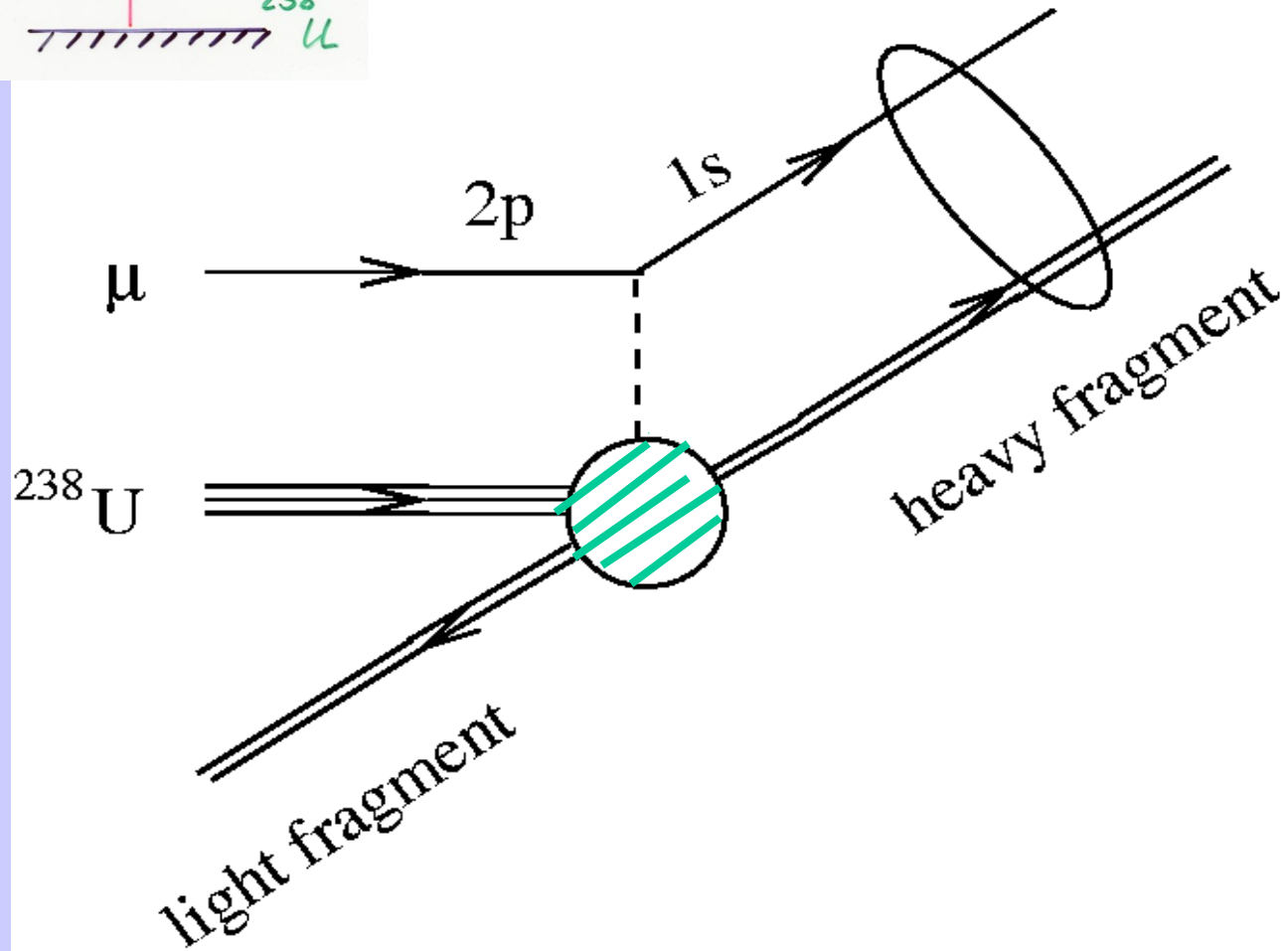
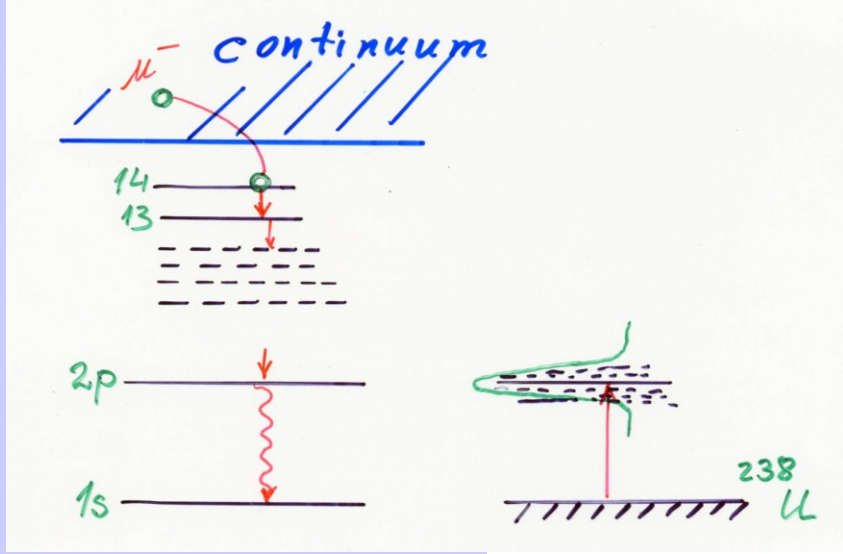


MANIFESTATION OF THE FISSION DYNAMICS IN MUON-INDUCED PROMPT FISSION

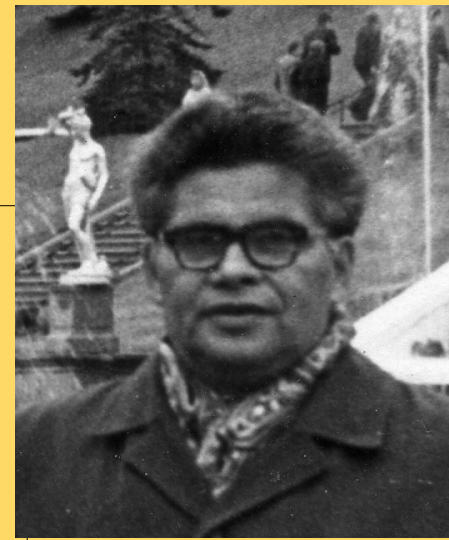
F.F. Karpeshin

D.I.Mendeleev Institute for Metrology, Saint-Petersburg, Russia

- In muonic atoms of ^{238}U , the nuclei can undergo prompt fission through non-radiative transitions of the muon: $2p - 1s$, $3p - 1s$, $3d - 1s$ etc. Main features of the fission dynamics are studied in prompt fission: augmentation of the barrier, dynamics of the saddle-to-scission descent, muonic conversion and characteristic X-rays from fission fragments supply information on the multipolarity of electromagnetic transitions and charge distribution, structure of nuclear transition currents. Revision of the non-radiative transition probabilities comprises the present purpose.



μ -Atom Transitions

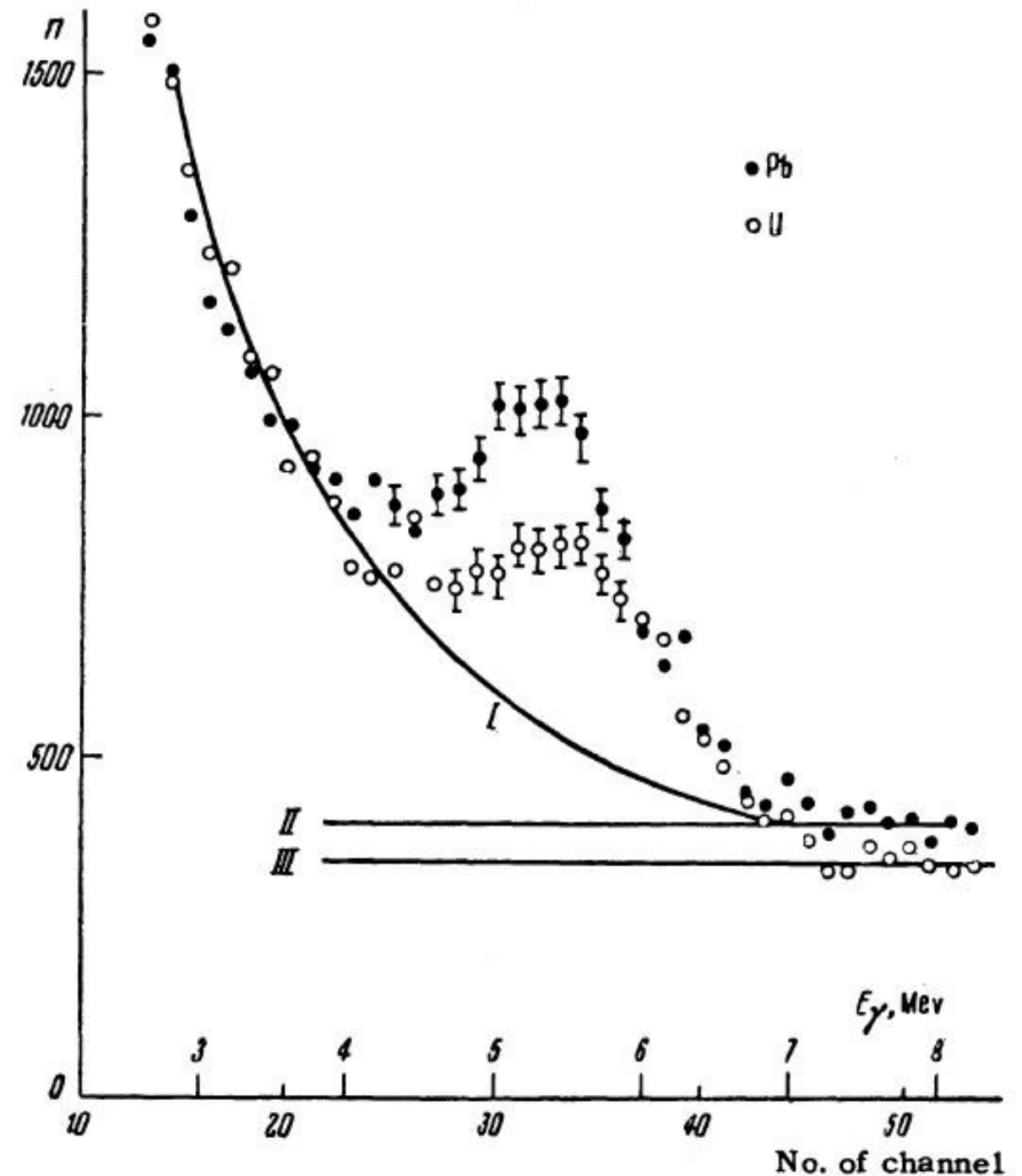


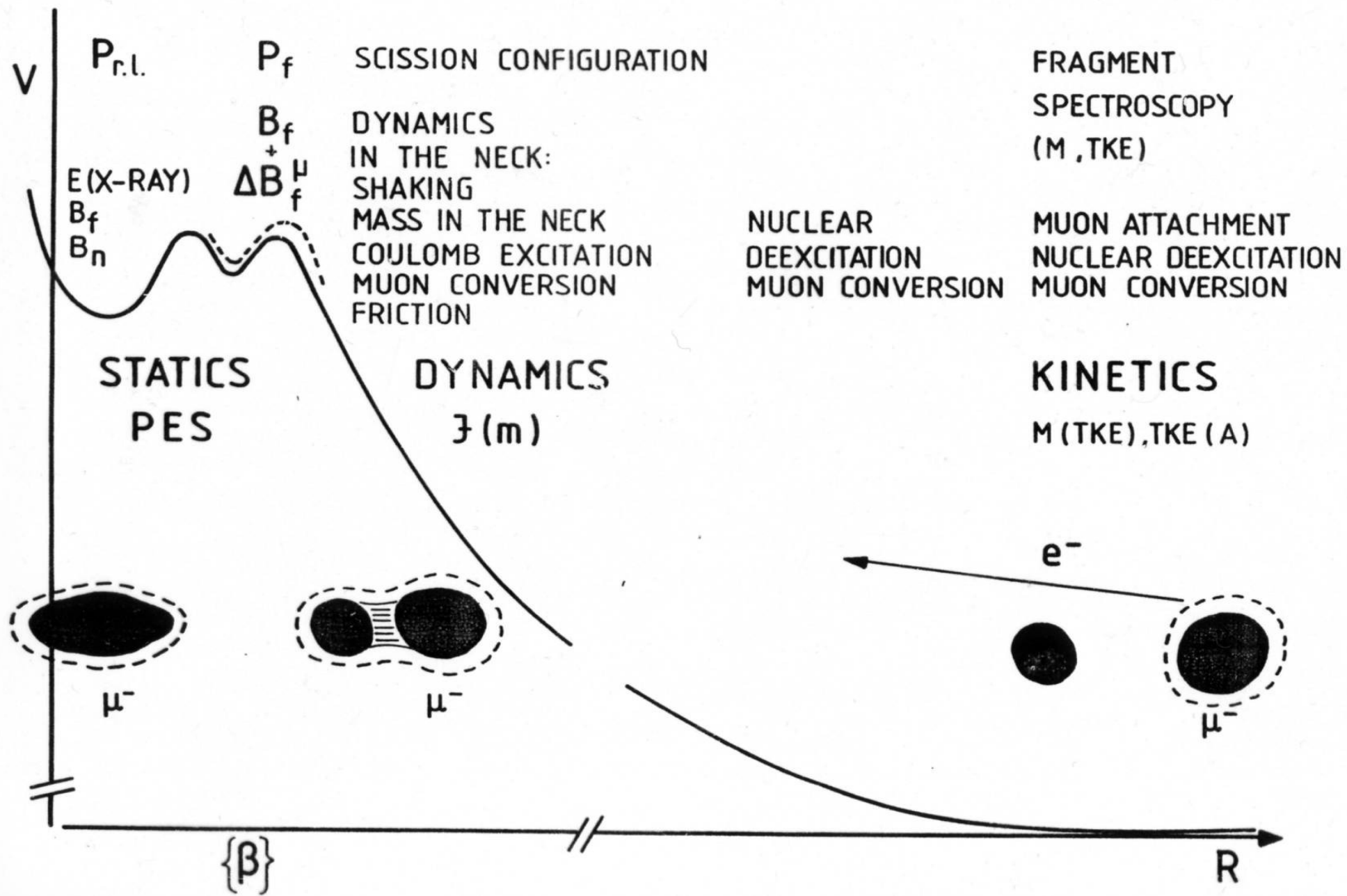
15%

NON-RADIATIVE TRANSITIONS IN MUONIC ^{238}U ATOMS



*M. Ya. Balatz, L. N. Kondrat'ev,
L. G. Landsberg, P. I. Lebedev,
Yu. V. Obukhov, and B. Pontecorvo.*
Soviet Phys. ZETP, **11**, 1239
(1960)





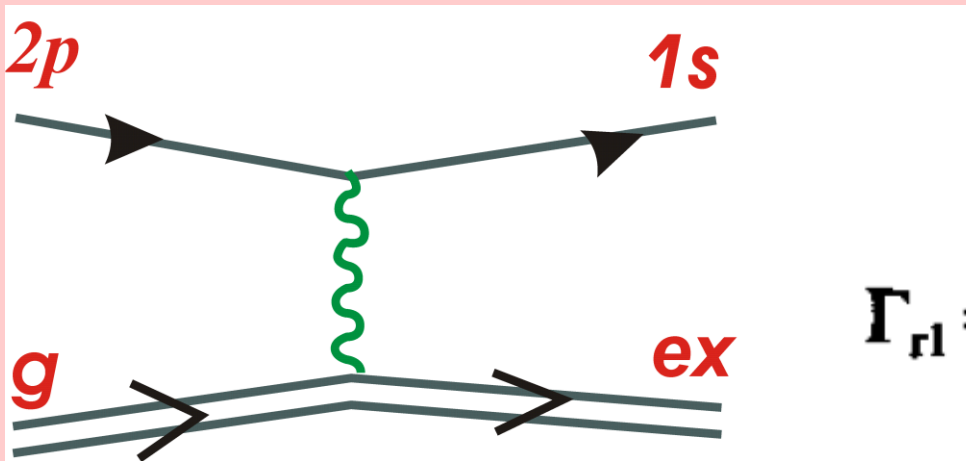


EXPERIMENTAL NON-RADIATIVE TRANSITION PROBABILITIES

*Ch. Roesel, P. David, H. Folger, H. Haenscheid et al.
Radiationless transition probabilities in muonic ^{208}Pb , ^{232}Th ,
and ^{238}U . Z. Phys. A - Hadrons and Nuclei **340**, 199-208 (1991)*

	$2p \rightarrow 1s$	$2p^l \rightarrow 1s$	$2p^h \rightarrow 1s$	$3p \rightarrow 1s$	$3d \rightarrow 1s$
This work	26.2 ± 2.6	$21.6 \pm 3.2(1.6)$	$31.1 \pm 2.8(1.3)$	88.9 ± 4.3	12.8 ± 1.4
Zaretski & Novikov [2, 3]					
with σ_r from [27]	22.4	21.1	24.2	64.7	
with σ_r from [28]	29.8	28.4	32.0	68.5	
Teller & Weiss [7]	20.7	20.0	21.7	59	9.6
Karpeshin & Nesterenko [8]	$11 \div 15^a$			$55 \div 65^c$	$19 \div 24^e$
	$19 \div 26^b$			$57 \div 69^d$	$25 \div 32^f$

MICROSCOPIC NON-RADIATIVE TRANSITION PROBABILITIES



*FF Karpeshin and V O Nesterenko,
J. Phys. B 17, 705 (1991)*

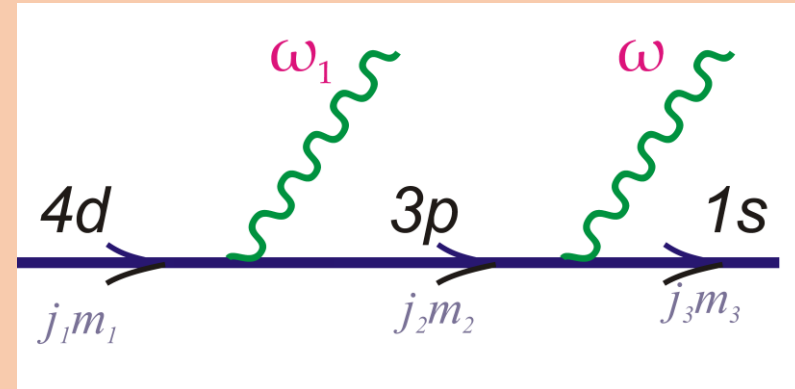
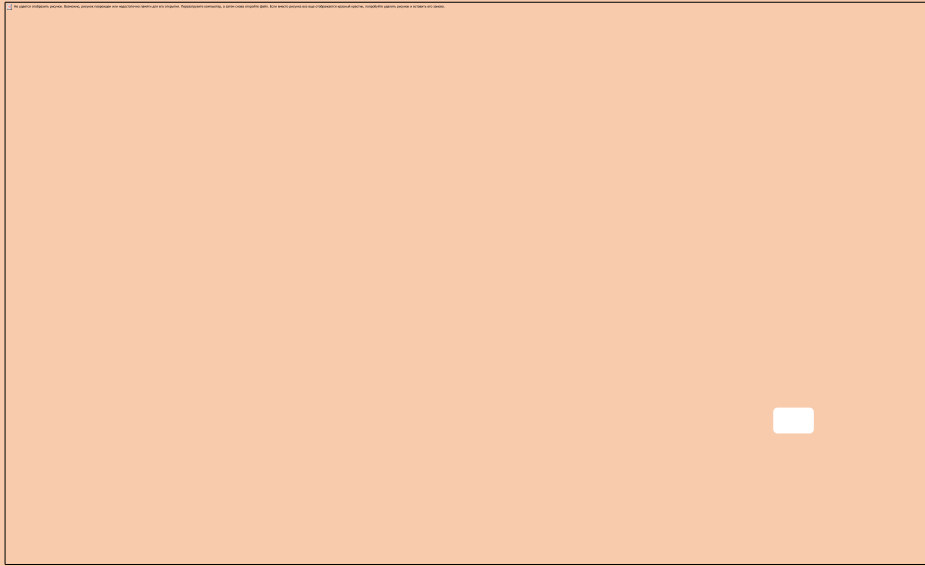
$$\Gamma_{\text{nl}} = \alpha_{\mu}^{(d)}(i \rightarrow f) \frac{8\pi(L+1)}{[(2L+1)!!]^2} \omega^{2L+1} b(\text{EL}; 0 \rightarrow \omega)$$

QPNM:

$$b(\text{EL}; 0 \rightarrow \omega) = \sum_{\mathbf{g}} B(\text{EL}; 0 \rightarrow \omega) \frac{\Delta/2\pi}{(\omega - \omega_{\mathbf{g}})^2 + (\Delta/2)^2}$$

Theory of nonradiative transitions

F.F. Karpeshin, I.M. Band, M.B.Trzhaskovskaya, M.A. Listengarten, Phys. Lett. B 372, 1 (1996).

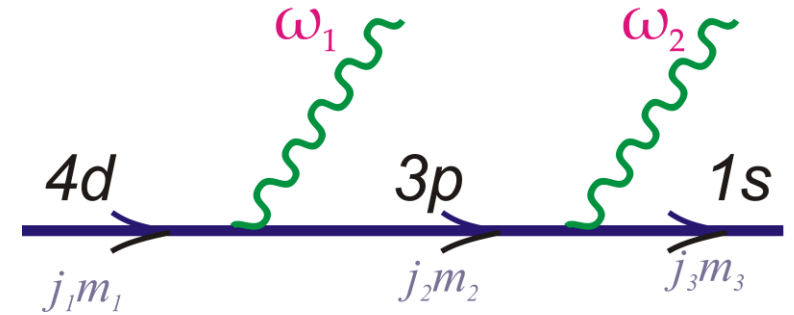


$$F_{nr} = \frac{C(I_2 M_2 \lambda \mu | I_1 M_1)}{\sqrt{2I_1 + 1}} \frac{C(j_2 m_2 \lambda \mu | j_3 m_3)}{\sqrt{2j_3 + 1}} \langle I_2 j_2 || H'_c(\lambda) || I_1 j_3 \rangle \frac{C(j_1 m_1 L M | j_2 m_2)}{\sqrt{2j_2 + 1}} \langle j_2 || H'_\gamma(L) || j_1 \rangle$$

$$\frac{d\Gamma_{nr}(2 \rightarrow 3)}{d\omega_1} = \alpha_d(2 \rightarrow 3; \epsilon_1 - \omega_1) \left(\frac{\omega}{\pi}\right)^2 \sigma_{ph}(\omega) \frac{\Gamma_\gamma(1 \rightarrow 2; \omega_1)/2\pi}{\epsilon_1 - \epsilon_2 - \omega_1)^2 + (\Gamma_2/2)^2}$$

$$\Gamma_{nr}(3p \rightarrow 1s) = \alpha_d(2 \rightarrow 3; \epsilon_1 - \omega_1) \sigma_{ph}(\omega) \left(\frac{\omega}{\pi}\right)^2 \frac{\Gamma_\gamma^{(a)}(1 \rightarrow 2)}{\Gamma_2}$$

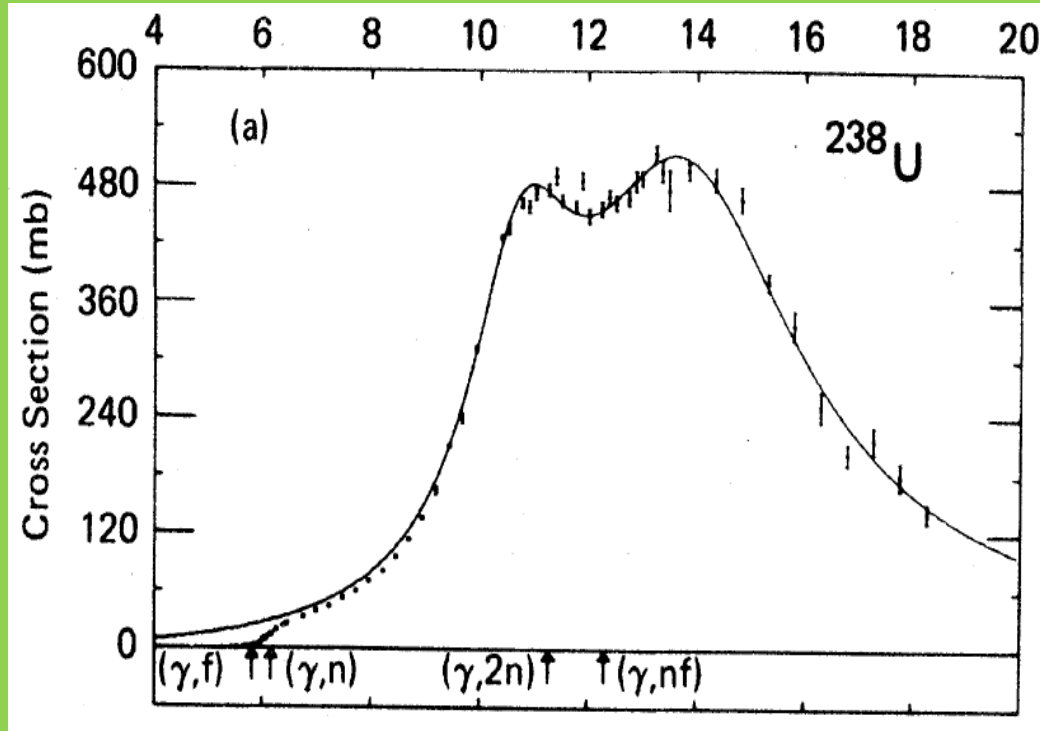
2-photon radiative width:



$$\Gamma_{2\gamma}(1 - 2 - 3) = \Gamma_{\gamma}^{(a)}(1 \rightarrow 2) \frac{\Gamma_{\gamma}^{(a)}(2 \rightarrow 3)}{\Gamma_2}$$

$$\frac{\Gamma_{nr}}{\Gamma_{2\gamma}} = \sigma_{ph}(\omega) \frac{\alpha_d(2 \rightarrow 3; \omega_1)}{\Gamma_{\gamma}^{(a)}(2 \rightarrow 3)} \left(\frac{\omega}{\pi}\right)^2$$

NON-RADIATIVE EXCITATION OF GDR



J. T. Caldwell et al. Phys. Rev. C **21**, 1215 (1980)



Transition	$W_{nr}, \%$		$\Gamma_{nr}/\Gamma_{\gamma}^{(a)}$	
	theory	experim.	theory	experim.
$2p_{3/2} \rightarrow 1s$	31.6	31.1 ± 2.8	0.46	0.45
$3p_{3/2} \rightarrow 1s$	80	88.9 ± 4.3	3.95	5.49 – 13.7



CONCLUSION

- Prompt fission provides multilateral information about fission dynamics. We understand a lot of data results concerning $2p - 1s$ nonradiative transitions. Augmentation of fission barrier and suppression of the fission mode is of great interest.
- Former contradiction between theory and experiment for the $3p - 1s$ transition has been considerably smoothed. At the same time, such a factor like suppression of the fission decay channel in muonic atoms is not involved in the analysis. This channel comprises as much as 20% of the total decay width. Such ruling out a decay channel may increase the photoabsorption cross-section by the same factor. This also depends on the decay-to-spread width. Therefore, this circumstance remains challenging for future investigations.

Thanks!