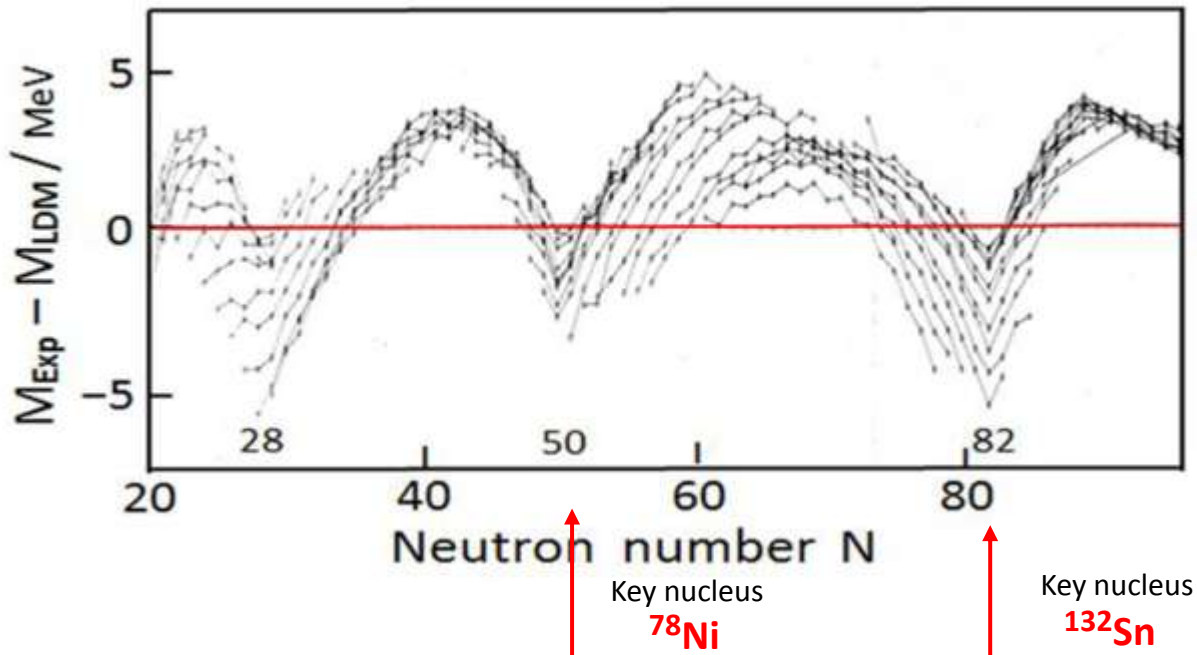
Microscopic nuclear structure necessitates corrections.

$$\delta W = M_{\text{exp}} - M_{\text{LDM}}$$

Myers-Swiatecki 1966: neutron and proton ranges with  $\delta W < 0$  and  $\delta W > 0$  alternate

$\delta W < 0$ : nuclei are more tightly bound than in LDM :  
"SHELLS"

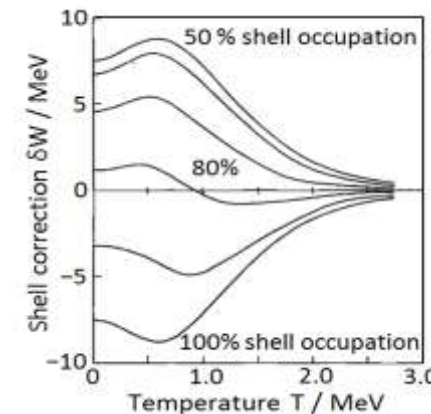
$\delta W > 0$ : nuclei are less tightly bound than in LDM :  
"ANTI-SHELLS"



Influence of shell effects around fragment masses  $A = 132$  u in fission is dramatic: It explains the prominent asymmetry in the mass distributions of fission fragments in the actinides

Influence of shell effects around fragment masses  $A = 78$  is less spectacular. It becomes only visible in detailed analysis of mass and kinetic energy distributions of fission fragments

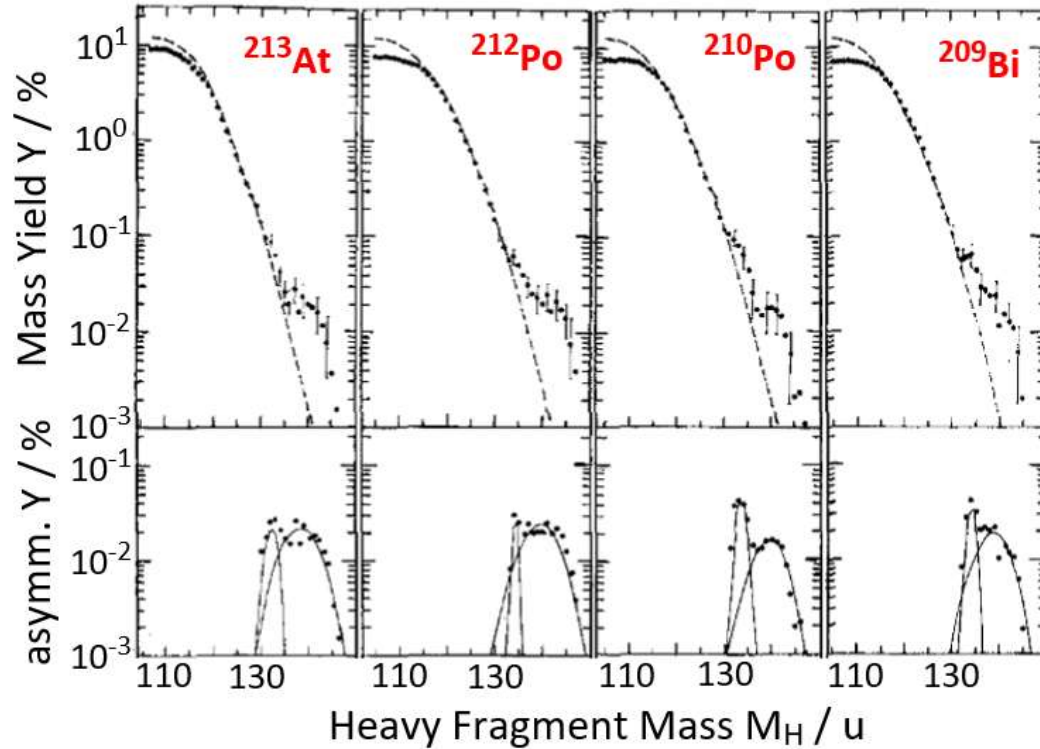
Characteristic property of shells:



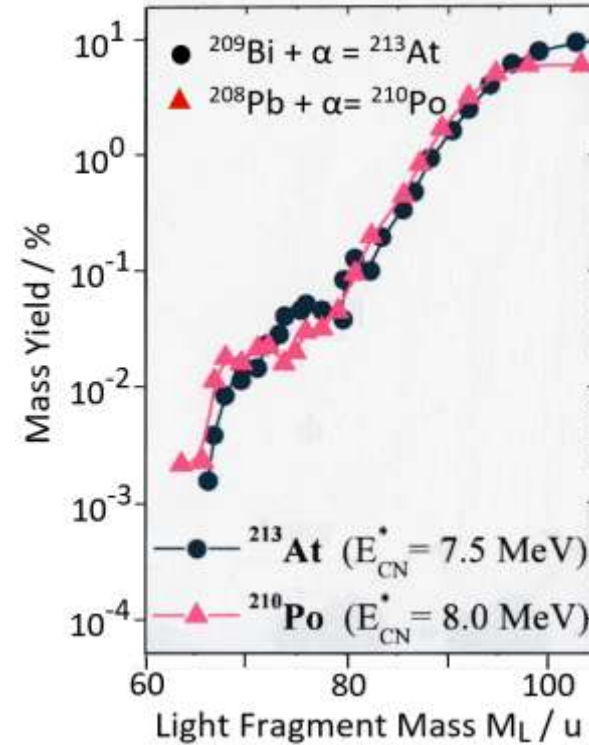
Shell effects disappear when excitation energy is increased

# Supersymmetric fission in the preactinides

$\alpha$ -induced reactions leading to



Itkis 1988



Itkis 1988

Mass distribution is symmetric but in tails appears at large masses shell structure (Brosa modes):

Standard I at  $\langle A_{ST\ I} \rangle = 134\ u$

Standard II at  $\langle A_{ST\ II} \rangle = 139\ u$

KEY NUCLEUS  
 **$^{132}\text{Sn}$**

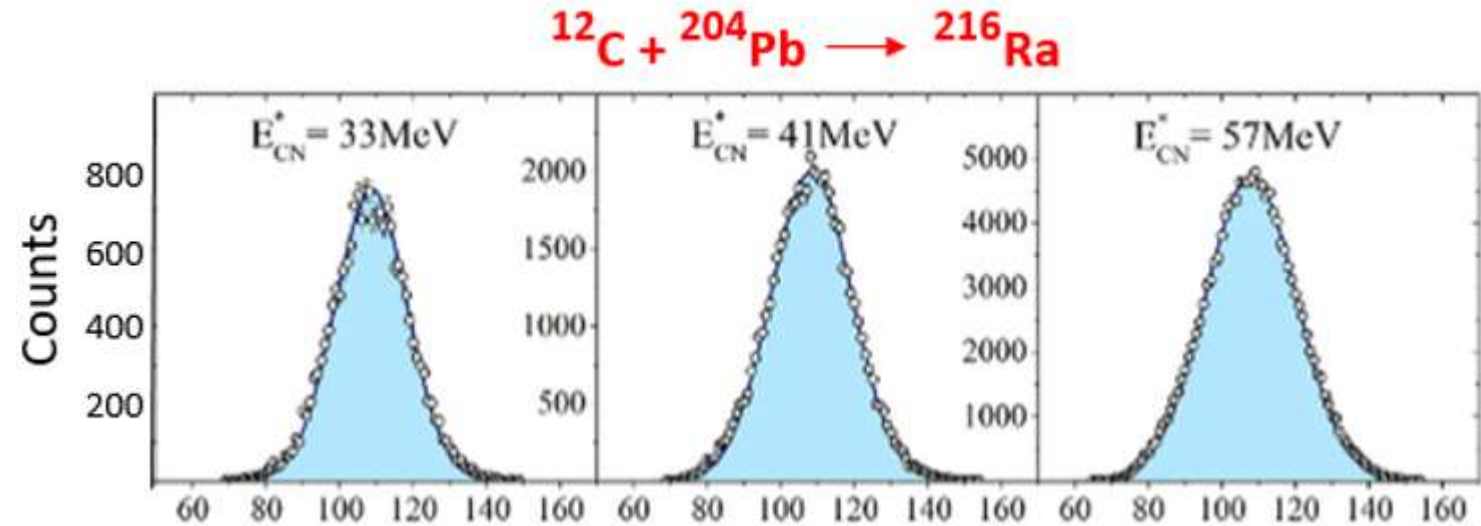
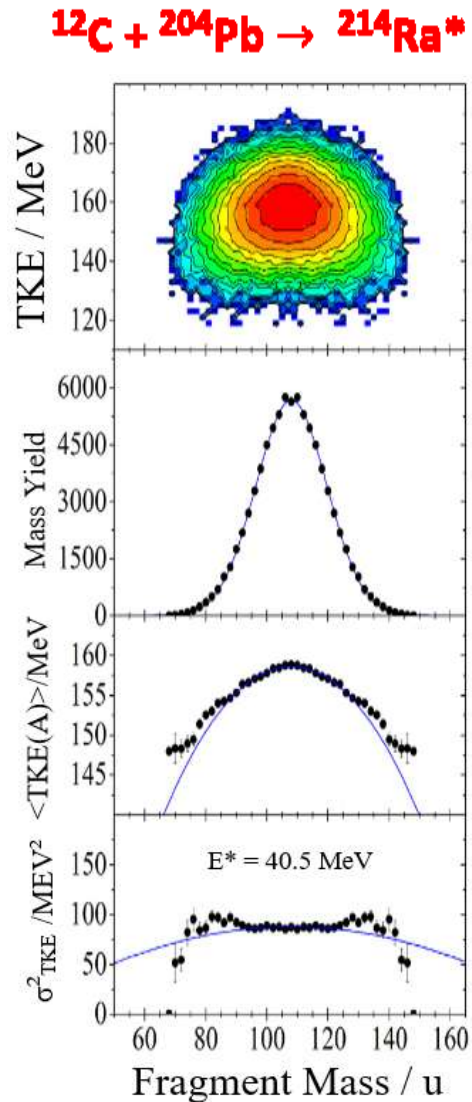
At small fragment masses structure

Standard III at  $\langle A_{ST\ III} \rangle = 78\ u$

KEY NUCLEUS  
 **$^{78}\text{Ni}$**

**Super-asymmetric  
Fission**

# Super-asymmetric Fission of $^{214}\text{Ra}$



**Surprisingly, in the fragment mass distribution of  $^{12}\text{C} + ^{204}\text{Pb} \rightarrow ^{214}\text{Ra}$  there is no indication for shell structure in the restricted mass range  $A_{\text{CN}}/2 \pm 25 \text{ u}$ . By this restriction in mass the ST III mode is cut away. However, in the  $\text{TKE}(A)$  and the  $\sigma^2_{\text{TKE}}(A)$  plots the mode St III is clearly present.**

**Noteworthy : St III stronger than St I and St II**

Fragment Mass / u

# Super Asymmetric Fission in the standard Actinides

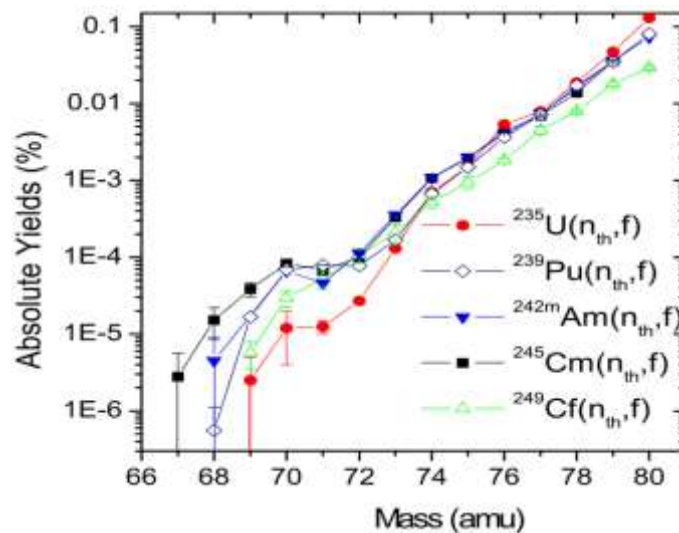
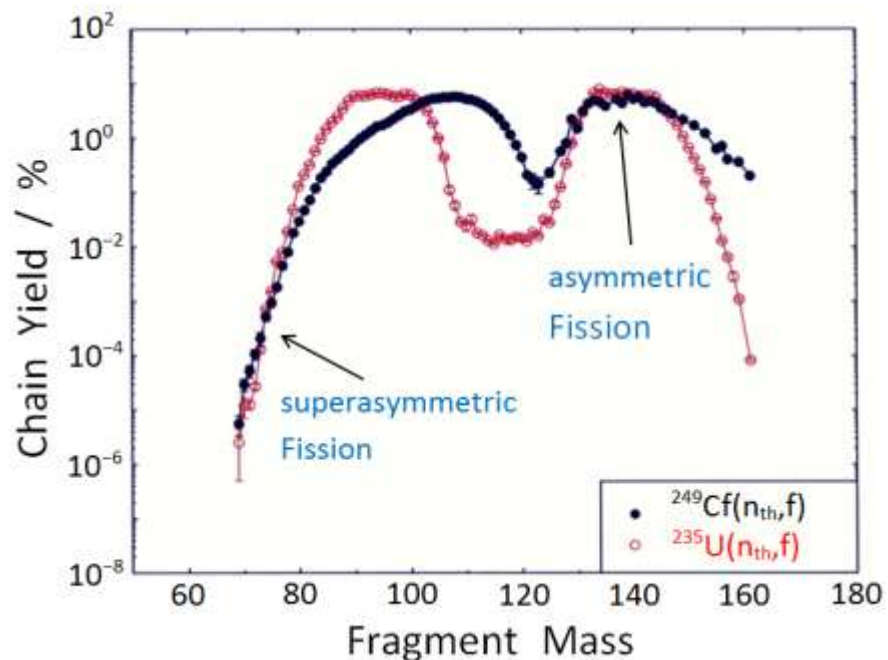
thermal Neutron induced

Shell modes St I and St II fix asymmetric fission.

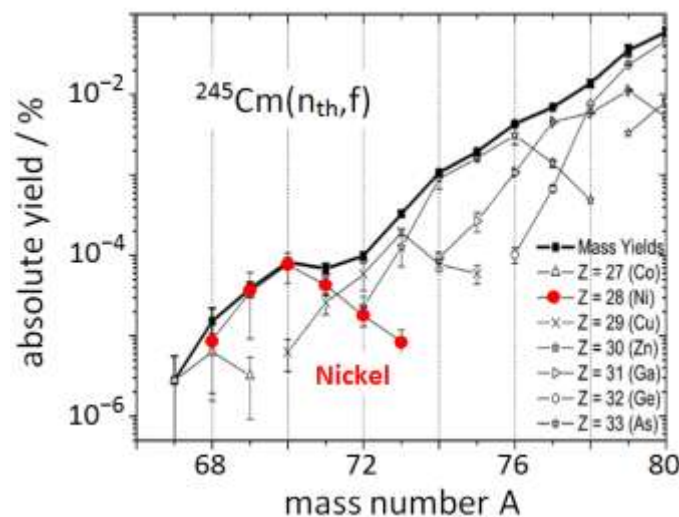
KEY  $^{132}\text{Sn}$

Shell mode St III fixes Super-asymmetric fission

KEY  $^{78}\text{Ni}$



Standard III is present in all  $(n_{th}, f)$  reactions



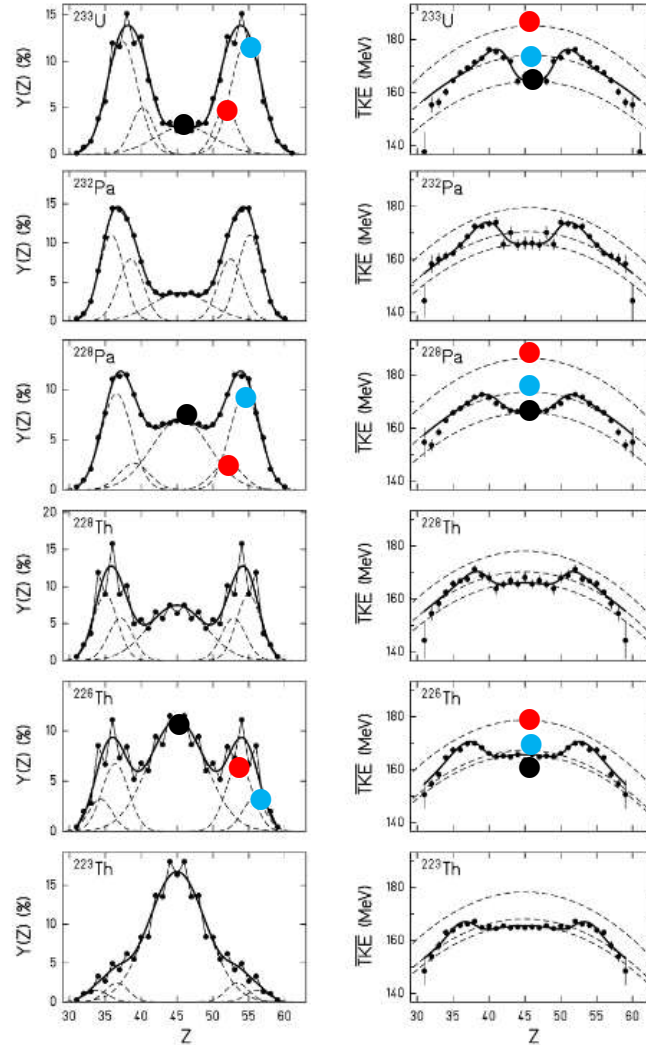
The key nucleus  $^{78}\text{Ni}$  is not observed. But the element Ni is prominent

# Super-asymmetric Fission in electromagnetic Fission ?

●  
Liquid Drop  
Model

●  
Standard I

●  
Standard II



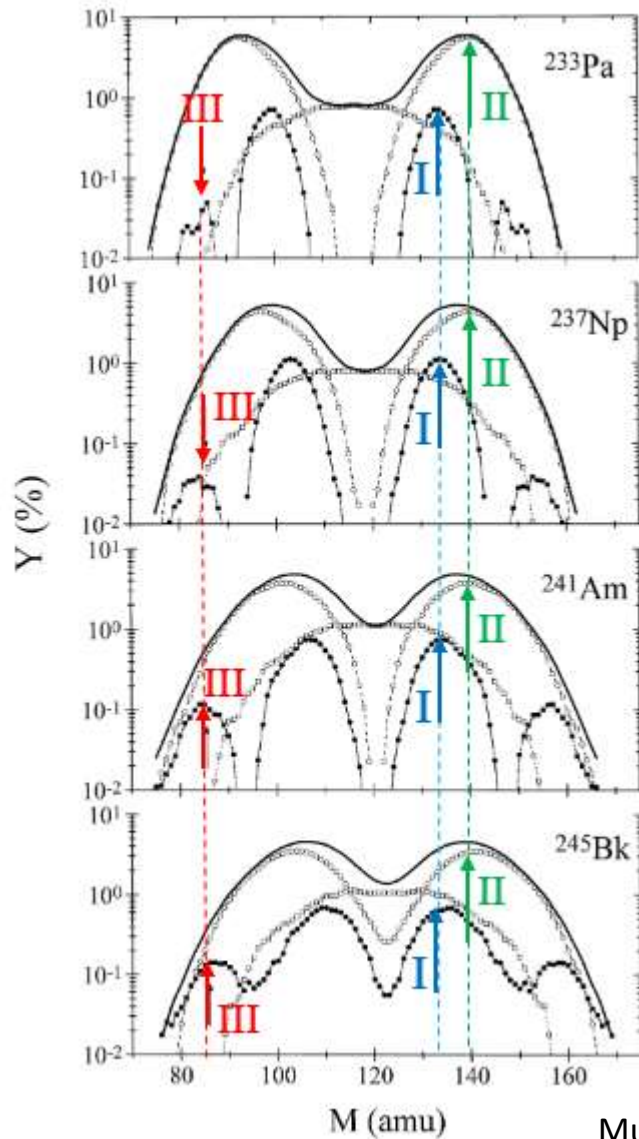
No super-asymmetric St III  
In electromagnetic Fission

Average excitation energy  $E^* \approx 11$  MeV

# Super Asymmetric Fission in the standard Actinides

at intermediate excitation

## 10.3 MeV proton induced Fission



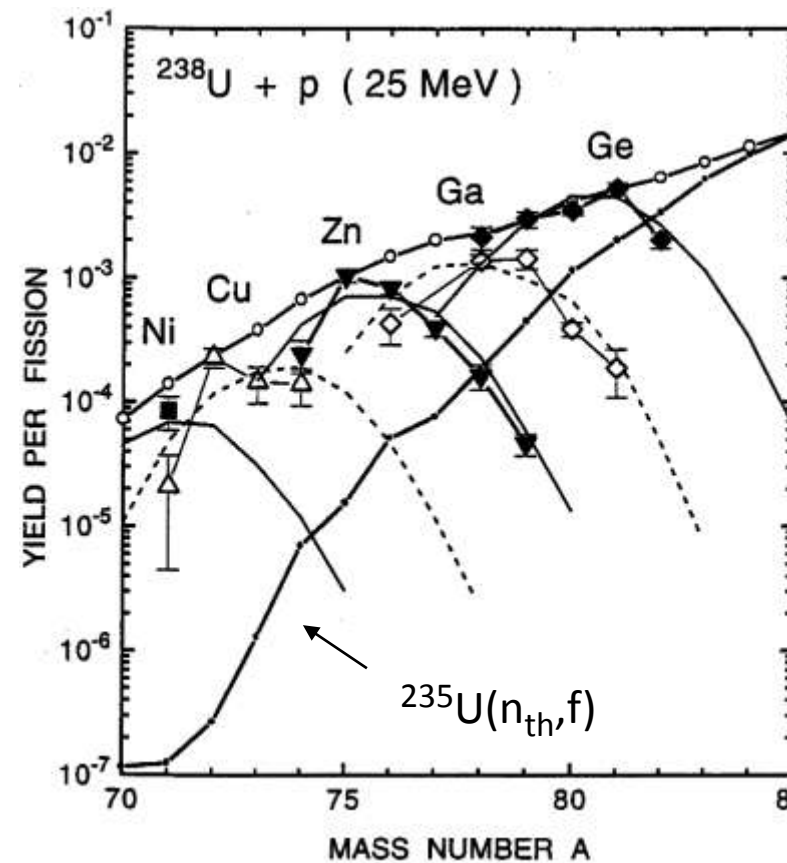
Constant mass position  
for St I and St II  
in heavy mass group

Constant mass position  
for St III  
in light mass group

Mulgin 1999

## Element by element zoom for St III in $p + ^{238}\text{U} \rightarrow ^{239}\text{Np}$

**Yield of St III increases dramatically with  $T$ .**  
Yet also in Liquid Drop Model  $\sigma^2_M \sim T$   
This increase has to be taken into account

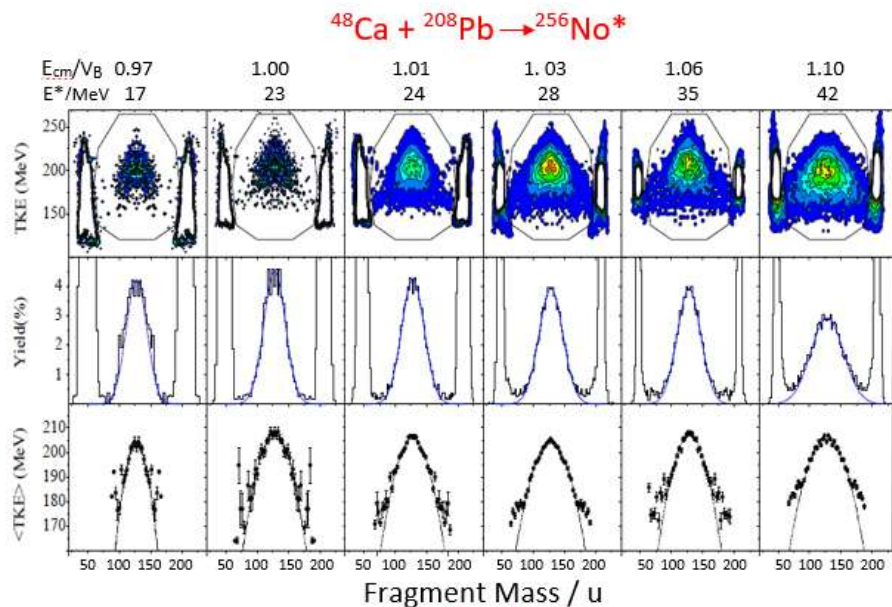


Huhta 1997

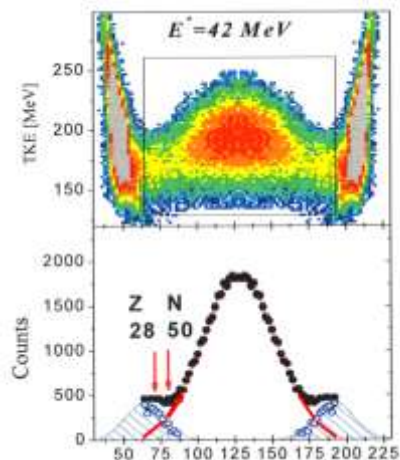
PLB 405,230 (1997)

# Super-asymmetric Fission in $^{256}\text{No}$

Sup.asym. Fission at high excitation



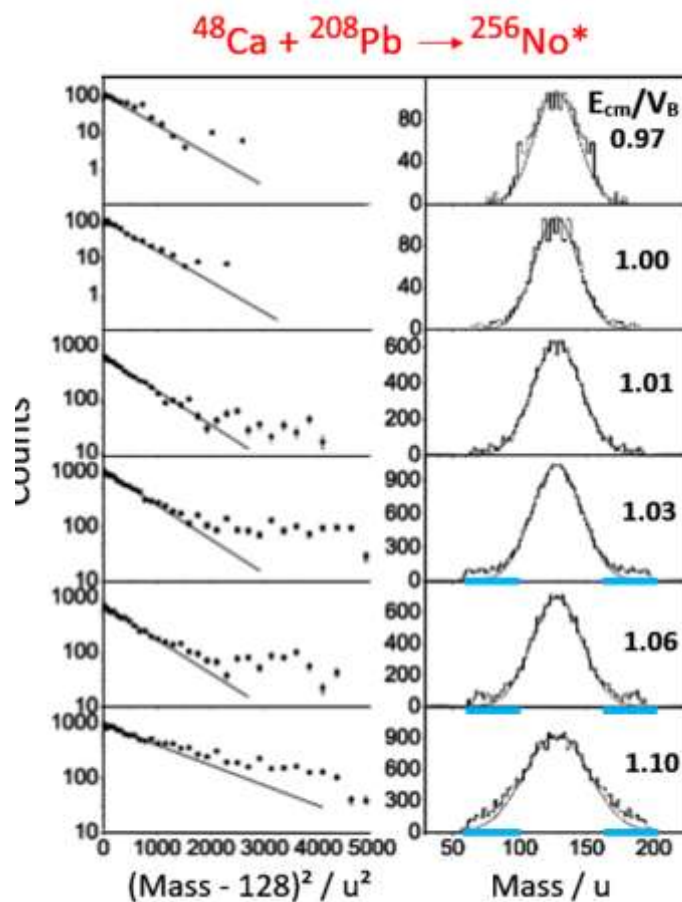
$ZpZ_T = 1640$



Zoom of MED  
at highest  $E^*$

Structure (in blue)  
attributed to  
Super-asymmetric fission  
Key nucleus  $^{78}\text{Ni}$

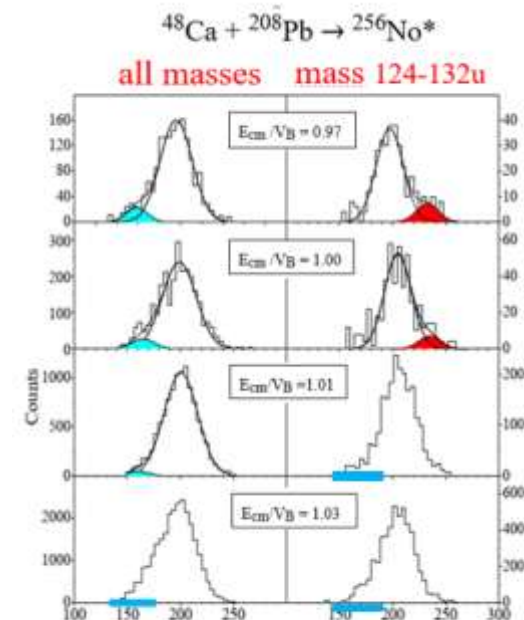
Fragment Mass / u



Flattening of Gaussian fit:  $\sigma_M^2 \sim T$  LDM

Super-asymmetric Fission  $\nearrow$  excitation energy  $E^*$

SAF is Quasi-Fission



TKE / MeV

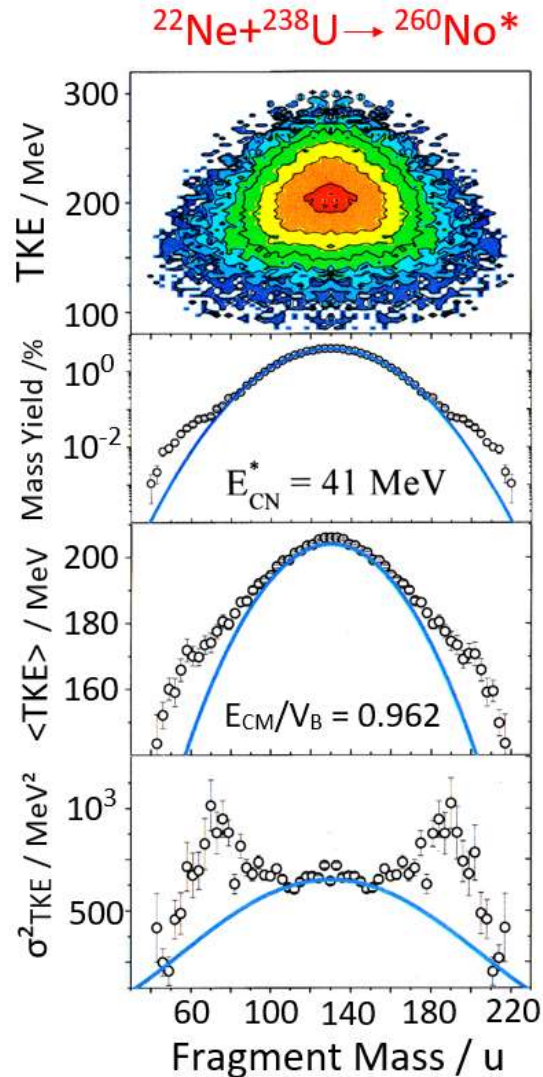
St I and St II at high TKE

St III and SAF at low TKE



# Super-asymmetric Fission in $^{260}\text{No}^*$

Sup.asym. Fission is dominant

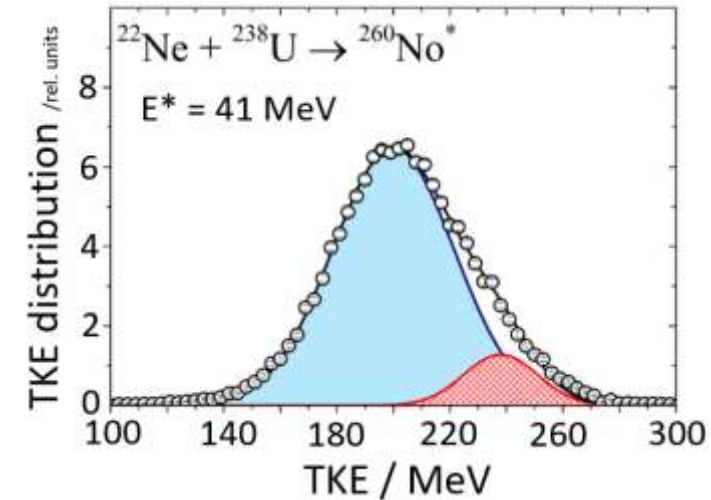


MED plot at  $E_{\text{cm}}/V_{\text{B}} = 0.962$  shows no asym. Quasi-Fission

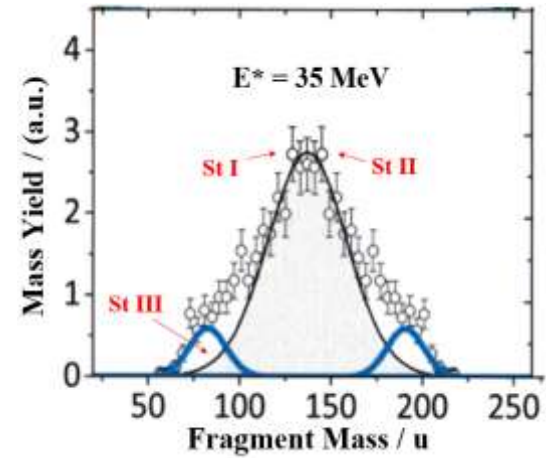
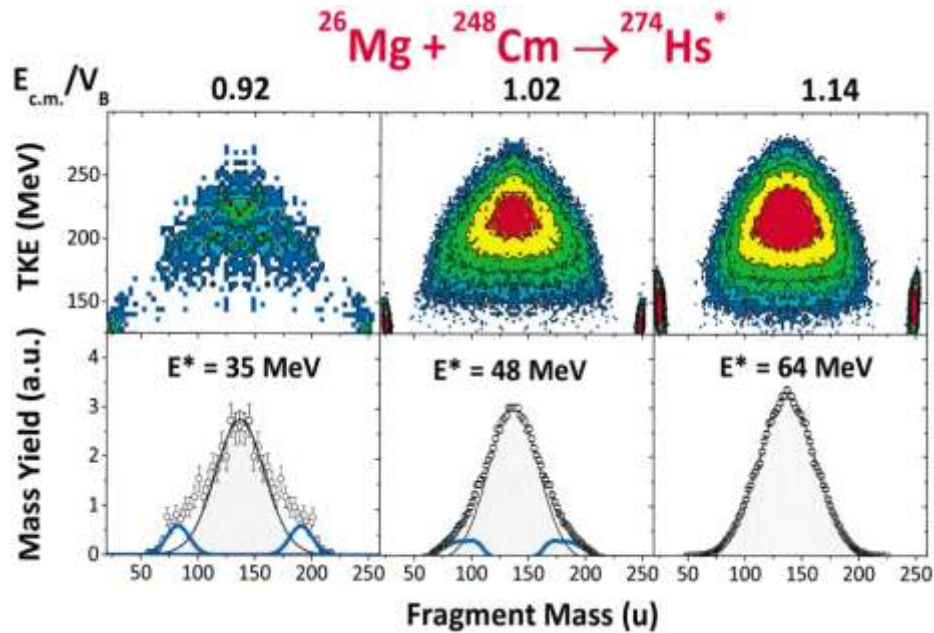
Pronounced structure in far asymmetric fission

Structure traced to **St III** super-asym. Fission

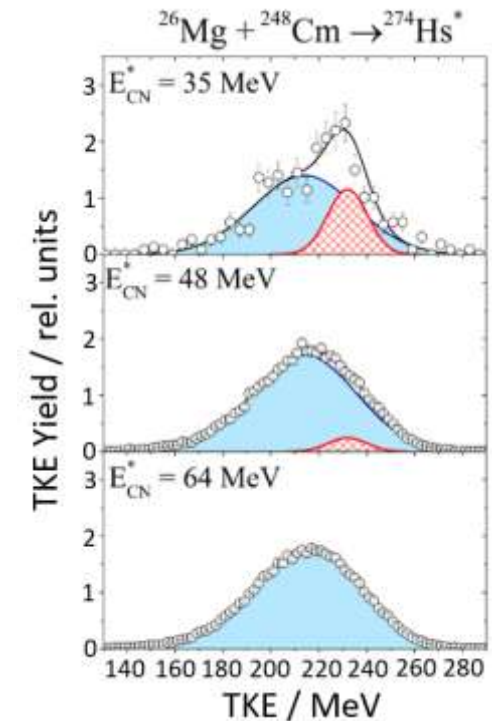
Structure best evidenced by  $\sigma^2_{\text{TKE}}(M)$  with peaks near  $^{78}\text{Ni}$



In heavy ion induced fission of the heavy Actinide  $^{260}\text{No}$  the mode St III is dominant. A trace of the more common modes St I and St II is only visible in the TKE distribution at high TKE energies for fragments near the KEY nucleus  $^{132}\text{Sn}$



Zoom of mass yield  
Strongly pronounced St III  
In CN reaction



St I and St II best seen in TKE  
distribution near symmetry  
 $A_{\text{CN}}/2 \pm 20 \text{ u}$

One of the few examples of Heavy Ion reactions where MED is not overwhelmed by Quasi-Fission. Only for this type of reactions it is possible to start a shell effect analysis of Compound Nucleus Fission.

**In CNF all shell effects fade away at higher excitation**

# Theory of Mass Distributions for FF + QF

Theory of Heavy Ion Reactions identifies the different sources of fission.

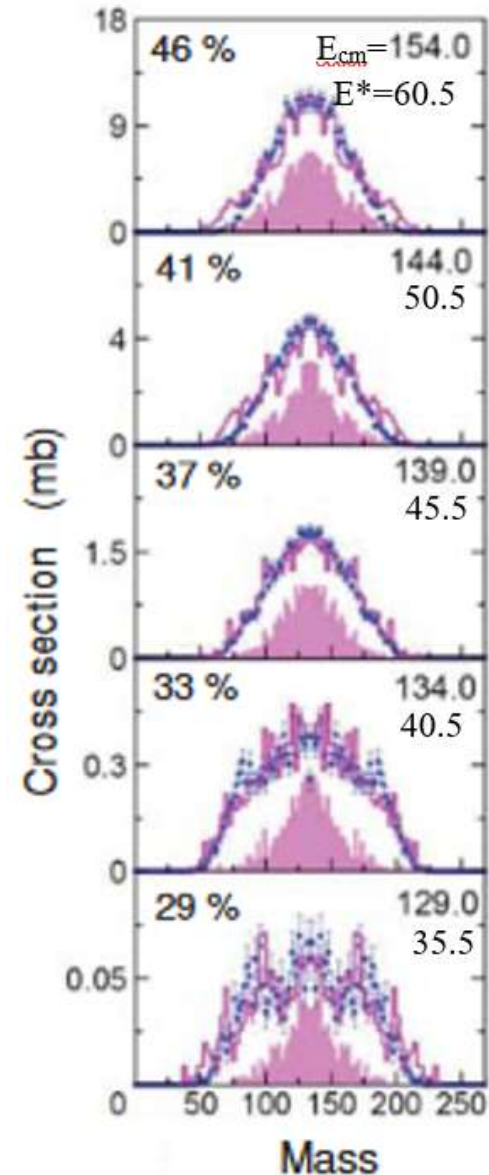
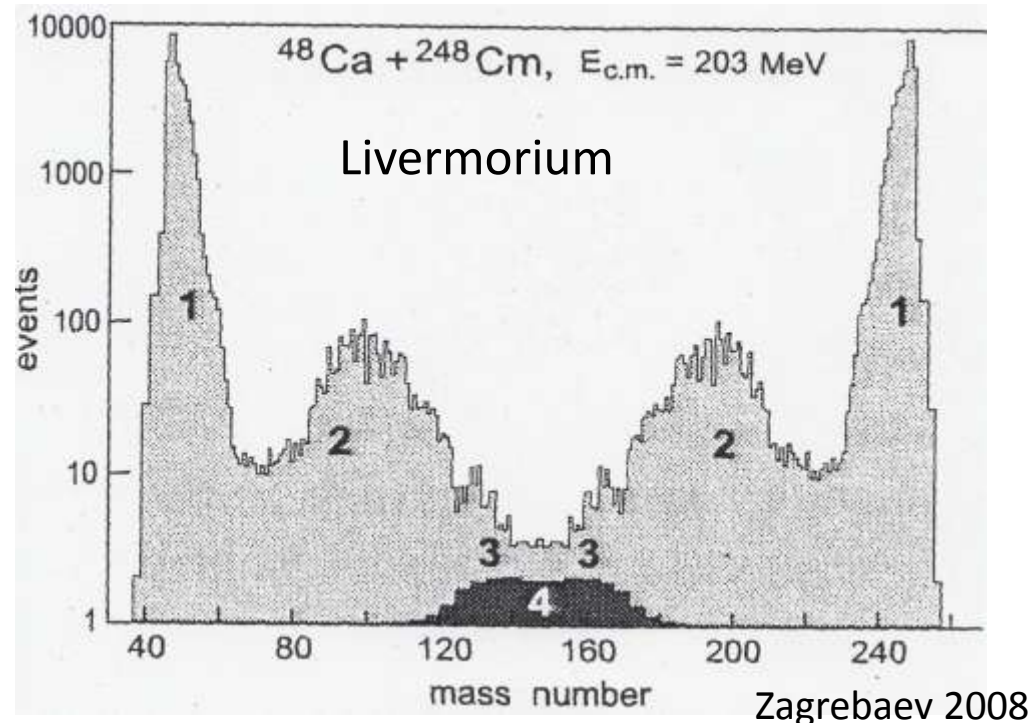
Mass distribution  $Y(A)$  for Livermorium ( $Z = 116$ )

1= Deep Inelastic Scattering

2= Fast asymmetric QF I

3= Slow QF II near to symmetry (not going through CN)

4= Compound Nucleus Fission with symmetric  $Y(A)$



Remarkable:  
Theory predicts QF near symmetry with  $\approx$  constant contribution up to high excitation energies  $E^*$ .

From experiment it is tentatively conjectured that the QF is mainly

**SAMF.**

**Super-Asymmetric Mass Distribution**

Aritomo-Nishio 2008