

**$^{236}\text{sNp}$  isomer yields in  $^{237}\text{Np}(n,2n)$   
and  $^{238}\text{U}(p,3n)$  reactions.**

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**• Vladimir Maslov  
220025, Minsk, Byelorussia**

# Collaboration

- **ADVANCED EVALUATION OF 237Np and 243Am NEUTRON DATA'**
- **COLLABORATION**
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- *V.M. Maslov<sup>1)</sup>, V.P.Pronyaev<sup>2)</sup>, N.A. Tetereva<sup>1)</sup>,*
- *T. Granier<sup>3)</sup>, F.-J. Hambsch<sup>4)</sup>*
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- 1) Minsk-Sosny, Byelorussia
- 2) Institute of Physics and Power Engineering, 249033, Obninsk, Russia
- 3) CEA, Centre DAM-Ile de France, 91927, Arpajon, Cedex, France
- 4) EU-JRC Institute for Reference Materials and Measurements, Geel, Belgium

- Neptunium-237 is a major constituent of the spent nuclear fuel. Main chains
- $^{235}\text{U}(n, \gamma) ^{236}\text{U}(n, \gamma) ^{237}\text{U}(\beta^-) ^{237}\text{Np}$  and  $^{238}\text{U}(n, 2n) ^{237}\text{U}(\beta^-) ^{237}\text{Np}$ .
- The transmutation of the  $^{237}\text{Np}$  in VVER is affected by
- $^{237}\text{Np}(n, \gamma) ^{238}\text{Np}(\beta^-) ^{238}\text{Pu}(n, \gamma) ^{239}\text{Pu}$ .
- $^{236(s)}\text{Np} - ^{237}\text{Np}(n, 2n) ^{236s}\text{Np}(\beta^-) ^{236}\text{Pu}(\alpha) ^{232}\text{U}$

# SCOPE

Reaction chain  $^{237}\text{Np}(n,2n)^{236\text{s}}\text{Np}(\beta^-)$   
 $^{236}\text{Pu}(\alpha)^{232}\text{U}$  is one of the major sources of the accumulation of  $^{232}\text{U}$  in spent fuel.

- The half-life of  $^{236\text{s}}\text{Np}$   $^{236\text{s}}\text{Np}$  is 22.4 h
- The half-life of long-lived state
- $^{237}\text{Np}(n,2n)^{236\text{l}}\text{Np}$  is 155000 y
- $^{236\text{l}}\text{Np}$  thermal fission cross section 2000 barn influence the core neutronics.
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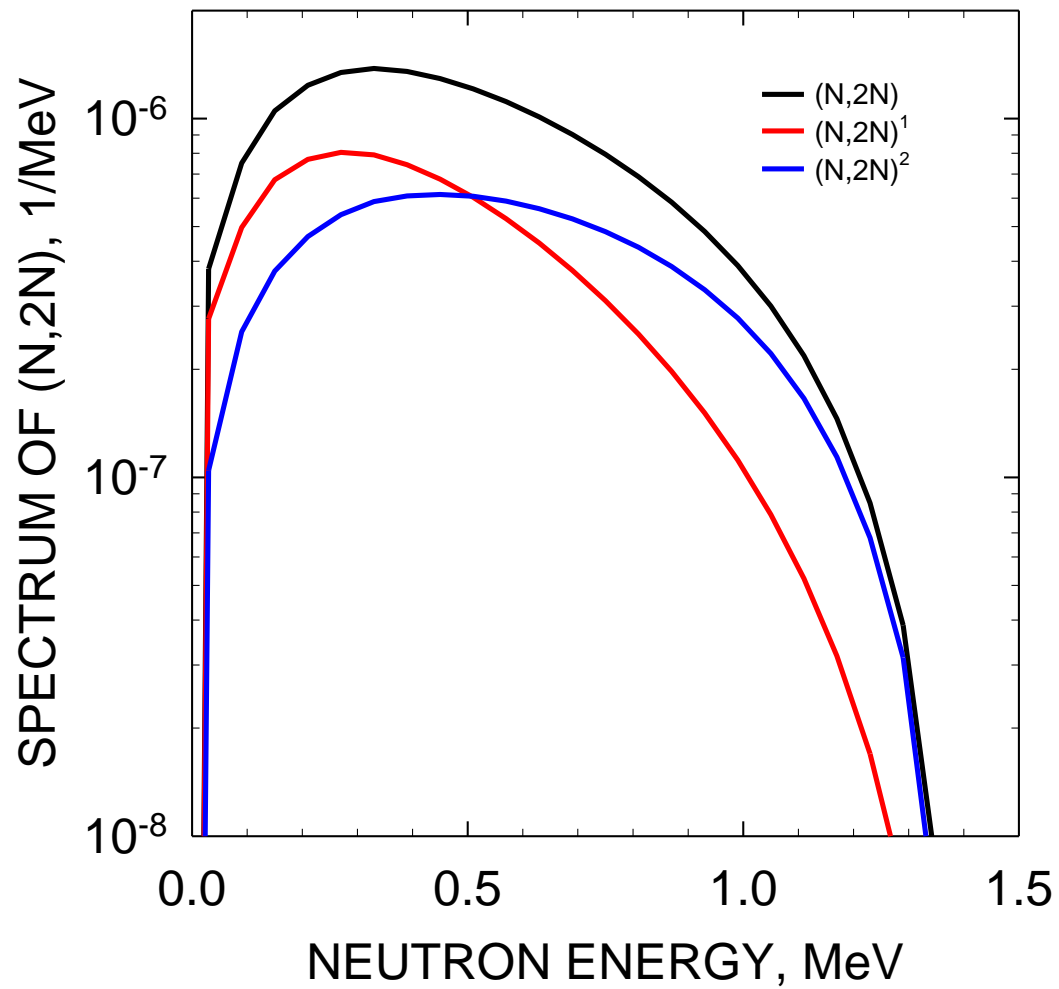
$$\frac{d\sigma_{n2nx}^1}{d\varepsilon} = \frac{d\sigma_{n2nx}^1(\varepsilon)}{d\varepsilon} \frac{\Gamma_n^A(E_n - \varepsilon)}{\Gamma^A(E_n - \varepsilon)}$$

$$\frac{d\sigma_{n2n}^1}{d\varepsilon} = \int_0^{E_n - B_n^A} \frac{d\sigma_{n2nx}^1(\varepsilon)}{d\varepsilon} \frac{\Gamma_n^{A-1}(E_n - B_n^A - \varepsilon - \varepsilon_1)}{\Gamma^{A-1}(E_n - B_n^A - \varepsilon - \varepsilon_1)} d\varepsilon_1$$

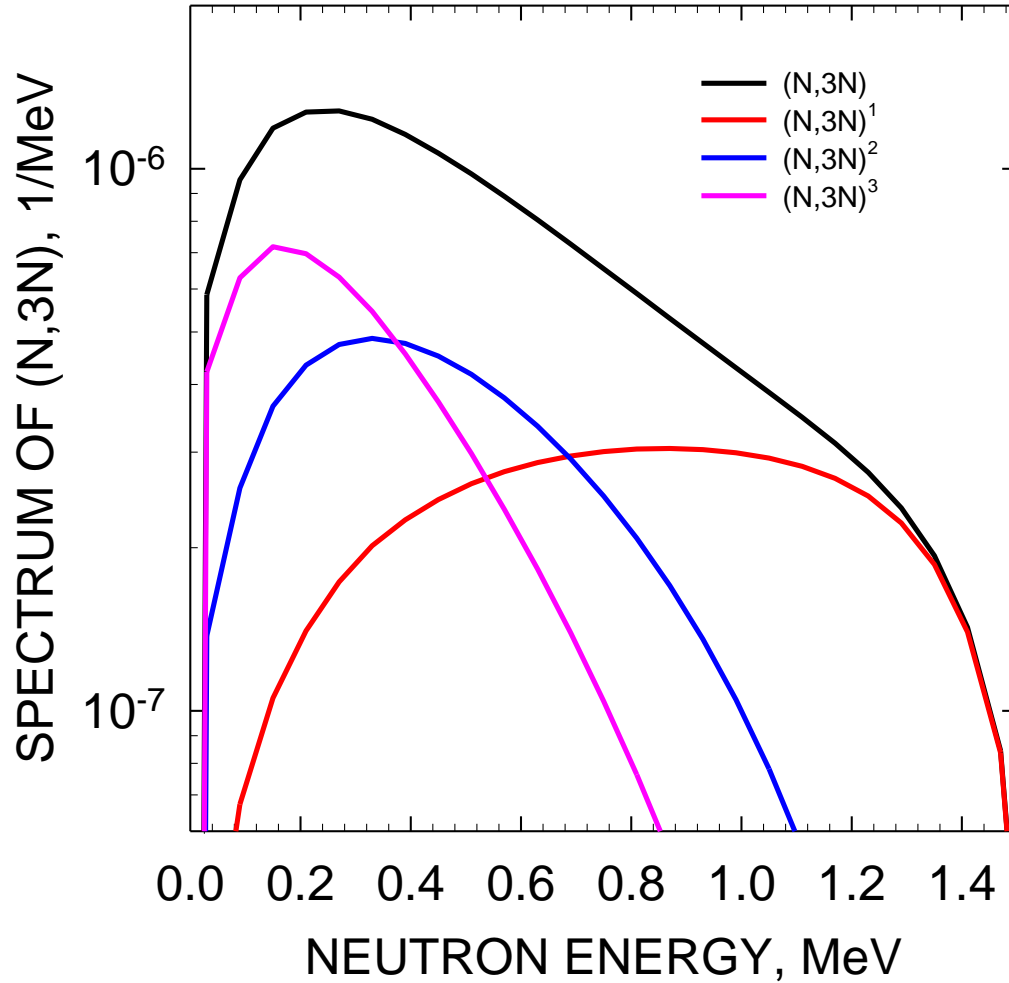
$$\frac{d\sigma_{n2nx}^2}{d\varepsilon} = \int_0^{E - B_n^A - \varepsilon} \frac{d\sigma_{n2nx}^1(\varepsilon)}{d\varepsilon} \frac{\Gamma_n^A(E_n - B_n^A - \varepsilon - \varepsilon_1)}{\Gamma^A(E_n - B_n^A - \varepsilon - \varepsilon_1)} d\varepsilon_1$$

$$\frac{d\sigma_{n2n}^2}{d\varepsilon} = \int_0^{E - B_n} \frac{d\sigma_{n2nx}^2(\varepsilon)}{d\varepsilon} \frac{\Gamma_n^{A-1}(E_n - B_n^A - \varepsilon_1 - \varepsilon_2)}{\Gamma^{A-1}(E_n - B_n^A - \varepsilon_1 - \varepsilon_2)} d\varepsilon_1$$

$^{237}\text{Np}$ ,  $E_n = 8 \text{ MeV}$



$^{237}\text{Np}$ ,  $E_n = 14 \text{ MeV}$



The  $\gamma$ -decay of the excited nucleus is described by the kinetic equation

$$\frac{\partial \omega_k(U, J^\pi, t)}{\partial t} = \sum_{J'\pi'} \int_0^{U_g} \omega_{k-1}(U', J^{\pi'}, t) \frac{\Gamma_\gamma(U', J^{\pi'}, U, J^\pi)}{\Gamma(U', J^{\pi'})} dt - \omega_k(U, J^\pi, t) \frac{\Gamma_\gamma(U, J^\pi)}{\Gamma(U, J^\pi)}$$

$\omega_k(U, J^\pi)$  is the population of state  $J^\pi$  at  $U$  after emission of  $k$   $\gamma$  – quanta

Integrating over  $t$ , in the long run, one gets  $W(U, J^\pi)$  after emission of  $k$   $\gamma$ -quanta

$$\omega_k(U, J^\pi, \infty) - \omega_k(U, J^\pi, 0) = \sum_{J'\pi'} \int_U^{U_g} \frac{\Gamma_\gamma(U', J^{\pi'}, U, J^\pi)}{\Gamma(U', J^{\pi'})} \int_0^\infty \omega_{k-1}(U', J^{\pi'}, t) dt dU' - \frac{\Gamma_\gamma(U, J^\pi)}{\Gamma(U, J^\pi)} \int_0^\infty \omega_k(U, J^\pi, t) dt$$



$$\frac{\partial \omega_k(U, J^\pi, t)}{\partial t} = \sum_{J'^\pi} \int_0^{U_g} \omega_{k-1}(U', J'^\pi, t) \frac{\Gamma_\gamma(U', J'^\pi, U, J^\pi)}{\Gamma(U', J'^\pi)} dt - \omega_k(U, J^\pi, t) \frac{\Gamma_\gamma(U, J^\pi)}{\Gamma(U, J^\pi)}$$

$$\omega_k(U, J^\pi, t=0) = \delta_{k0} \omega_0(U, J^\pi)$$

$$\omega_k(U, J^\pi, \infty) - \omega_k(U, J^\pi, 0) = \sum_{J'^\pi} \int_U^{U_g} \frac{\Gamma_\gamma(U', J'^\pi, U, J^\pi)}{\Gamma(U', J'^\pi)} \int_0^\infty \omega_{k-1}(U', J'^\pi, t) dt dU' -$$

$$\frac{\Gamma_\gamma(U, J^\pi)}{\Gamma(U, J^\pi)} \int_0^\infty \omega_k(U, J^\pi, t) dt$$

$$W_k(U, J^\pi) = \frac{\Gamma_\gamma(U, J^\pi)}{\Gamma(U, J^\pi)} \int_0^\infty \omega_k(U, J^\pi, t) dt$$

$W_k(U, J^\pi)$  population of state after emission of k gamma-quanta

$$W_k(U, J^\pi) = \sum_{J'\pi'} \int_U^{U_g} \frac{\Gamma_\gamma(U', J^{\pi'}, U, J^\pi)}{\Gamma(U', J^{\pi'})} W_{k-1}(U', J^{\pi'}) dU' + \omega_k(U, J^\pi, 0)$$

$$W(U, J^\pi) = \sum_k W_k(U, J^\pi)$$

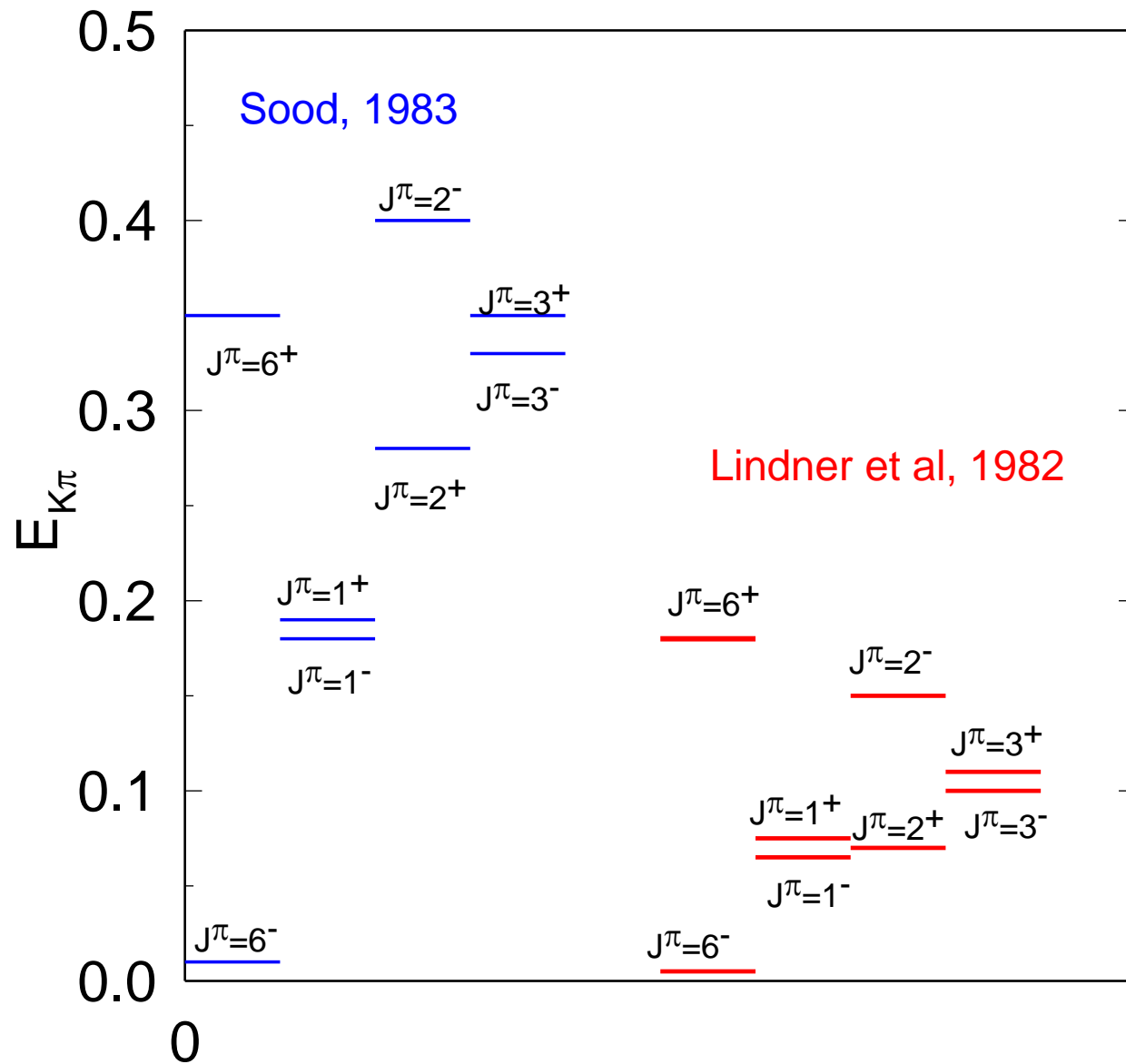
$$W(U, J^\pi) = \sum_{J'\pi'} \int_U^{U_g} \frac{\Gamma_\gamma(U', J^{\pi'}, U, J^\pi)}{\Gamma(U', J^{\pi'})} W(U', J^{\pi'}) dU' + W_0(U, J^\pi)$$

$$r(E_n) = \frac{\sum_{J > (J_l + J_s)/2} W(U, J^\pi)}{\sum_{J \leq (J_l + J_s)/2} W(U, J^\pi)}$$

$$r(E_n) = \sigma_{n2n}^l(E_n) / \sigma_{n2n}^s(E_n)$$

$$\sigma_{n2n}^s(E_n) = \sigma_{n2n}(E_n) / (1 + r(E_n))$$

# $^{236}\text{Np}$ levels



Splitted Gallaher-Moshkowski doublets

$$K^+ = \left| K_n + K_p \right| \quad K^- = \left| K_n - K_p \right|$$

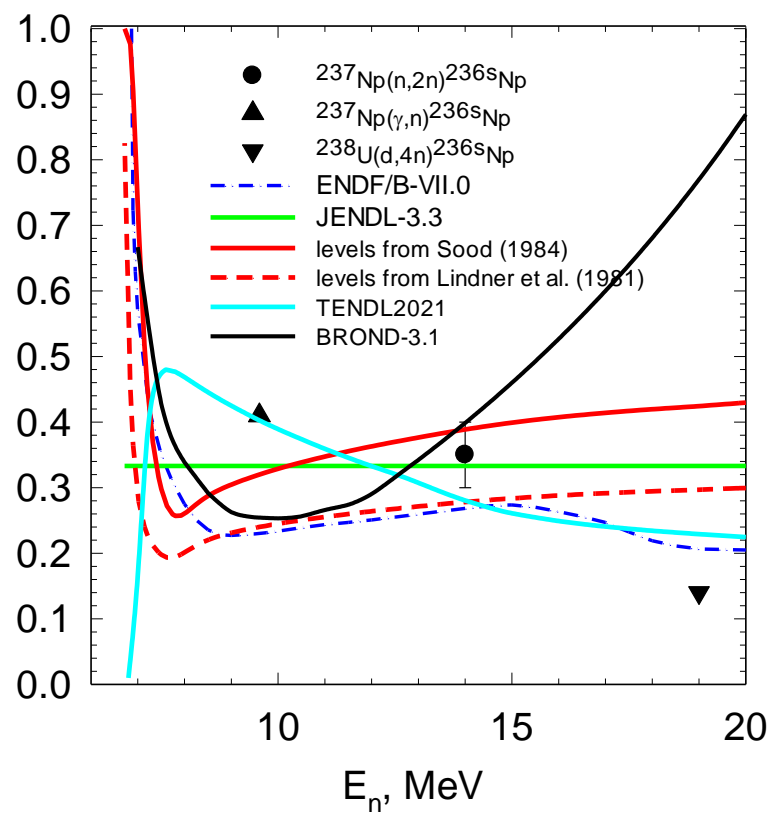
Rotational bands are build on two-quasiparticle states

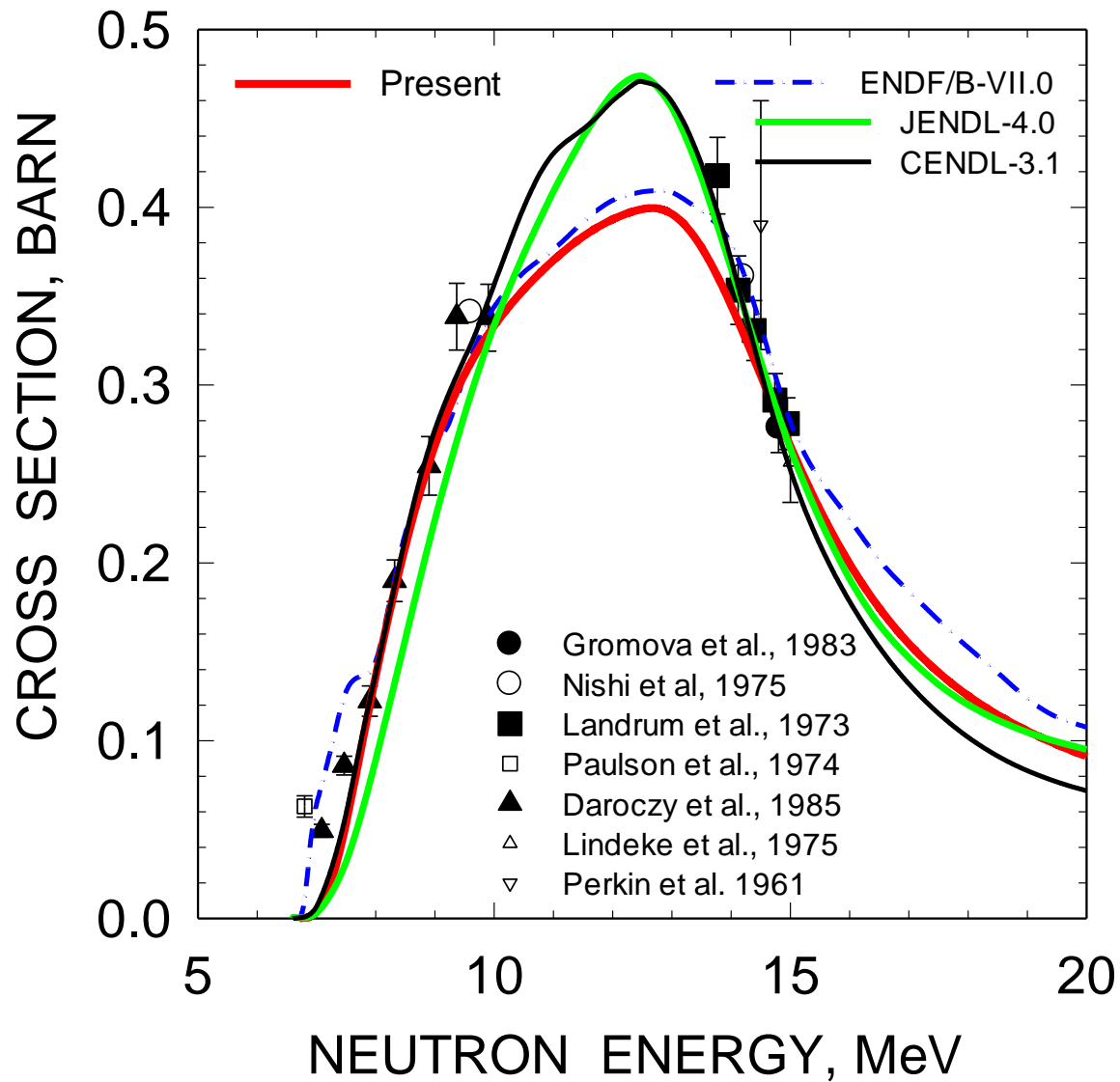
$$E_{JK\pi} = E_{JK} + 5.5 \left[ J(J+1) - K(K+1) \right]$$

$$N(U) = e^{2\Delta_0/T} (e^{U/T} - 1)$$

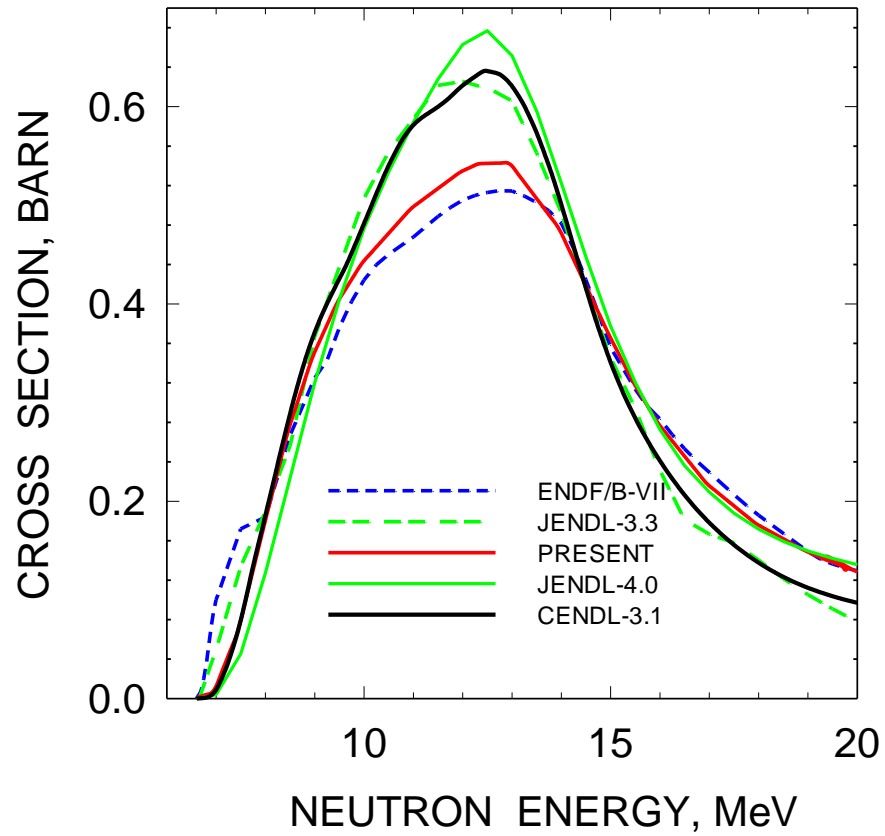
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### Isomer Ratio $^{237}\text{I}_{\text{Np}}/^{237}\text{S}_{\text{Np}}$



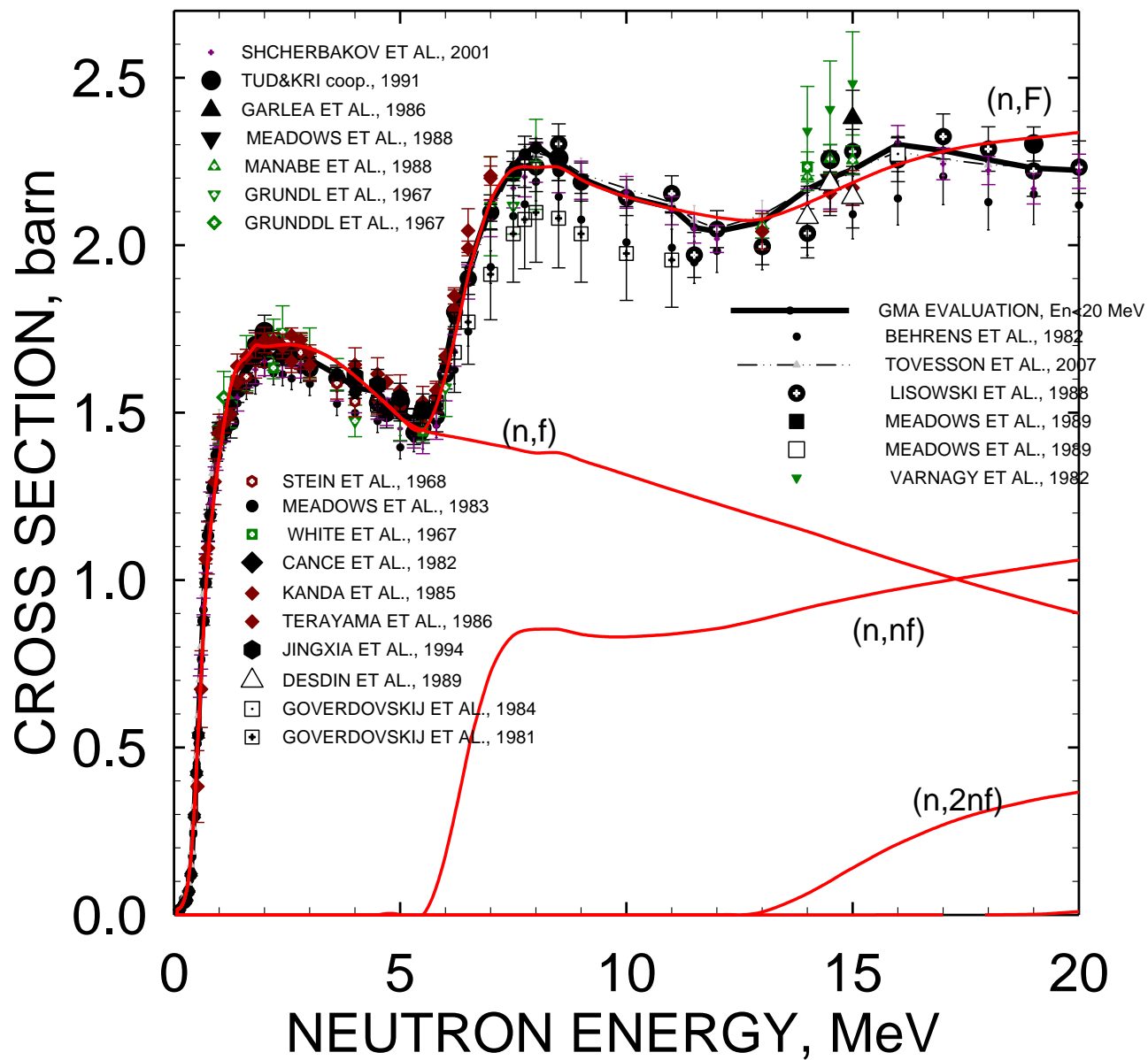
$^{237}\text{Np}(n,2n)$ 

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 $^{237}\text{Np}$  (N,2N) CROSS SECTION

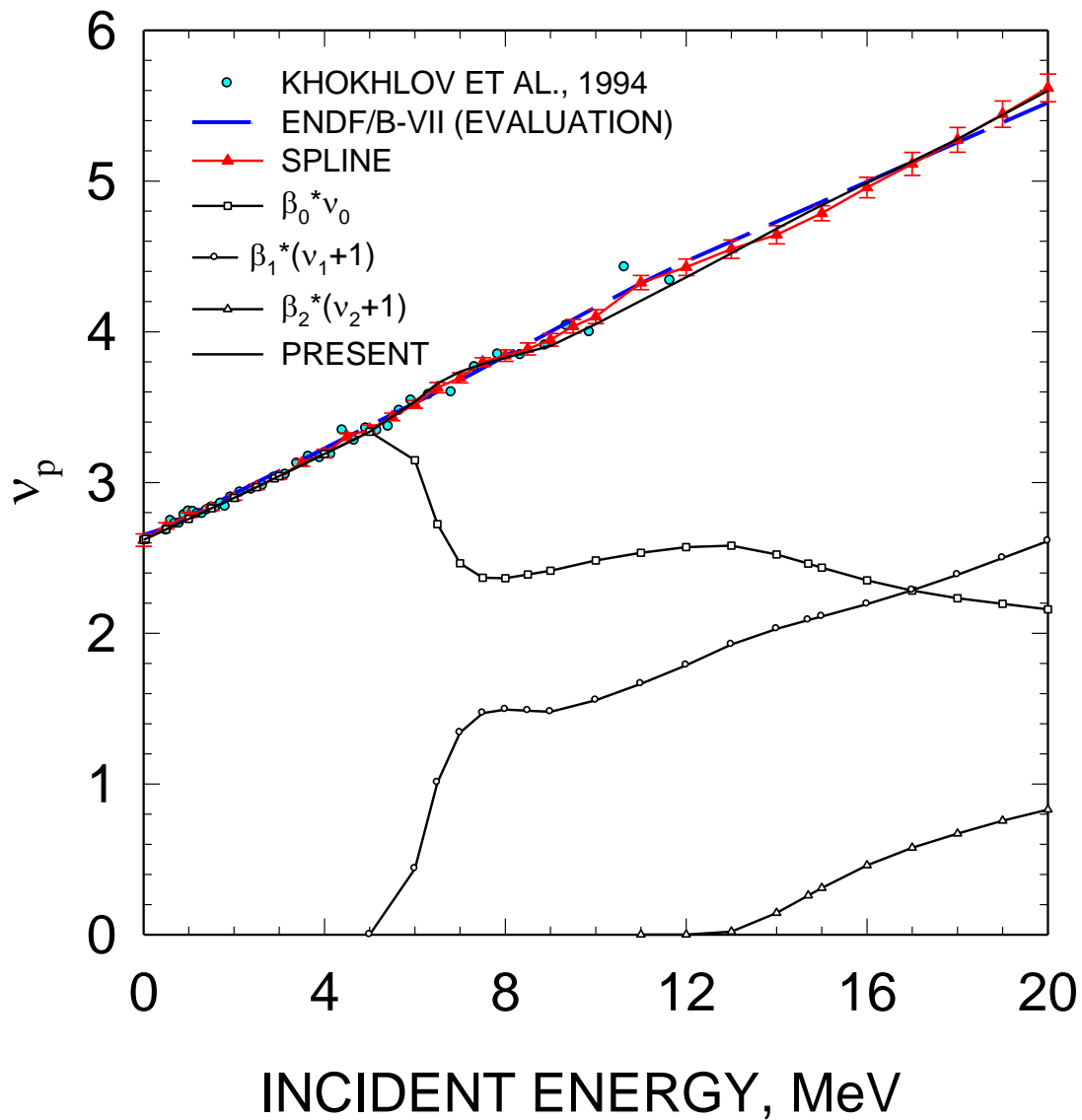




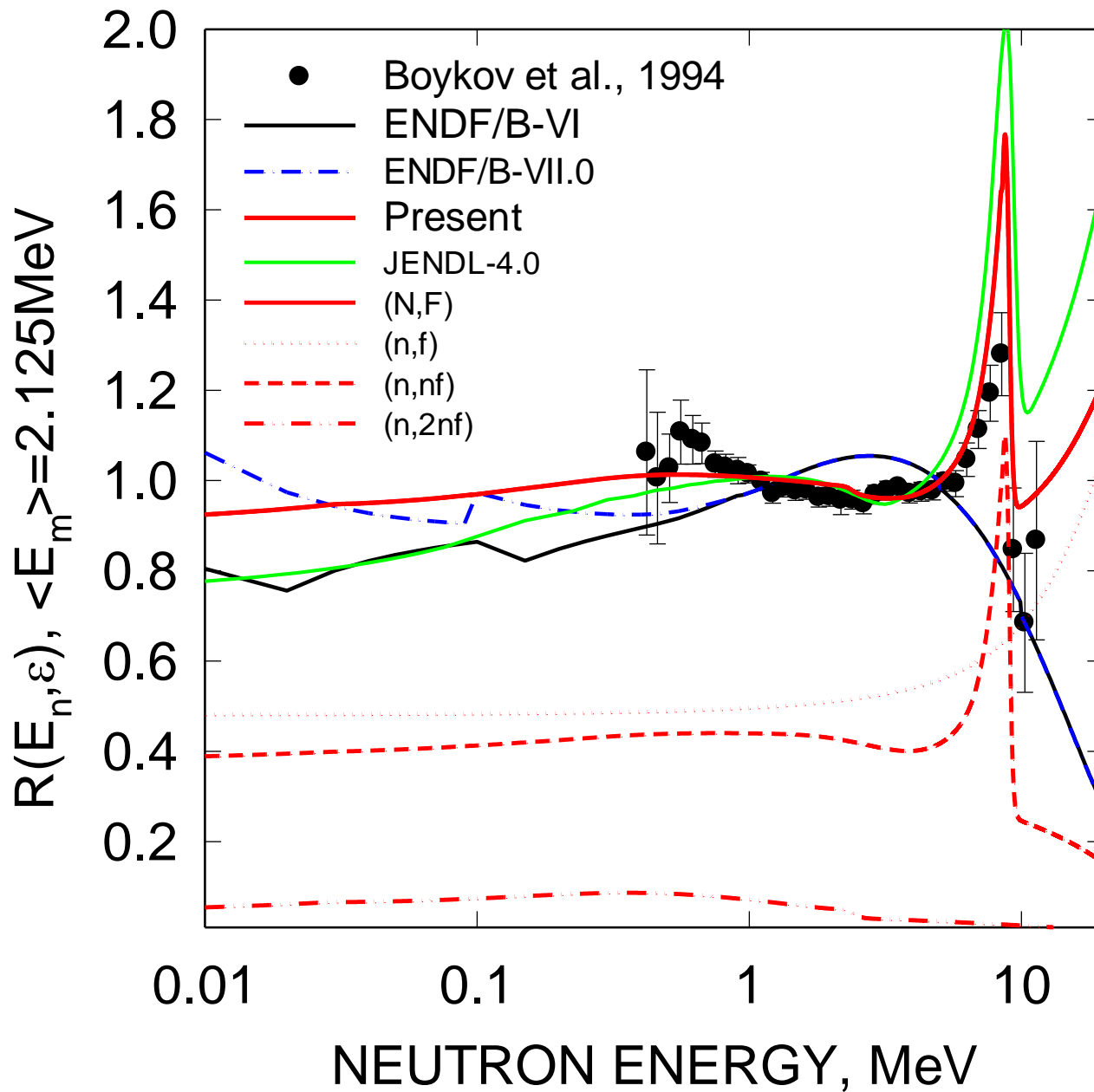
# $^{237}\text{Np}$ FISSION CROSS SECTION



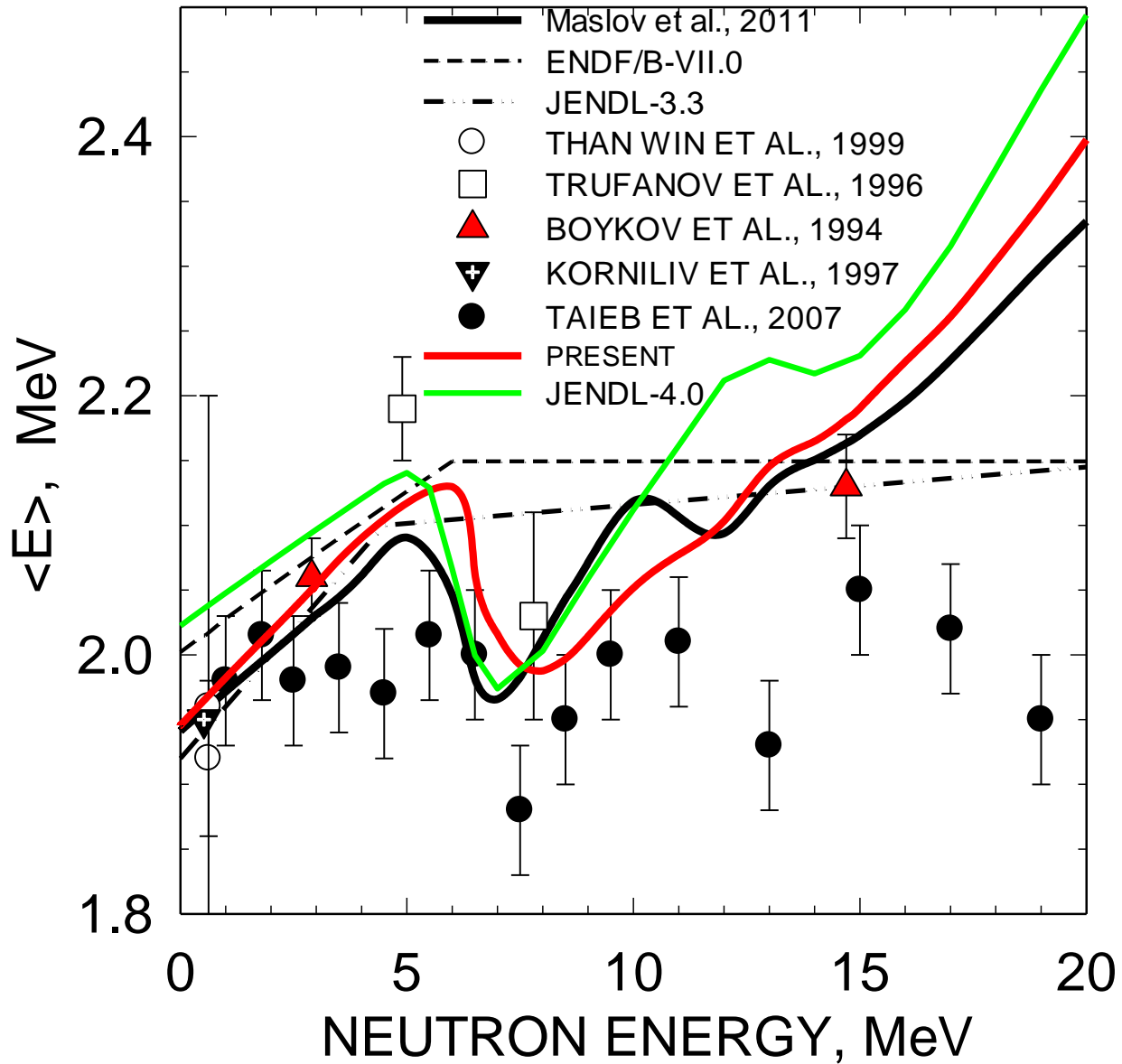
# $^{237}\text{Np}$ NEUTRON MULTIPLICITY



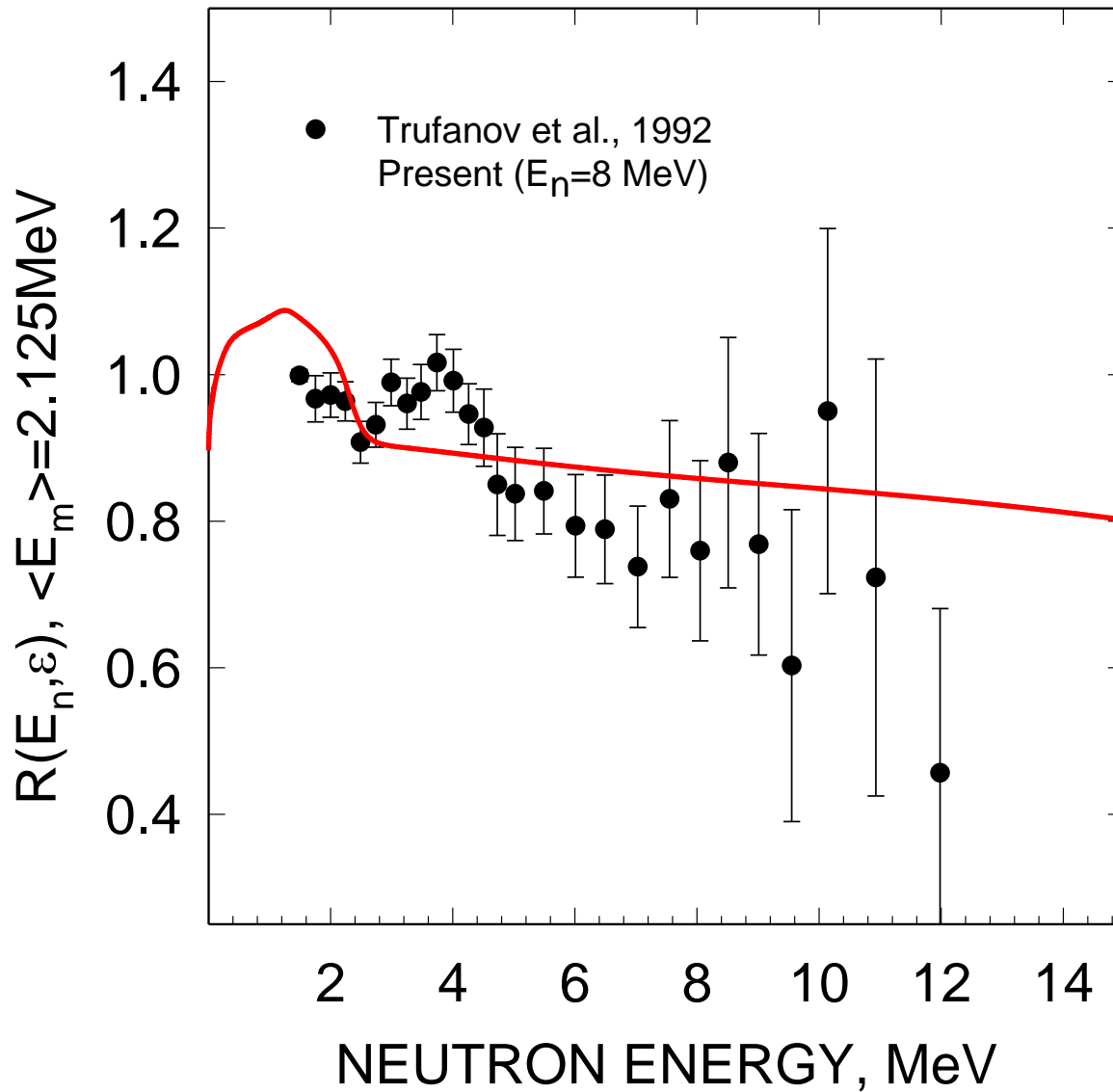
$^{237}\text{Np}$  FISSION NEUTRON SPECTRUM,  
 $E_n=14.7\text{MeV}$



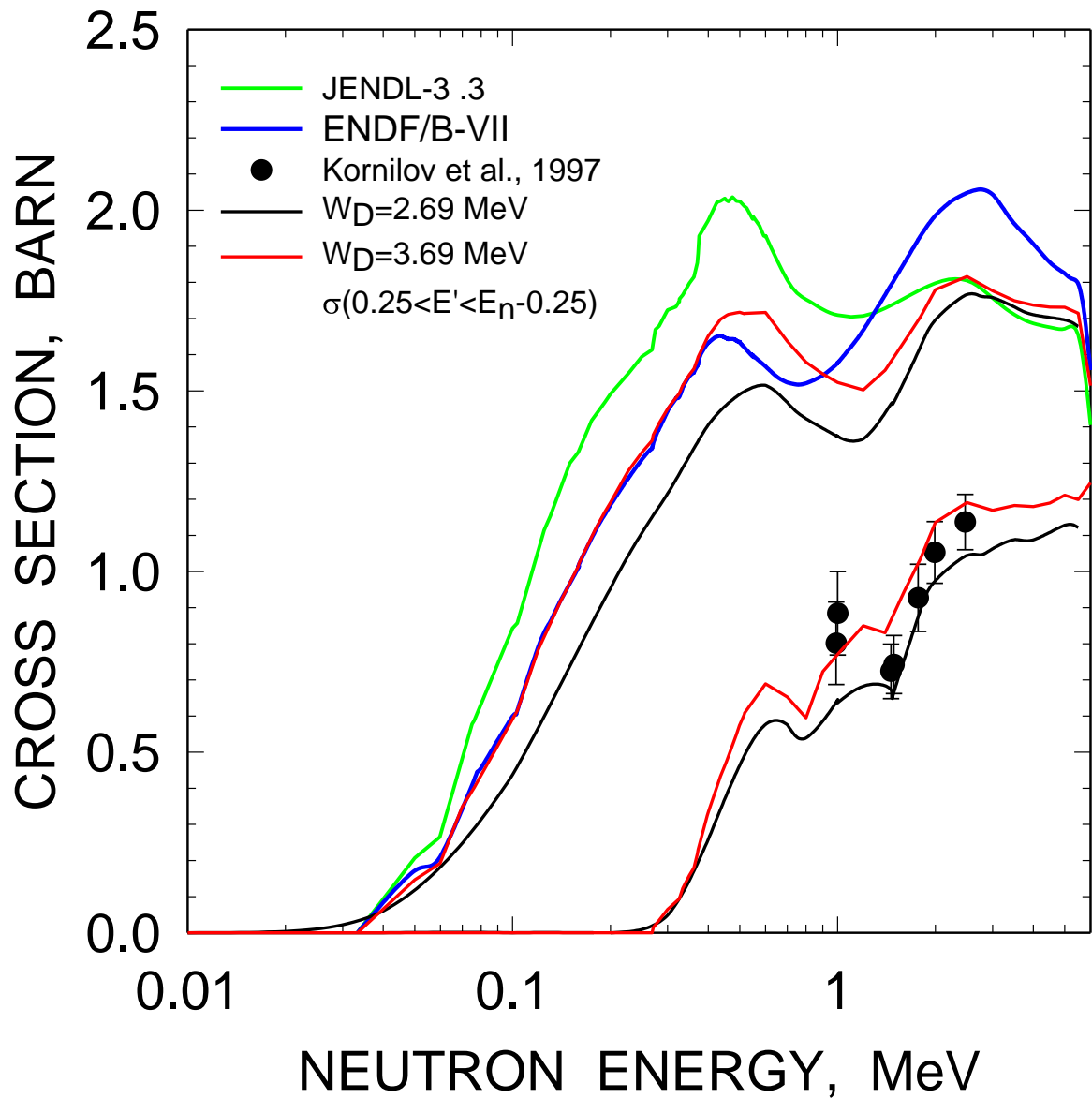
# $^{237}\text{Np}$ , AVERAGE PROMPT FISSION NEUTRON ENERGY



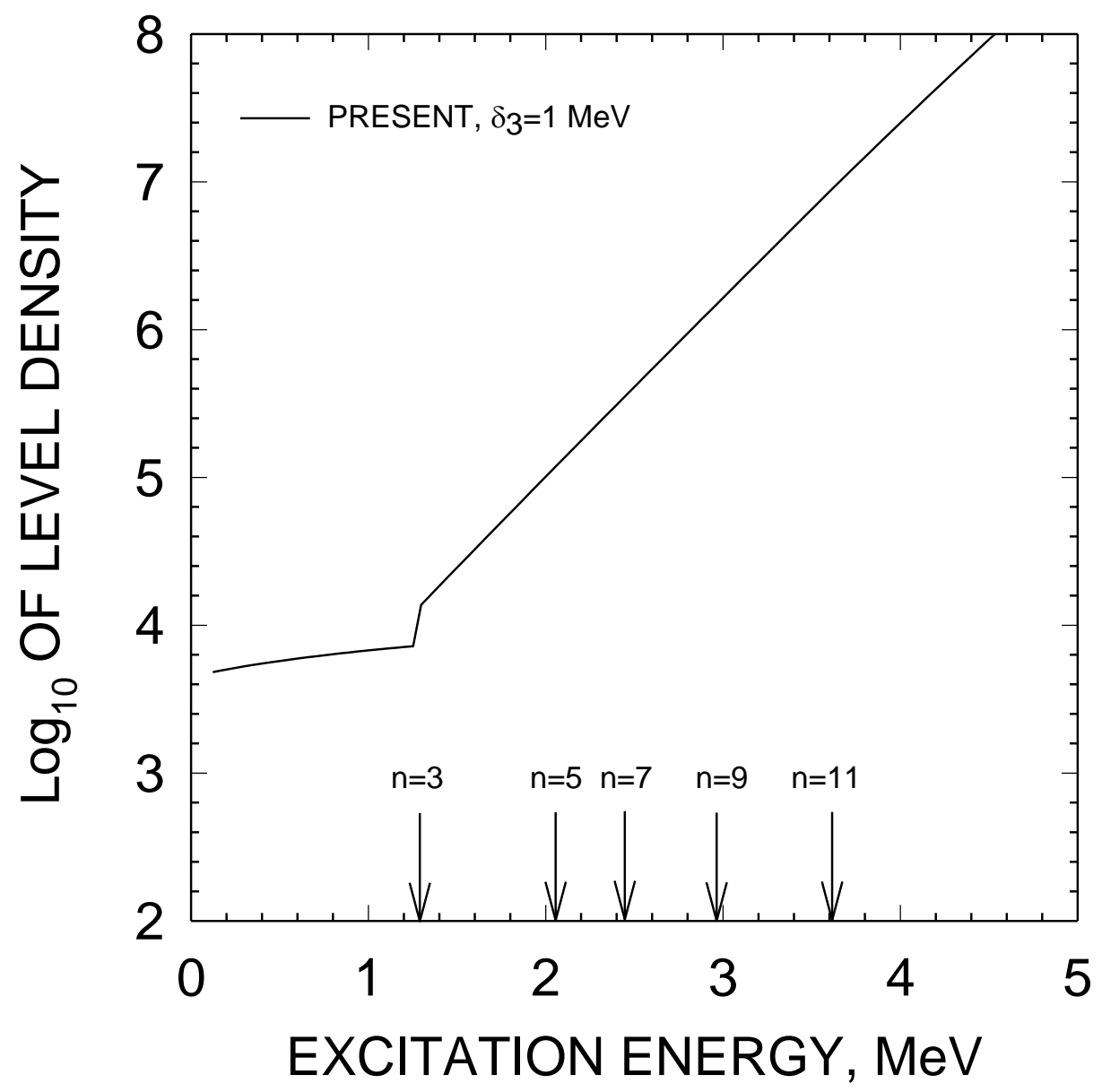
$^{237}\text{Np}$  FISSION NEUTRON SPECTRUM  
 $E_n=7.8$  MeV



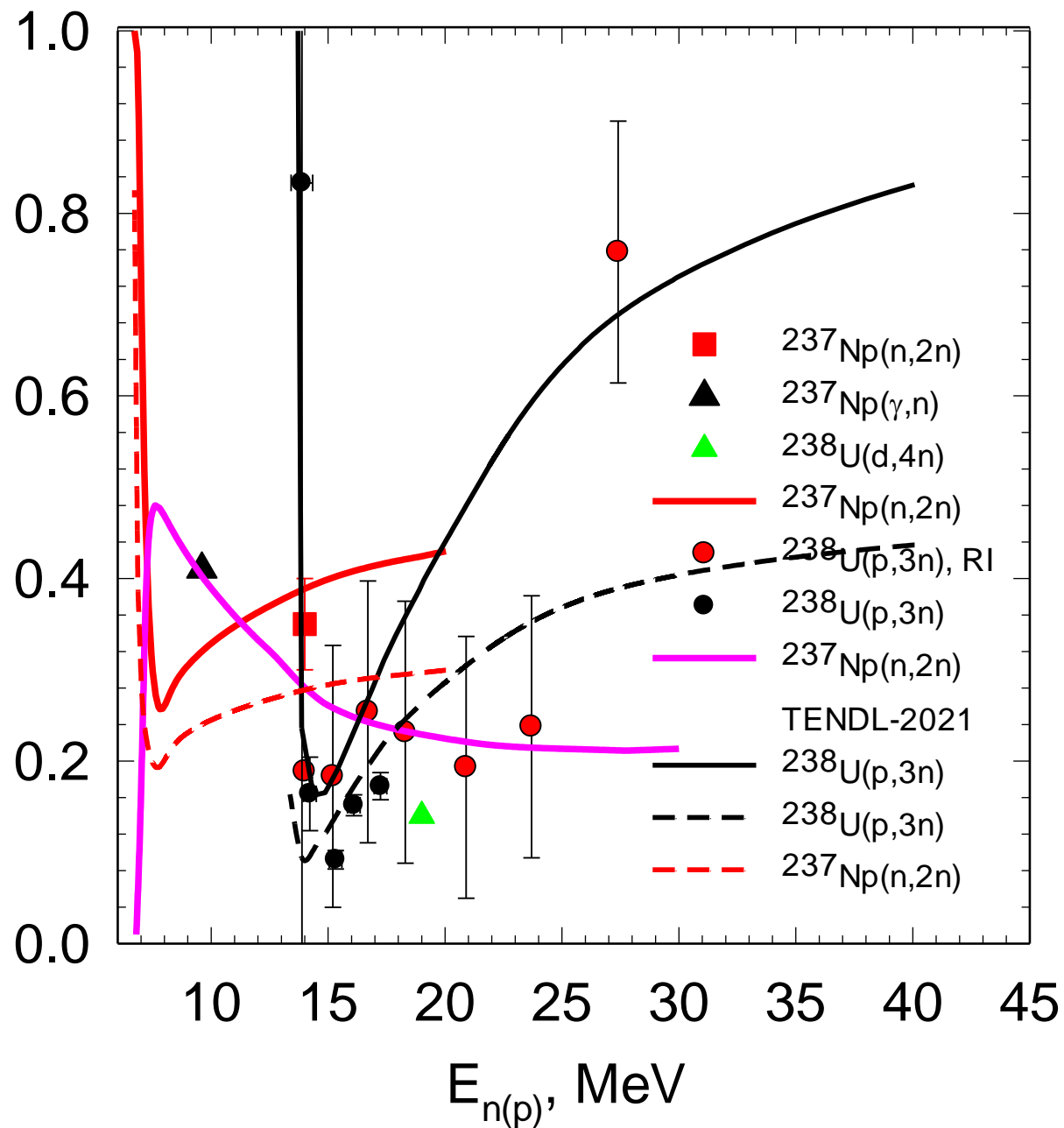
# $^{237}\text{Np}$ INELASTIC CROSS SECTION



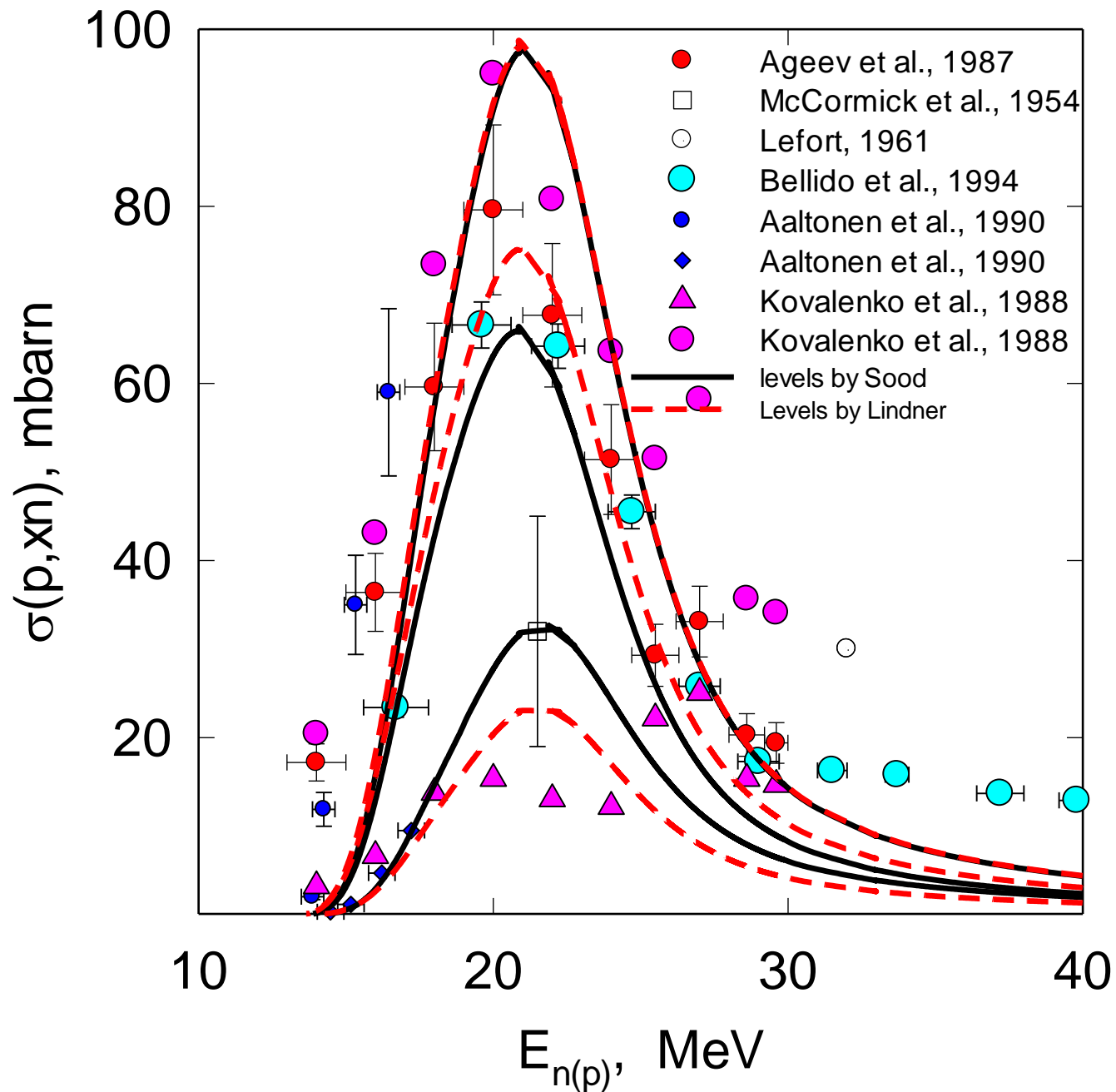
$^{237}\text{Np}$

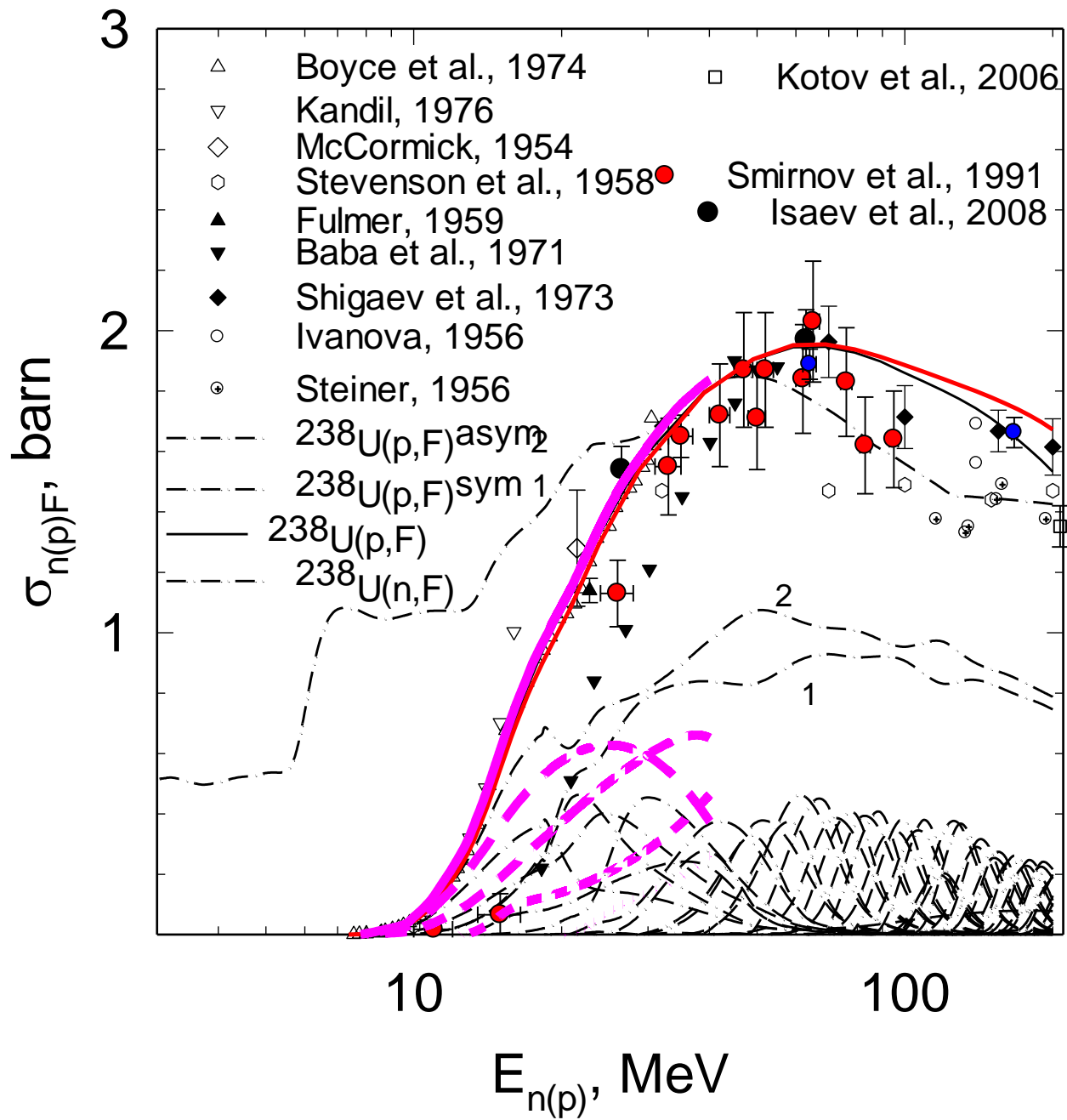


# Isomer Ratio $^{236\text{l}}\text{Np}/^{236\text{s}}\text{Np}$

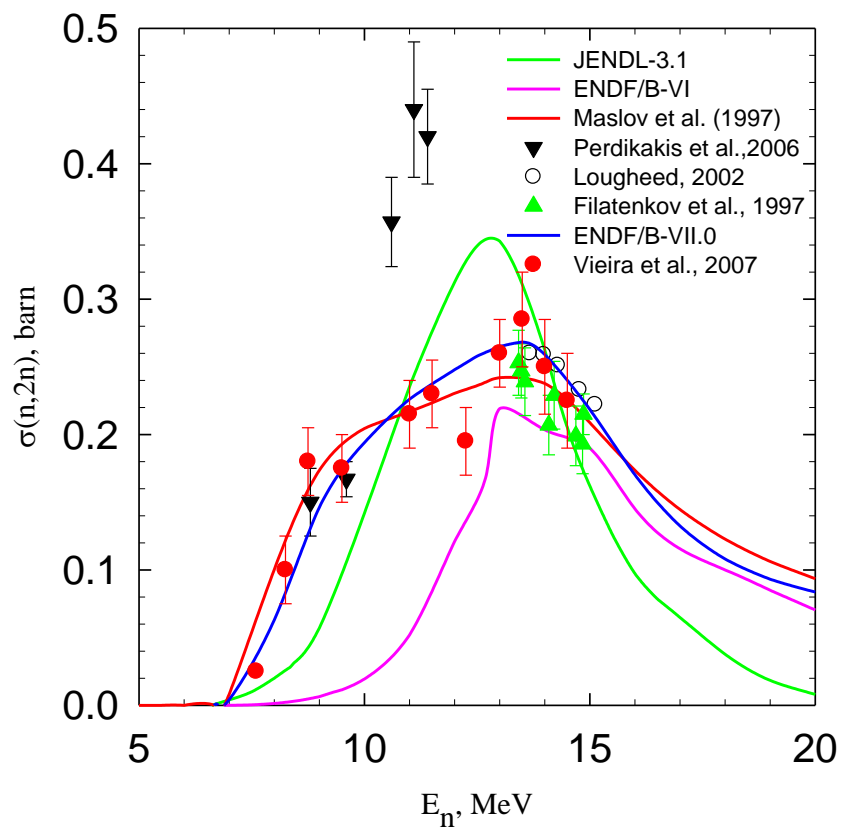




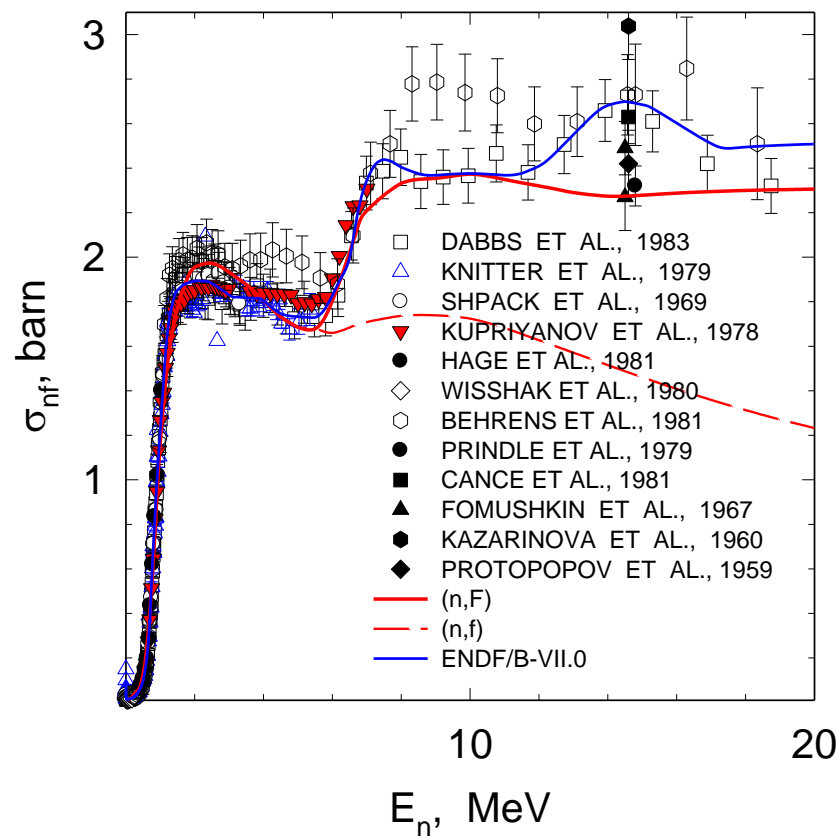
$^{238}\text{U}(p,3n)^{236}\text{Np}$  &  $^{232}\text{Th}(p,3n)$ 



$^{241}\text{Am}(n,2n)$  CROSS SECTION



$^{241}\text{Am}$  FISSION CROSS SECTION



## ISOMER RATIOS FOR $^{241}\text{Am}(n, \gamma)$ AND $^{243}\text{Am}(n, 2n)$ REACTIONS

- Obtained in the same manner

# $^{243}\text{Am}$ FISSION CROSS SECTION

